## Superconducting

Magnet Division

## Lecture IV

## Magnetic Design Coil Optimization

Ramesh Gupta<br>Superconducting Magnet Division<br>Brookhaven National Laboratory<br>US Particle Accelerator School<br>University of California - Santa Barbara June 23-27, 2003

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## Coil Designs for Real Magnets

 Magnet DivisionLHC Dipole


RHIC Dipole

contanvent vessel

Tevatron Dipole


Figure 4.9: The Tevatron 'warm-iron' dipole (Tollestrup 1979).

HERA Dipole

-All magnets use
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- All designs use cosine theta coil geometry

Do they really look like having a cosine theta current distribution?

Or that matter, even an elliptical geometry for conductor having a constant current density?

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## Coil Cross-section Optimization (1)

The optimization of a coil cross-section for a good magnetic design involves:

- Minimizing field harmonics
- Maximizing field (Transfer Function) for a lower number of turns
- Minimizing Peak Field (Max. field on the conductor for given central field)

At first, it appears to be fairly straight forward process, thanks, in part, to the modern automated codes like ROXIE and PAR2DOPT, etc.,
In fact, one can built a magnet based on the optimized coil structure obtained by a relatively new user.

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## Coil Cross-section Optimization (2)

But the advanced cross-section optimization is a bit more involved:

- One must avoid designs that create mechanical difficulties
- One should look for flexibility to allow future adjustments
- Also look for special requirements in each application.

One approach fits all, may not always be a good strategy.

## My Experience:

-The initial design, quiet often sets, the eventual performance of the magnet apart from degree of difficulties in manufacturing the magnet.
-As compared to building magnets, design process takes a relatively small resources. Spend time in looking for as many possibilities/options as possible.

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## Field Harmonic Definitions

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The field quality in magnets is expressed in terms of the normal and skew harmonic coefficients, $b_{n}$ and $a_{n}$ by the following expansion:

$$
B_{y}+i B_{x}=10^{-4} \times B_{R} \sum_{n=1}^{\infty}\left(b_{n}+i a_{n}\right)[(x+i y) / R]^{n-1}
$$

where $x$ and $y$ are the horizontal and vertical coordinates, $B_{\mathrm{R}}$ is the field strength of the primary harmonics at the "reference radius" $R$. The values of the field harmonic are given in the units of $10^{-4}$.

The definition used above (European convention) differs from that used in many U.S. publications (US convention), where $n-1$ is replaced by $n$ and the summation starts from $n=0$.

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## Field and Harmonics From A Current Block (No Iron)

Consider a "Radial Block" between radii $\rho_{1} \& \rho_{2}$ and angle $\phi_{1} \& \phi_{2}$ having a current density of $J$. The total current $(I)$ and harmonics $\left(A_{n}, B_{n}\right)$ are given by:

$$
I=\frac{1}{2} J\left(\rho_{2}^{2}-\rho_{1}^{2}\right)\left(\phi_{2}-\phi_{1}\right)
$$

$$
\begin{aligned}
& A_{1}=-\frac{\mu_{o} J}{2 \pi}\left(\rho_{2}-\rho_{1}\right)\left[\cos \left(\phi_{2}\right)-\cos \left(\phi_{1}\right)\right], \\
& A_{2}=-\frac{\mu_{o} J R_{o}}{2 \pi} \ln \left(\frac{\rho_{2}}{\rho_{1}}\right)\left[\cos \left(2 \phi_{2}\right)-\cos \left(2 \phi_{1}\right)\right],
\end{aligned}
$$

$$
n=1 \text { refers to dipole. }
$$

$R_{o}$ is the reference radius.
for $\mathrm{n} \geq 3$

$$
A_{n}=\frac{\mu_{o} J}{2 \pi} \frac{R_{o}^{n-1}}{n(n-2)}\left(\frac{1}{\rho_{2}^{n-2}}-\frac{1}{\rho_{1}^{n-2}}\right)\left[\cos \left(n \phi_{2}\right)-\cos \left(n \phi_{1}\right)\right]
$$

$$
\begin{aligned}
& B_{1}=-\frac{\mu_{o} J}{2 \pi}\left(\rho_{2}-\rho_{1}\right)\left[\sin \left(\phi_{2}\right)-\sin \left(\phi_{1}\right)\right], \\
& B_{2}=-\frac{\mu_{o} J R_{o}}{2 \pi} \ln \left(\frac{\rho_{2}}{\rho_{1}}\right)\left[\sin \left(2 \phi_{2}\right)-\sin \left(2 \phi_{1}\right)\right],
\end{aligned}
$$

See my thesis for these and other derivations.
for $\mathrm{n} \geq 3$

$$
B_{n}=\frac{\mu_{o} J}{2 \pi} \sum_{n=3}^{\infty} \frac{R_{o}^{n-1}}{n(n-2)}\left(\frac{1}{\rho_{2}^{n-2}}-\frac{1}{\rho_{1}^{n-2}}\right)\left[\sin \left(n \phi_{2}\right)-\sin \left(n \phi_{1}\right)\right]
$$

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Assume that a current block is between radii 10 cm and 12 cm and it starts at an angle $\Phi=0$.
Find the subtended angle (or the angle at which the block must end) for

- The normal sextupole harmonic to be zero. Is it a unique solution?
- Assume that you have a dipole symmetry. How many other blocks must be present to generate this symmetry (give number and angles).
- Compute the values of the first three non-zero allowed harmonics in Tesla ( $A_{n}$ and/or $B_{n}$ ) at a reference radius of $6 \mathbf{c m}$.
- Do the above for decapole harmonic to be zero. Is it a unique solution?
- Can you find a solution for which both sextupole and decapole harmonics are zero? The block does not have to start from an angle $\Phi=0$. What happens if it is in a cylindrical iron cavity having a permeability of (a) $\mathbf{1 0}$, (b) 100 and (c) 1000. Hint: You can use the method of images.


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## Homework Assignment (Two blocks)

Assume that there are two current blocks. First between radii 10 cm and 11 cm and second between 11 cm and 12 cm . Both starts at angle $\Phi=0$.

Find the subtended angle (the angle at which block must end) for the normal sextupole harmonic and decapole harmonics to be zero.

Is it a unique solution?

What happens if it is in a cylindrical iron cavity having a permeability of
(a) 10
(b) 100 and
(b) 1,000.

Hint: You can use the method of images.

## How to Look for Optimal X-section (1)

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1. Look for designs that look similar to cosine theta distribution (experience)
2. Use special techniques/software (artificial intelligence)
3. Cover a large range of combinations and find the best

My recommendation, go for the $3^{\text {rd }}$ option (several thousand cases):

It does not take long to look a large number of possibilities - less than a few seconds per case if the peak field is not computed.

To save time compute peak field only in promising solutions.

## How to Look for Optimal X-section (2)

Develop a front-end program to automatically create several cases for a series optimization run.
In this optimization:

- Vary number of blocks and number of turn in those blocks.
- Vary starting condition of wedges, etc.

Post-process results to select a limited cases with filters for harmonics, etc.

- Compute peak field for these selected cases.


Go back and carefully evaluate and further optimize these few cases for performance, mechanical layout, flexibility, sensitivity, efficiency and any other requirement.

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## Superconducting Magnet Division <br> A D07GEN Run

## An Example of Pre-processor for Generating Cases for a Series Run


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30.


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## An Example of Post-processor for Selecting Cases

## A PARSLCT Run



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Table 4.3.1: The effect of a $+25 \mu \mathrm{~m}$ change in a wedge thickness or pole width on the transfer function and the field harmonics in the SSC 50 mm aperture dipole magnet. The field harmonics are calculated with a 10 mm reference radius. The numbering of the wedges is counter-clockwise from the midplane. The pole width is measured from the vertical axis.

| Parameter | $\delta T F$ | $\delta b_{2}$ | $\delta b_{4}$ | $\delta b_{6}$ |
| :---: | :---: | :---: | :---: | :---: |
| changed | $10^{-4} \frac{T}{k A}$ | $10^{-4}$ | $10^{-4}$ | $10^{-4}$ |
| Wedge No. 1 | -0.78 | -0.24 | 0.01 | 0.005 |
| Wedge No. 2 | 0.42 | 0.30 | 0.03 | -0.005 |
| Wedge No. 3 | 1.16 | 0.36 | -0.02 | 0.00 |
| Wedge No. 4 | -0.29 | -0.06 | 0.00 | 0.00 |
| Pole Width (inner) | 2.0 | 0.23 | -0.03 | 0.005 |
| Pole Width (outer) | 1.13 | 0.21 | 0.00 | 0.000 |

Wedge No. 4 is in outer layer.

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Table 4.3.2: The computed change in the transfer function and field harmonics produced by a $+25 \mu \mathrm{~m}$ ( $0.001^{\prime \prime}$ ) change in the wedge thickness, pole width or midplane gap in the RHCC 80 mm aperture arc dipoles. The field harmonics are calculated with a 25 mm reference radius. The numbering of the wedges starts at the midplane. The pole width and midplane gap are measured from the vertical and horizontal axis, respectively.

| Parameter <br> changed | $\delta T F$ <br> $10^{-4} \frac{T}{k A}$ | $\delta b_{2}$ <br> $10^{-4}$ | $\delta b_{4}$ <br> $10^{-4}$ | $\delta b_{6}$ <br> $10^{-4}$ | $\delta b_{8}$ <br> $10^{-4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Wedge 1 | -0.6 | -0.98 | -0.122 | 0.061 | 0.043 |
| Wedge 2 | 0.1 | 0.69 | 0.423 | 0.022 | -0.050 |
| Wedge 3 | 1.1 | 1.42 | -0.090 | -0.068 | 0.041 |
| Pole Width | 1.7 | 1.11 | -0.154 | 0.039 | -0.014 |
| Midplane Gap | -0.9 | -1.68 | -0.557 | -0.156 | -0.0 .50 |

Notice that the magnitude of change in harmonics by a $25 \mu$ change, is much large in RHIC dipole than in SSC dipole. WHY?

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## Change in Harmonics for 0.1 mm Change in Dimension in DO Dipole

RHIC Insertion Dipole D0 with single layer coil Coil inner radius $=\mathbf{1 0 0} \mathbf{~ m m}$,
Harmonic reference radius $=\mathbf{6 5 m m}$
"D0 MAGNET" : Rough Cross section Optimization Spread sheet


To iterate cross section for b2 \& b4 with midplane and pole shims and fixed wedge changes, go to line b65:b72

| Expected Parameters of the Iterated Design |  |  |  |  |  | b2 | b4 | b6 | b8 | b10 | b12 | b14 | b16 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Coil Prestress | 1 | mil from the target. |  | Expected Harmonic |  | 6.23 | 2.24 | 0.54 | 0.240 | -0.012 | 0.132 | -0.047 | -0.099 |
| Coil Size |  |  |  |  |  | db2 | db4 | db6 | db8 | db10 | db12 | db14 | db16 |
| Target Increase | 5 | in Coil size | Mutiply | Change(mil) | Target b2n | -8.00 | -2.30 | -0.16 | -0.15 | 0.016 | -0.15 | 0.039 | 0.096 |
| Total Increase | 6 | mil in compression |  |  |  | -1.77 | -0.06 | 0.38 | 0.09 | 0.00 | -0.02 | -0.01 | 0.00 |
| Midplane(mil) | 1 | Fixed Pole | 0.25 | 4 | -0.00333 | -5.79 | -1.85 | -0.48 | -0.16 | -0.06 | -0.02 | -0.01 | 0.00 |
| PoleShim(mil) | 1 | Decrease Pole | 0.25 | 4 | 0.006 | 3.98 | -0.53 | 0.09 | -0.05 | 0.01 | 0.00 | 0.01 | 0.00 |
| Wedge1(mil) | 4 | Fixed Pole | 1 | 4 | -0.0012 | -1.32 | 0.53 | 0.48 | 0.14 | 0.02 | -0.01 | -0.01 | 0.00 |
| Wedge2(mil) | 0 | Fixed Pole | 0 | 4 | 0.00119 | 3.05 | 1.28 | -0.06 | -0.16 | -0.01 | 0.02 | 0.01 | 0.00 |
| Wedge3(mil) | 0 | Fixed Pole | 0 | 4 | 0.00314 | 4.80 | 0.30 | -0.37 | 0.02 | 0.03 | -0.01 | 0.00 | 0.00 |
| Wedge4(mil) | 0 | Fixed Pole | 0 | 4 | 0.00508 | 4.89 | -0.65 | -0.01 | 0.04 | -0.04 | 0.01 | 0.0 | 0.0 |
|  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| Midplane +4 mil | 0 | Increase Pole | 0 | 4 | -0.0097 | -9.76 | -1.31 | -0.5 | -0.11 | -0.07 | -0.02 | -0.0 | 0.00 |
|  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| Wedge1+4mil | 0 |  | 0 | 4 | -0.00758 -0.0052 | $\begin{array}{r}-5.28 \\ -0.91 \\ \hline\end{array}$ | 1.06 | 0.389 -0.155 | $\begin{array}{r}0.187 \\ -0.117 \\ \hline\end{array}$ | $\begin{array}{r}0.006 \\ -0.018 \\ \hline\end{array}$ | -0.006 | -0.013 <br> -0.002 | 0 |
| Wedge2+4mil | 0 | Increase Pole | 0 | 4 | -0.0052 -0.00323 | -0.91 | 1.81 0.82 | -0.155 | -0.117 | -0.018 | -0.025 | -0.002 | 0.004 |
| Wedge4+4mil | 0 | Increase Pole | 0 | 4 | -0.0013 | 4.07 | -1.37 | 0.439 | -0.127 | 0.174 | -0.184 | -0.057 | 0 |

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## A Flexible Design from the Beginning

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## Design Philosophy:

- Start out a design from the beginning itself, that allows significant adjustability for field harmonics and mechanical parameters (cable thickness, wedges, etc.).
- A flexible design is generally economical, efficient and produces magnets with better performance. I think it's a prudent approach.


Geometric: Start with a larger than required shim and midplane cap. Then adjust it, as required without changing the cross-section of the cured coil. One can also adjust the layers of wedge/cable insulation, if needed. These three parameters can adjust, first two allowed harmonics and pre-stress or cable insulation. This approach was used extensively in various RHIC magnets.

Saturation: Start out with holes and fill them with iron rods. Or, punch holes in laminations later.
In KEK, C.F. dipole, one may consider adjusting for quad and sextupole component.

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## Change in Midplane Gap to Adjust Harmonics

 (can be easily done by changing the size of the ground-plane insulation cap)

## Coil-to-coil

 midplane gap (1 in dipole)Table 4.4.1: The computed and measured change in field harmonics at 25 mm reference radius due to a change in the coil midplane gap. The midplane gap was increased from 0.114 mm to 0.16 mm in the rebuilt 80 mm aperture RHIC model dipole magnet DRS009. In the production magnets, the midplane gap was changed back to 0.114 mm from 0.16 mm to adjust the $b_{4}$ harmonic.


|  | $\Delta b_{2}$ | $\Delta b_{4}$ | $\Delta b_{6}$ | $\Delta b_{8}$ |
| :--- | :---: | :---: | :---: | :---: |
| Computed | -3.0 | -1.0 | -0.28 | -0.09 |
| Measured | -3.0 | -1.0 | -0.29 | -0.12 |

Table 4.5.1: The measured and computed change in field harmonics caused by an asymmetric increase in the coil-to-midplane gap in the prototype 130 mm aperture RHIC interaction quadrupole QRI002. The gap was increased by 0.1 mm in the horizontal plane only. The harmonics are given at a reference radius of 40 mm .

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RHIC 100 mm Aperture Insertion Dipole: The first magnet gets the body harmonics right

## Geometric Field Errors on the X-axis of DRZ101 Body

First magnet and first attempt in RHIC 100 mm aperture insertion dipole A number of things were done in the test assembly to get pre-stress \& harmonics right

Harmonics at $\mathbf{2 k A}$ (mostly geometric).


Note: Field errors are within $10^{-4}$ at $60 \%$ of coil radius and $\sim 4^{* 1} 10^{-4}$ at $80 \%$ radius.

Later magnets had adjustments for integral field and saturation control. The coil cross-section never changed.

Measured in 0.23 m long straigth section.

Reference radius $=31 \mathrm{~mm}$

| b1 | -0.39 | a2 | -1.06 |
| :---: | :---: | :---: | :---: |
| b2 | -0.39 | a3 | -0.19 |
| b3 | -0.07 | a4 | 0.21 |
| b4 | 0.78 | a5 | 0.05 |
| b5 | -0.05 | a6 | -0.20 |
| b6 | 0.13 | a7 | 0.02 |
| b7 | -0.03 | a8 | -0.16 |
| b8 | 0.14 | a9 | -0.01 |
| b9 | 0.02 | a10 | 0.01 |
| b10 | -0.04 | a11 | -0.06 |
| b11 | 0.03 | a12 | -0.01 |
| b12 | 0.16 | a13 | 0.06 |
| b13 | -0.03 | a14 | 0.03 |
| b14 | -0.10 | a15 | 0.02 |

All harmonics are within or close to one sigma of RHIC arc dipoles.

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## Average Field errors ~10-4 up to $80 \%$ of the coil radius

Geometric Field Errors on the X-axis of RHIC DRZ magnets (108-125)
Coil X-section was not changed between $1^{\text {st }}$ prototype and final production magnet A Flexible \& Experimental Design Approach Allowed Right Pre-stress \& Right Harmonics
 Estimated Integral Mean in Final Set (Warm-cold correlation used in estimating) Harmonics at 3 kA (mostly geometric)
Reference radius is 31 mm (Coil 50 mm )

| b1 | -0.28 | a1 | -0.03 |
| :---: | :---: | :---: | :---: |
| b2 | -0.26 | a2 | -3.36 |
| b3 | -0.07 | a3 | 0.03 |
| b4 | 0.15 | a4 | 0.48 |
| b5 | 0.00 | a5 | 0.04 |
| b6 | 0.32 | a6 | -0.24 |
| b7 | 0.00 | a7 | 0.01 |
| b8 | -0.08 | a8 | 0.05 |
| b9 | 0.00 | a9 | 0.00 |
| b10 | -0.12 | a10 | -0.02 |
| b11 | 0.03 | a11 | -0.01 |
| b12 | 0.16 | a12 | 0.06 |
| b13 | -0.03 | a13 | 0.03 |
| b14 | -0.10 | a14 | 0.02 |

*Raw Data Provided by Animesh Jain at BNL
*Field errors are $10^{-4}$ to $80 \%$ of the aperture at midplane. ${ }^{*}$
(Extrapolation used in going from 34 mm to 40 mm ; reliability decreases)

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## Computed Peak Fields in SSC Dipole

Table 6.2.6: Peak fields in the SSC 50 mm dipole as computed


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## Computed Performance of SSC Dipole

Table 6.2.7: Expected quench performance of the SSC 50 mm dipole with $5 \%$ cable degradation $\left(J_{c}=2612.5 \mathrm{~A} / \mathrm{mm}^{2}\right)$ and at 4.35 K temperature. $S_{\text {quench }}$ is the computed current density in the copper at quench and $S_{6.7 T}$ at the design field of 6.7 Tesla.

| Layer | $\mathrm{Cu} / \mathrm{Sc}$ | $B_{s s}$ | $I_{c}$ | $B_{\text {margin }}$ | $T_{\text {margin }}$ | $S_{\text {quench }}$ | $S_{6.7 T}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\downarrow$ | Ratio | tesla | $A$ | \%over 6.7T | kelvin | $A / \mathrm{cm}^{2}$ | $A / \mathrm{cm}^{2}$ |
| Inner | 1.7 | 7.149 | 7126 | 6.7 | 0.519 | 736 | 681 |
|  | 1.5 | 7.273 | 7273 | 8.6 | 0.625 | 788 | 715 |
|  | 1.3 | 7.399 | 7411 | 10.4 | 0.730 | 853 | 759 |
| Outer | 2.0 | 7.268 | 7267 | 8.7 | 0.580 | 919 | 834 |
|  | 1.8 | 7.445 | 7470 | 11.1 | 0.709 | 980 | 865 |

the magnet will quench at the design central field ( $B_{\text {design }}=6.7$ tesla). The field margin is defined as follows

$$
B_{\operatorname{margin}}(\%)=\frac{B_{s s}-B_{d e s i g n}}{B_{d e s i g n}} \times 100
$$

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## Error in Parts and Influence in Field Harmonics in SSC Dipole

Table 6.2.8: The effect of a 0.05 mm increase in the given parameter on the transfer function and the field harmonics.

| Parameter | TF | $b_{2}$ | $b_{4}$ | $b_{6}$ |
| :---: | :---: | :---: | :---: | :---: |
| changed | $\mathrm{T} / \mathrm{kA}$ | $10^{-4}$ | $10^{-4}$ | $10^{-4}$ |
| Radius of Block No. 1 | 0.31 | -0.25 | -0.10 | -0.01 |
| Radius of Block No. 2 | -0.32 | 0.31 | 0.12 | 0.01 |
| Radius of Block No. 3 | -0.12 | 0.36 | -0.02 | -0.01 |
| Radius of Block No. 4 | -0.20 | 0.33 | -0.08 | 0.01 |
| Radius of Block No. 5 | -0.11 | -0.04 | -0.01 | 0.00 |
| Radius of Block No. 6 | -0.78 | 0.22 | 0.03 | 0.00 |
| RMS Blocks | 0.38 | 0.27 | 0.07 | 0.01 |


| Thickness of Wedge No. 1 | -1.56 | -0.48 | 0.02 | 0.01 |
| :---: | :---: | :---: | :---: | :---: |
| Thickness of Wedge No. 2 | 0.83 | 0.59 | 0.05 | -0.01 |
| Thickness of Wedge No. 3 | 2.32 | 0.71 | -0.04 | 0.00 |
| Thickness of Wedge No. 4 | -0.57 | -0.11 | 0.00 | 0.00 |
| RMS Wedges | 1.48 | 0.52 | 0.03 | 0.01 |
| Cable thickness inner | 2.63 | 1.08 | 0.05 | -0.01 |
| Cable thickness outer | 1.99 | 0.48 | 0.02 | 0.00 |
| RMS Cable thickness | 2.33 | 0.83 | 0.04 | 0.01 |
| Pole angle inner | -4.01 | -0.45 | 0.06 | -0.01 |
| Pole angle outer | -2.26 | -0.42 | 0.00 | 0.00 |
| RMS Pole angles | 3.25 | 0.43 | 0.04 | 0.01 |

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$$
\text { Stored Energy }=\frac{1}{2} \text { Inductance } \times(\text { Current })^{2}
$$

The inductance decreases at high field as the iron yoke saturates.
The results of POISSON computations for the SSC 50 mm aperture dipole are given at 6.5 kA in Table 6.2 .9 for the stored energy per unit length and the inductance per unit length. The total stored energy and the inductance for a 15 m long dipole are also given.

Table 6.2.9: Stored Energy and Inductance at 6.5 kA as computed with the code POISSON for the SSC 50 mm aperture dipole.

2-d Modeling program compute stored energy. Don't forget to multiply by the program symmetry.

| Stored Energy per unit length, kJ/m | 105.0 |
| :--- | :--- |
| Stored Energy for 15 m long Dipole, kJ | 1575.6 |
| Inductance per unit length, $\mathrm{mH} / \mathrm{m}$ | 4.972 |
| Inductance for 15 m long Dipole, mH | 74.585 |

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## Lorentz Forces on the Individual Turn of the SSC Dipole

Figure 6.2.5: The magnitude of the components of the Lorentz force on the individual turns in a SSC 50 mm prototype magnet. The radial component of the force ( $F_{r}$ ) pushes the coil outward and the azimuthal component ( $F_{a}$ ) compresses the coil towards the midplane (horizontal plane). There are 19 turns in the inner layer and 26 turns in the outer layer of each quadrant.

Computed Lorentz forces at the design field of 6.6 T ( 6.6 kA ).



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## Forces in the Current Blocks of SSC Dipole

The table summarizes the forces on the blocks in the example SSC dipole.

| BLOCK | $F_{X}, \mathrm{lb} . / \mathrm{in}$ | $F_{Y}, \mathrm{lb} / \mathrm{in}$ | $F_{X}, \mathrm{~N} / \mathrm{m}$ | $\mathrm{F}_{\mathrm{Y},}, \mathrm{N} / \mathrm{m}$ |
| :---: | ---: | ---: | ---: | :--- |
| 1 | 1005.46 | -108.98 | $1.76 \mathrm{E}+05$ | $-1.91 \mathrm{E}+04$ |
| 2 | 1312.68 | -237.56 | $2.30 \mathrm{E}+05$ | $-4.16 \mathrm{E}+04$ |
| 3 | 612.52 | -151.97 | $1.07 \mathrm{E}+05$ | $-2.66 \mathrm{E}+04$ |
| 4 | 650.99 | -116.62 | $1.14 \mathrm{E}+05$ | $-2.04 \mathrm{E}+04$ |
| 5 | 231.73 | -384.60 | $4.06 \mathrm{E}+04$ | $-6.74 \mathrm{E}+04$ |
| 6 | 1208.86 | -1371.22 | $2.12 \mathrm{E}+05$ | $-2.40 \mathrm{E}+05$ |
| Total | 5022.24 | -2370.96 | $8.80 \mathrm{E}+05$ | $-4.15 \mathrm{E}+05$ |



Octupole in Quadrupoles When Quad Assembled Like Dipoles

Measured $b_{3}$ (octupole) $\neq 0$ )

also seen in Sc mm Arc Quad $(7$ unit in both magnets) Octupole From Quad

BASIC Problem
(a) Trying to assemble quad like dipole or
(b) Starting $Y_{0}$ ke i.d is circular (which becomes non-circular after assembly)
Fix

1. Stent out-of round yoke;
so that when assembled, it is round
$\rightarrow$ Too late for that (andtoo sypensive of changes)

- 2. Create deliberate asymmetry in coil
$\overrightarrow{A N D}$ Easy, Tested and worked $(b, y m$,


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## Goals of End Design

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## Magnetic Design

-Optimize for low integrated harmonics
-Guide design towards lower peak field without large increase in length
-Compute cross talk and fringe fields
Mechanical Layout
-Minimize strain and tilt of the cable in the end. Minimize large changes
-Cable and entire ends should be well supported (constrained)
In low field magnets, magnetic design drives the end design, whereas, in high field (high force) magnets, the mechanical design must!
These guiding principle are common to our all high force magnet designs (including 12 T common coil dipole design).

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## Ends of Cosine Theta Cable Magnets

## Superconducting

Magnet Division


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## Ends of Cosine Theta Cable Magnets

$\qquad$
DSA-207 RETURN END LEFT SIDE


 1 NO.CGO4RE $\mathbf{2}$ THELS.STARRETTCO. $\mathbf{3}$ ATHOLMASS.USSA. $\mathbf{4}$ TEMPERED $\mathbf{5}$


## Superconducting

## End Harmonic Optimization (conceptual)

## Magnet Division

Ends with spacer
(integrated harmonics \& peak field reduced)



Ends without spacer (large harmonics and peak field)

- End spacers increase the straight section length of some turns (turns at midplane go further out)
- Now consider the integral field generated by each turn. The harmonic component generated by a turn will depend on the angular location of it. The integral strength will depend on the length.
- A proper choice of end spacer can make integral end-harmonics small. However, note that the local values are large.
- Spacer also reduce the maximum value of field on the conductor (peak field) in the end.


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## Superconducting

## Block Structure

Magnet Division
Straight section (6 blocks, 70 turns): 302010433 (counting from midplane) 334102030 (counting from pole)
End section (8 blocks, 70 turns):
10584134620
(counting from pole)
Straight section => pole

$$
\begin{array}{ll}
3,3,4 & =>10 \\
4,10,20 & =>5,8,4,13 \\
30 & =>4,6,20
\end{array}
$$



Equal spacing in "Red Color" blocks is used as harmonic optimization parameters

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## Superconducting

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## Tilt of Turns in Various End Blocks (at far out position)



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## The AKF Parameter of Turns in Various End Blocks

AKF indicates the deviation from constant perimeter (hence strain on the cable)
Large Deviation from 1.0 is bad


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## Superconducting

## Coil End: Design A

Magnet Division


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## Peak Field Minimization

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A high peak field reduces the magnet quench performance.
> A large effort must be undertaken in 2-d optimization.


Usually about thousand cases are examined to :
-Minimize harmonics -Find a solution with lower peak field
-Good mechanical turn configuration (wedges, tilt angle, etc).

A series of computer programs have been written to carry out the above optimization in an exhaustive and systematic manner.

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## Peak Field in the Body of the Magnet

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## Peak Field in the End

## An example of an End Design



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## Peak Field in the End



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## Peak Field in the End How does it compare to Body?

In this example, the peak value is larger in the end than in the body of the magnet.

Component: BMOD
70.0

## Peak Field in the Ends

In cosine theta magnets, the conductors in the Ends are more strained and the mechanical design is generally less robust.

Therefore, one would like the peak field in the Ends to be less than that in the body of the magnet, to give conductor a larger margin.

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## Ends with 3 Re-adjusted Spacers



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## Superconducting

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## Peak Field Straight Section vs. Ends



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## Superconducting

## Block Structure

Magnet Division
Straight section (6 blocks, 70 turns): 302010433 (counting from midplane) 334102030 (counting from pole)
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Straight section => pole

$$
\begin{array}{ll}
3,3,4 & =>10 \\
4,10,20 & =>5,8,4,13 \\
30 & =>4,6,20
\end{array}
$$

## Must avoid large Ultum spacers

 (subdivide, if necessary)

Equal spacing in "Red Color" blocks is used as harmonic optimization parameters

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## Block configuration:

(8 blocks, 70 turns):
$10,5,8,4,13,4,6,20$
(counting from pole)

## End Harmonic Optimization: SMINSQ

End spacers in block \#2 (with 5 turns) and end spacer in block \#7 (with 4 turns).
All spacers with in a block have the same size.

Changing the size of two group of end spacers was adequate to get all harmonics small.
Computed values:

$$
\begin{aligned}
& B_{5}<1 \text { unit-meter; } \\
& B_{9} \text { and } B_{13}<0.1 \text { unit-m }
\end{aligned}
$$

Effective Magnetic Length $\sim 15.6 \mathrm{~cm}$
Mechanical Length $\sim 28 \mathrm{~cm}+$ End Saddle

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## Field through the Coil Ends

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## Field Harmonics through the End $\mathrm{b}_{5}$ : dodecapole

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## Field Harmonics through the End : $b_{9}$



Slide No. 48 of Lecture IV
Ramesh Gupta, BNL

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## Superconducting

## Field Harmonics through the End : $b_{13}$




## Common Coil Design (The Basic Concept)

- Simple 2-d geometry with large bend radius (no complex 3-d ends)
- Conductor friendly (suitable for brittle materials - most are - $\mathbf{N b}_{3} \mathbf{S n}$, HTS tapes and HTS cables)
- Compact (compared to single aperture LBL's D20 magnet, half the yoke size for two apertures)
- Block design (for large Lorentz forces at high fields)
- Efficient and methodical R\&D due to simple \& modular design
- Minimum requirements on big expensive tooling and labor
- Lower cost magnets expected

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## Superconducting

 Magnet DivisionUp-down asymmetry gives large skew harmonics, if done nothing. Integrate By.dl 10 mm above and 10 mm below midplane.

## ${ }^{\text {onen }}$ An up-down asymmetry in the ends with "no spacer"


$B_{y} 10 \mathrm{~mm}$ above and below midplane on magnet axis


Up-down asymmetry can be compensated with end spacers. One spacer is used below to match integral By.dl 10 mm above \& below midplane.


## $B_{y} 10 \mathrm{~mm}$ above and below midplane on magnet axis



A large Bz.dl in two ends ( $\sim 1$ T.m in 15 T magnet).

- Is it a problem?
- Examine AP issues.
- Zero integral.
- Lead end of one magnet + Return of the next magnet will make it cancel in about $\sim 1$ meter (cell length $\sim 200$ meters).
- Small v X B.


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## Superconducting

 Magnet Division$\qquad$

## An Example of End Optimization with ROXIE (iron not included)

End harmonics can be made small in a common coil design.

Contribution to integral $\left(a_{n}, b_{n}\right)$ in a 14 m long dipole $\left(<10^{-6}\right)$

End harmonics in Unit-m


| $\mathbf{n}$ | bn | an |
| :---: | :---: | :---: |
| $\mathbf{2}$ | 0.000 | 0.001 |
| $\mathbf{3}$ | 0.002 | 0.000 |
| $\mathbf{4}$ | 0.000 | -0.005 |
| $\mathbf{5}$ | 0.019 | 0.000 |
| $\mathbf{6}$ | 0.000 | -0.014 |
| $\mathbf{7}$ | 0.025 | 0.000 |
| $\mathbf{8}$ | 0.000 | -0.008 |
| $\mathbf{9}$ | -0.001 | 0.000 |
| $\mathbf{1 0}$ | 0.000 | -0.001 |
| $\mathbf{1 1}$ | -0.001 | 0.000 |
| $\mathbf{1 2}$ | 0.000 | 0.000 |



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## Superconducting

## Reduction in Peak Field in the Ends of Common Coil



Field and Field lines as computed by OPERA 2-d
15.1


3-d Model
$L$ L-300.0


## Peak field in the ends is minimized by:

- Removing iron over the ends
- Using end spacers between the turns.


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## Superconducting

 Magnet Division
## A Helical Magnet for AGS at BNL (1)

This magnet uses helical coils to maintain the polarization of the beam as it passes spin resonances in AGS.

Note: Particle Tracking


V VECTOR FIELDS

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## Superconducting

 Magnet Division
## A Helical Magnet for AGS at BNL (2)



$z(\mathrm{~mm})$
V- VECTOR FIELDS




## SUMMARY

## We are now expert in:

2d coil design
Requires good mechanical, magnetic and flexible design

3d coil design
Requires good mechanical and magnetic design

