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STABILITY OF A Nb₃Sn LOW-BETA QUADRUPOLE IN THE LHC RADIATION ENVIRONMENT

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Abstract

Use of Nb₃Sn coils can significantly improve the performances of the *low* β quadrupoles of hadron colliders. In this paper we report the results of a study to evaluate the average and the peak power deposition on the Nb₃Sn coils of a quadrupole for the second generation of the *low* β insertion for the LHC. The power release into the coils by radiation escaping from the interaction point has been investigated, by means of FLUKA code, as a function of different gradient–aperture combinations (the basic values being 300 T/m – 70 mm). Consequently the superconducting stability of the impregnated coils has been evaluated both at 2 and 4.2 K operation temperatures of the magnet.

1 INTRODUCTION

In the frame of the CERN-INFN collaboration on superconducting magnets for the LHC, a study of a very high performance quadrupole for a second generation of inner triplet of the *low* β insertions is under way at LASA lab of INFN-Milan. This type of quadrupole needs to reach very high magnetic field gradient, beyond 235 T/m in the machine, in a single aperture bore of 70 mm [1].

CERN has designed and built, in collaboration with Oxford Instruments (GB) [2], a full section 1.3 m long model of a novel design quadrupole, that has been successfully tested up 93% of the I_{max} reaching 245 T/m [4]. This quadrupole model makes use of very high performance NbTi cable with all kapton insulation.

To break the barrier of 250 T/m in 70 mm coil aperture we have designed a quadrupole, based on Nb₃Sn conductor, whose schematic cross section is reported in fig. 1. Nb₃Sn has very high performance in term of I_c characteristic vs field but suffers of serious drawbacks, namely the necessity of high temperature reaction (650 °C or more) and a brittleness that makes coil handling hazardous.

Brittleness of the Nb₃Sn and its I_c degradation vs transverse stress in case of bare cable requires the coils be fully impregnated under vacuum after reaction. Full impregnation of the coils raises a concern about coil stability due to lack of direct cooling.

The main points on stability in our design are:

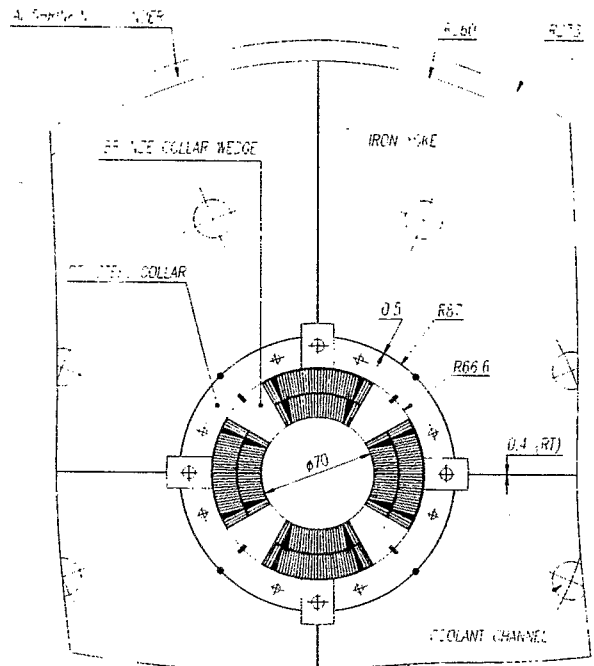


Figure 1: Schematic cross section of the Nb₃Sn quadrupole.

- the peak field on the coil, 11.5 T at 300 T/m, imply large operating stress in the coil, over 100 MPa in the high field region;
- impregnation of the coil will help mechanical stability on the coil head; this point is sometime neglected but it has been proved very important in the main dipole magnet [5], such that now is part of the design of the LHC dipoles;
- when working at 1.8 K, the disadvantage of specific heat, lower than at 4.2 K, is not balanced by the advantage of better heat transfer;
- the stability margin can be compromised by a considerable heat deposition into the coil due to radiation which can increase significantly the operating temperature.

A study of the power deposition has been carried out by CERN [6] with a simplified model and the power deposi-

Table 1: main features of quadrupole design

coil aperture (mm)	85	70
gradient (T/m)	250	300
operating temperature (K)	1.8	
superconductor	Nb ₃ Sn Int.Tin Diffusion	
J _c non Cu (A/mm ²)	1500 at 12 T, 4.2 K	
α=Cu:non Cu	1:1	
cable composition	36 strands, φ=0.825 mm	
cable size (mm)	1.34–1.60×15.0	
operating current (kA)	15.1	17.9
J _{overall} A/mm ²)	565	670
peak field on coils (T)	12.1	11.45
insulation type	R-glass+epoxy	
insulation thickness (mm)	0.125 azim., 0.250 rad.	
midplane shim (mm)	2 × 0.3mm	
temperature margin (K)	3.9	2.8
hot spot temperature (K)	≤ 120	≤ 130

tion has been found to 6 W/m (30 W per magnet). This point must be very carefully evaluated for coils which are almost adiabatic, especially because the power deposition into the coils – a substantial fraction of the total power – is strongly peaked at the coil midplanes.

2 QUADRUPOLE CHARACTERISTICS

We examined two designs with the same conductor: one that should generate 250 T/m in a 85 mm coil aperture [7] and the latest that aims to generate 300 T/m with 70 mm useful bore. The design stability margin is given by operation at 93–94 % of the I_{max} as given by short sample measurements.

The main parameters are reported in table 1.

The load line for the 300 T/m – 70 mm case together with the $J_{c,ov}$ at 4.2 and 1.8 K for the Nb₃Sn winding are reported in fig. 2, where for comparison is reported the $J_{c,ov}$ characteristic of the NbTi CERN-Oxford quadrupole.

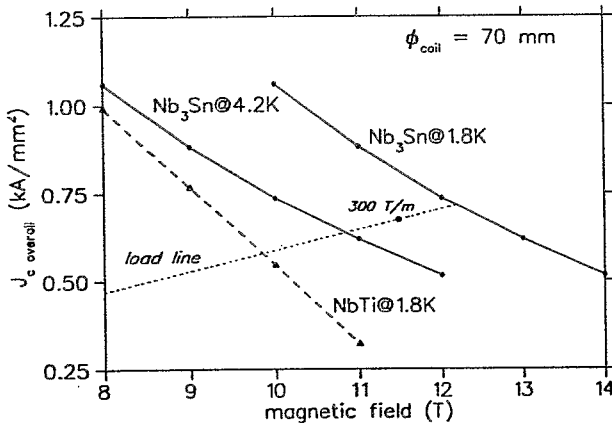


Figure 2: Coil J_c for our Nb₃Sn and for NbTi and magnet load line.

3 POWER DEPOSITED BY RADIATION

The source of the radiation power escaping from the interaction point (I.P.) has been evaluated by taking few hundreds of the 7+7 TeV events from DTUJET code. The interesting events are the inelastic and the single diffractive ones, whose cross sections are $\sigma_{in} = 60$ mb and $\sigma_{sd} = 12$ mb respectively. The neutral particles released in the interactions were directly processed by FLUKA [9], while the charged ones were transported along the beam pipe through the detector field and through the inner quadrupole triplet of the *low β* insertion. Low energy neutrons have been neglected. Whenever a charged particle hits the beam line structure it was processed by FLUKA as a single event.

The inner triplet starts with Q_1 placed at 23 m from I.P., with a free space of 2.5 m between Q_1 and Q_2 , 1 m between Q_2 and Q_3 , 2.3 m between Q_3 and Q_4 . Between the I.P. and Q_1 a $\phi_{in} = 30$ mm, 1.8 m long, copper collimator is placed at 19 m from I.P. Moreover the effect of four stainless steel absorbers inside the beam pipe along the quadrupoles has been evaluated. The thickness of the absorbers is 5 mm [8].

The results of the calculation for the nominal luminosity of 10^{34} cm⁻² s⁻¹ are summarized in table 2 where data refers to Q_2 that is by far the hottest quadrupole.

The peak power deposited into the coils ranges from 1.3 (for 250 T/m – 85 mm) up to 2.1 mW/cm³ (for 300 T/m – 70 mm). We think that the grid used is too wide (about $1 \times 1 \times 50$ cm³) and these should be regarded as minimum values. For these reasons and to be conservative we preferred, in thermal analysis, to concentrate the whole power released in the coils near the midplane. In this way, in case of 300 T/m – 70 mm, the power density along the midplane for the hottest coil becomes 5.8 mW/cm³.

The results so far obtained are very preliminary but sufficient to evaluate the stability of the Nb₃Sn coil. More accurate calculation, with bigger statistic and higher number of events is underway.

4 2-D THERMAL MODEL

A 2-D thermal analysis was carried out by means of ANSYS code on an octant. Main features of the thermal model are:

1. steady state;
2. linear analysis, i.e. material properties are constant at the initial temperature. This hypothesis is confirmed a posteriori by the very small temperature increase;
3. regions are (iron is neglected since is thermally separated):

- (a) two coil shells, each one containing a longitudinal copper wedge, where the material properties are the effective values of the conductor unit cell; typical mesh dimension in the coil is 2×2 mm²;

Table 2: heat deposition (four poles) and coil temperature rise above 2 K

apert. \times grad. (mm \times T/m)	power into coils (W/m)	hottest coil (W/m)	power in th. calc. (mW/cm ³)	ΔT_{max} coils (K)	total power into cold mass (W/m)	power into adsorber (W/m)
70 \times 250	1.27	0.36	4.7	0.42	3.84	2.31
70 \times 300	1.44	0.44	5.8	0.50	4.30	2.56
85 \times 250	1.27	0.35	4.6	0.42	3.98	2.42

Table 3: stability margin of the coils (without radiation)

apert. \times grad.	T_{batt} (K)	ΔT_m (K)	ΔH_m (kJ/m ³)
70 \times 300	2.0	2.8	5.24
70 \times 300	4.2	2.4	11.3

- (b) radial insulation regions on the inner coil radius, between two coil shells and between outer coil radius and collar;
 - (c) pole wedge in bronze and, separated, the insulation between pole wedge and coils;
 - (d) stainless steel ring collar;
4. no heat exchange at midplane and along the 45° boundaries (symmetry conditions); convective heat exchange with cryogen, HeII or LHe, at fixed temperature: a conservative heat transfer coefficient of 100 W/m²-K has been taken.
 5. the total power released into the coil has been concentrated in the mesh elements along the midplane; computation has been done for the coil with the biggest heat deposition;

The inner bound insulation region has been magnified in radial dimension to be properly described and consequently its thermal conductivity amplified by the same factor.

The results of this preliminary analysis, computed at 2 K where the thermal conductivity are worse than in the case of 4.2 K, are summarized in the table 2.

The heat flow to the helium is very low, with a peak of about 17.7 W/m², generating a temperature rise of 0.17 K at the coil to helium interface.

5 STABILITY ANALYSIS

The numbers reported in the previous table must be compared with the stability margin for the coil operation. The stability can be reported as temperature margin, ΔT_m , or enthalpy margin (against distributed continuous perturbation), ΔH_m . These values are reported in table 3 without taking into account the radiation heat.

It can be seen on the temperature margin that the 0.5 K of temperature increase due to heat deposition, see table 2, has a small effect on the coil stability. Actually the first 0.5 K above 2 K gives negligible contribution to enthalpy margin which decreases from 5.24 down to 4.8 kJ/m³. The situation is even better at 4.2 K where the residual 1.9 K

temperature margin corresponds to a more comfortable 9.8 kJ/m³ enthalpy margin.

6 CONCLUSIONS

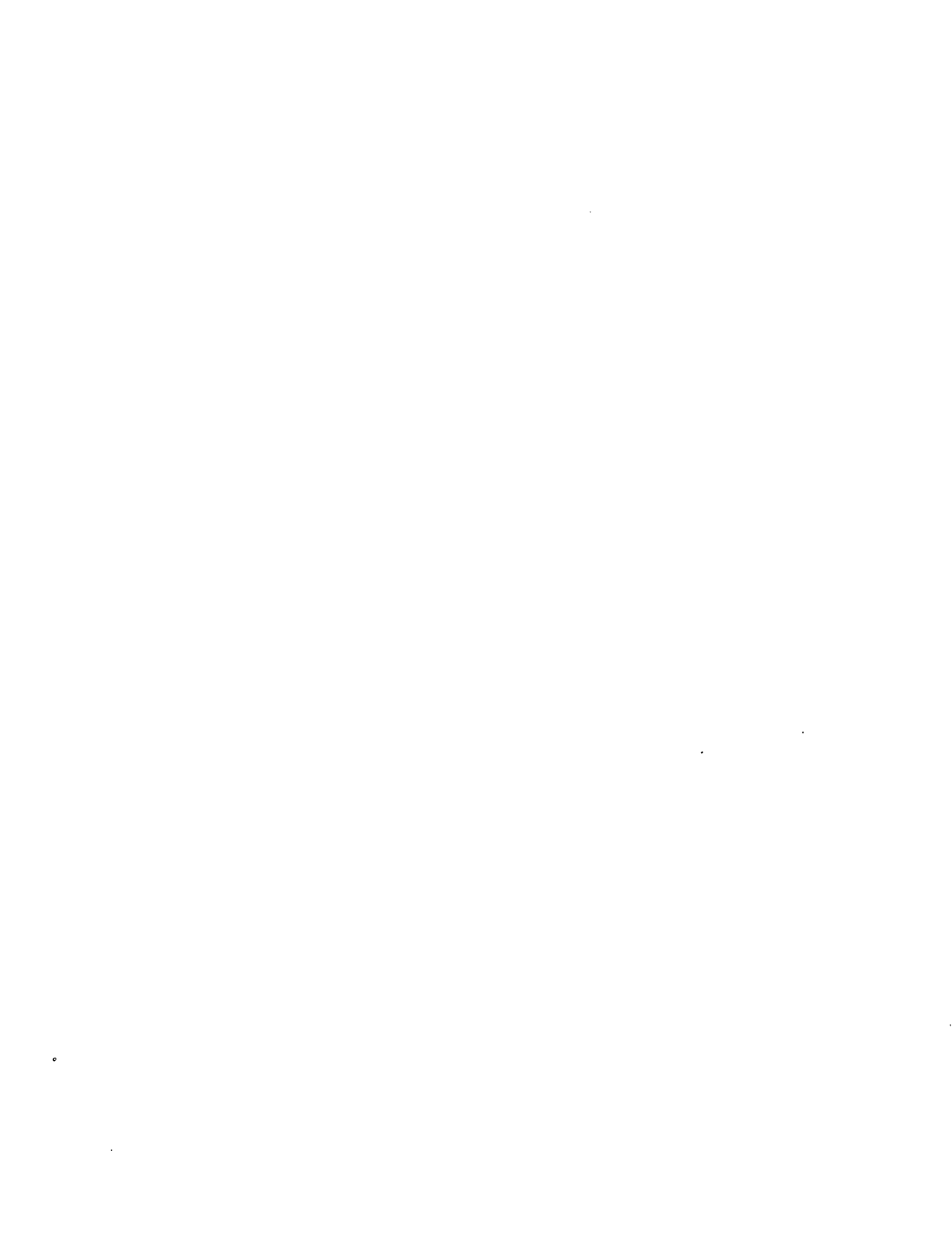
This preliminary calculation shows that Nb₃Sn technology can be employed for these very demanding magnets. From the point of view of stability the heat deposition by radiation is acceptable without endangering the low β quadrupole operation also at superfluid temperature.

7 ACKNOWLEDGEMENTS

Authors gratefully acknowledge G. Kirby, R. Ostojic and J. Strait of CERN for the useful discussion and information provided.

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A 300 T/m Nb₃Sn QUADRUPOLE FOR THE LOW-BETA INSERTIONS OF THE LHC

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Abstract

The NbTi quadrupoles for the LHC insertions have a design gradient of 250 T/m at 1.8 K with a 70 mm aperture. We are exploring the use of the Nb₃Sn technology for second generation *low β* quadrupoles for LHC. A conceptual design of a Nb₃Sn quadrupole operating at 1.8 K with 300 T/m gradient and 70 mm aperture is presented. Cable performance, winding technique, magnetic design, mechanical structure and magnet protection are discussed.

1 INTRODUCTION

In the LHC (Large Hadron Collider) the proton beams are focussed at each side of the IP (Interaction Point) by a set of four quadrupoles, series connected, forming the inner triplet. These quadrupoles, 5.5 m long, have single aperture (70 mm) and a design gradient of 250 T/m at 1.8 K.

A 1.3 m long model using an high quality NbTi conductor and a novel design [1] has been built by a CERN-Oxford Instruments (GB) collaboration and has reached a gradient of 245 T/m.

In the frame of a collaboration between CERN and INFN on superconducting magnets we have looked to alternative design using Nb₃Sn conductor. An initial investigation [2], based on simple scaling laws and assuming as a reference the CERN-Oxford design, has explored various layout (shell or rectangular block) and different bore apertures (70 or 85 mm) indicating the possibility to reach with the Nb₃Sn technology the 300 T/m with the small aperture or 250 T/m with the large bore.

A more detailed study [3], including the mechanical structure and magnet protection, has been completed for a large aperture (85 mm) quadrupole. Finally the choice has been made to maintain the nominal aperture (70 mm) and to aim at gradient of 300 T/m. In this paper the conceptual design of this quadrupole is presented and the relevant aspects are discussed in detail.

2 CONDUCTOR

We plan to use a conductor manufactured by Europa Metall (Florence) with the Internal Tin Diffusion method. It has 37 islands embedded in a copper matrix. Each island, which has a Sn core and 190 Nb filaments, is enclosed by a Nb barrier, whose thickness is about 20 μm, to prevent poisoning of the pure copper during heat treatment. The barrier is partially reacted into Nb₃Sn : this helps to obtain

Table 1: Superconducting cable.

superconductor	Nb ₃ Sn - Int. Tin Diffusion
barrier	niobium
effective ϕ_{fil}	20-30 μm
ϕ strand	0.825 mm
$\alpha = \text{Cu}:\text{non Cu}$	1
heat treat. temp.	up to 650°C, 350 hours
J_c non Cu	1500 A/mm ² @ 12 T, 4.2 K
n. of strands	36
cable size	1.34 ÷ 1.60 × 15.0 mm ²

very high I_c values but is detrimental for the effective filament diameter that ranges from 20 to 30 μm, depending on the layout [4].

The Nb₃Sn critical current density measured in virgin wires at 4.2 K vs applied magnetic field is reported in Fig 1 together with the scaling to 1.8 K. Also shown, for comparison, is the NbTi curve of the conductor used in the CERN-Oxford quadrupole. The cable characteristics are reported in Table 1.

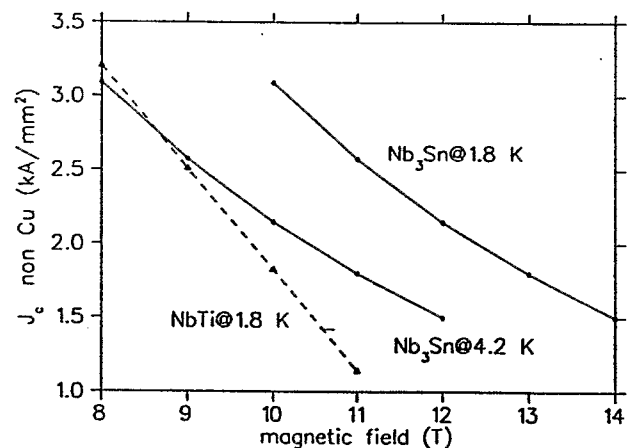


Figure 1: Non copper critical current density for Nb₃Sn and NbTi in virgin wires.

Disadvantages of Nb₃Sn technology with respect to the more standard NbTi one have been extensively described in [3]. In the meanwhile we have investigated I_c degradation due to cabling. In a prototype cable, 150 m long and with 7% lower current density, degradation after cabling (measured on wires) was found about 5%, very near to the stan-

ard 3% for good NbTi cable. Also the I_c degradation vs transverse compression, a key factor for this type of magnet, will be soon measured.

To avoid degradation of the outer copper quality and possible flake off of the barrier during heat treatment due to oxygen contamination through the thin copper layer, our conductor must be treated under vacuum or inert atmosphere during all the stages of reaction.

3 INSULATION AND WINDING

The coils are wound and then reacted. So far we have tested for insulation three different materials: R-glass tape, mica-glass tape from Isovolta (A) and a ceramic-glass composite foil from Toshiba (J). Each insulation has been treated at high temperature (700 °C) in Ar atmosphere and under vacuum in shape of tape or foil and then the insulation resistance has been measured up to 1 kV under a pressure of over 50 MPa.

Results have been very positive for the mica-glass tape and for the ceramic-glass foil, with $R \geq 10 \text{ M}\Omega$ at 1 kV and $\geq 1 \text{ G}\Omega$ at 500 V. The R-glass tape has shown, after a special chemical treatment [5] to remove the sizing agent, insulation resistance of typically 1 M Ω at 1 kV and $\geq 1 \text{ G}\Omega$ at 500 V. It should be mentioned that resistance on glass samples occasionally is much lower and that the chemical process is not yet tested on a real coil. We rely also on the fact that continuity of the insulation should be assured by proper impregnation of the coil under vacuum. For the time being we keep as interturn insulation a R-glass tape, that is easy available in different dimensions, butt-lap wrapped around the cable and enclosing, on one side only of the cable, a further glass tape. In this way the turn to turn gap is 240 μm .

The alternatives under study are:

i) use of glass tape wrapped around the cable and enclosing a mica-glass tape (probably it is the best solution for insulation),

ii) ceramic-glass tape used only as a spacer, if it will be available in the proper dimension. In this way the turn to turn distance is 150 μm .

The ceramic-glass tape is more robust than R-glass and more flexible than mica-glass but unfortunately the minimum available thickness is 150 μm while R-glass is 80 μm thick and mica-glass is 100 μm .

The coils are wound in double pancake technique to avoid the joint between the inner and outer layer of each pole which is normally a weak point of the magnet and can limit the performance.

4 COIL LAYOUT

The double pancake technique prevents the possibility of grading; we point out however that the grading is less effective in Nb₃Sn coils respect to NbTi because of the reduced slope of the J_c vs B line (see fig 1). No specific cable optimization has been pursued; we selected actually

Table 2: Quadrupole characteristics.

Operating gradient	(T/m)	300
Load line margin		94%
Peak field in conductor	(T)	11.5
Operating current	(kA)	17.9
$ b_6 , b_{10} $	(10^{-4} units)	< 0.01
No. of turns per coil		12+17
Inner/outer coil diameter	(mm)	70/132.2
Inner/outer yoke diameter	(mm)	174/520
Stored energy	(kJ/m)	453
Inductance	(mH/m)	2.8

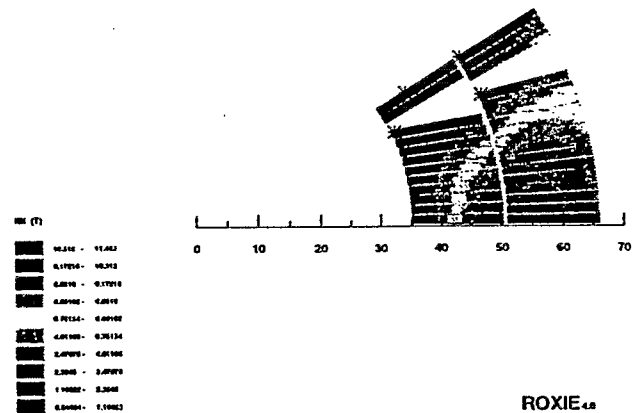


Figure 2: Coil layout (one octant) with field distribution. Peak field locations for each block are indicated.

strand diameter, cable size and keystone angle corresponding to the outer layer of the LHC dipole in order to avoid any cabling development.

The cable layout, calculated with ROXIE [6], is presented in fig 2 and the main parameters of the magnet design are given in Table 2. The 300 T/m gradient corresponds to 94% of the load line; so far no degradation for the azimuthal stress (110 MPa) has been considered pending measurements. The iron yoke is placed at 20 mm from the coil in order to have a rigid collars; an iron closer to the coils, 10 mm, gives only 1.4% increase in gradient. The operating current is large, 17.9 kA, but still manageable, to get the advantage of a low inductance which makes the protection of the magnet easier.

5 MECHANICAL STRUCTURE

The main characteristics of the mechanical structure, schematically shown in fig 3 are:

- i) a 20 mm stainless steel collar ring and four bronze wedges in the pole regions;
 - ii) an iron yoke, 174 mm thick, cut in four parts and in contact with the collar only near the coil midplanes;
 - iii) an aluminum shrinking cylinder 13 mm thick.
- The mechanical structure has been designed to mini-

mize the transverse compression in order to reduce the critical current degradation. The mechanical separation of the pole wedge from the ring collars allows to keep the wedge always in position from winding to yoking. The collar wedges are made out of bronze to fit the thermal contraction of the coil during the cooldown. Materials with a lower thermal contraction, like Ti alloys, give a further prestress during cooldown, and also longitudinal stress on coil heads, but they cause an elliptical deformation of coils increasing the peak stress during magnet excitation. Material and thickness of the collar ring have been designed to withstand the magnetic forces with negligible deformation.

Almost all coil prestress is given during collaring and the force needed is about 500 ton/m using a dipole type press. Two collaring techniques are under study: the first using a special press[3] acting simultaneously on X,Y planes, the second using clamps for the precollaring and then a dipole type press acting alternatively on the two planes and inserting pinning rods of increasing diameter. This last technique requires an expandable collaring mandrel.

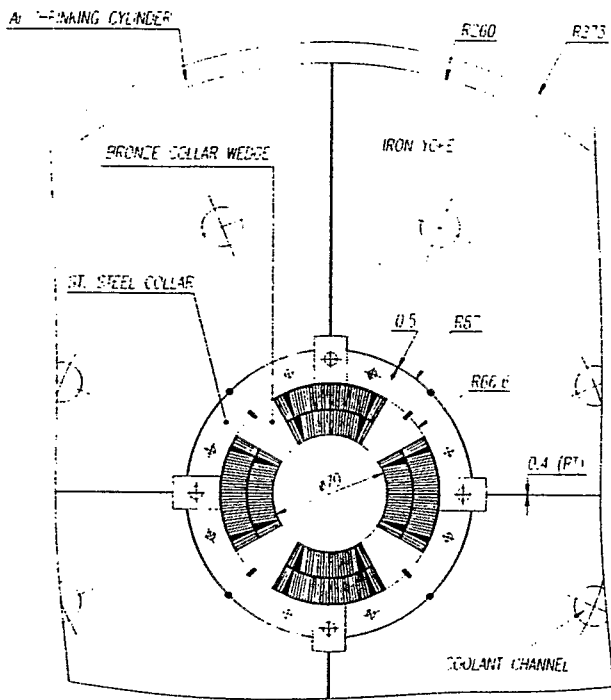


Figure 3: Cross section of the magnet

The yoke and the shrinking cylinder have been designed in order to minimize the coil stress and the coil deformation at full magnet excitation. In the first studies, trying different collar materials and thicknesses, the peak pressure in the coils was always found in the midplane at the inner side of the first layer (as usual in this type of magnet) with values between 140 and 160 MPa. This peak was due to the e.m. forces ($F_r = 93$ ton/m, $F_\theta = -289$ ton/m in each octant) and to a deformation of the collar that give a 0.06 mm deformation of the inner edges of the coils.

In this design the shrinking cylinder gives after cooldown

a force of 100 ton/m for octant that the yoke transfers to the collar only on the midplanes, causing an opposite deformation. In this way when the magnet is at maximum excitation this deformation is almost all compensated, the peak stress value (123 MPa) is still on the midplane but on the outer edge of the coil (where the field is much lower) while on the inner edge the peak stress is less than 110 MPa.

6 STABILITY AND PROTECTION

A preliminary evaluation [7] of the stability problem of the coils at 1.8 K with a peak power deposition by radiation of 2.1 mW/cm^3 indicates that impregnated coils will have a sound margin of stability, 1.9 K above the operating temperature.

The magnet can be protected with an active heaters system, as in the LHC dipoles. A passive system based on a dumping resistor can be used for a 1 m long model[3]. The extension of the passive system to a full length magnet is under investigation.

7 CONCLUSION

A 70 mm bore Nb_3Sn quadrupole with a gradient of 300 T/m at 1.8 K seems feasible with the present state of the art and without major cable improving.

A proposal will be submitted to INFN in 1996 to build and test a 1 m long model.

Test on insulation and critical current degradation under stress are under way. A small racetrack coil (12 turns, 0.4 m long) is under construction to verify all technical aspects of insulation, winding and impregnation.

8 ACKNOWLEDGEMENTS

We wish to thank T. Taylor, R. Ostojic and S. Russenschuck of CERN for fruitful discussions.

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THE 1m LONG SINGLE APERTURE DIPOLE COIL TEST PROGRAM FOR LHC

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Abstract

The CERN short model activity on main dipole magnets is centred around the design, in-house fabrication and testing of single and twin aperture 1m long magnets. In order to study the influence of individual coil parameters on the magnet behaviour with a fast turn around rate and to qualify the possible design solutions priority was given to the fabrication of a certain number of single aperture dipole models. The collared coils are assembled in a reusable yoke structure and tested in a vertical cryostat at 2 K. The present paper reviews the aims of the program, the design and fabrication to date of single aperture models, their instrumentation and the preliminary results and conclusions.

1 INTRODUCTION

The Large Hadron Collider (LHC) [1], approved by the CERN Council in December 1994, is a 7 TeV proton accelerator-collider operating at a nominal field of 8.3 T. Its main components are double aperture superconducting dipole magnets which have evolved for reasons of economy and machine requirements from 10 to 15 m in length, 50 to 56 mm in bore diameter, and 17 to 15 mm in cable width [2]. Of the former designs seven 10 m long double aperture dipoles and 14 short 1 m models (both single and double aperture with constructional variants) were built and tested, totalling nearly 70 test campaigns at cryogenic temperatures [3]. Results show that these magnets have similar performance and training behaviour, exhibiting comparable weak spots located especially in the coil ends, jump/splice regions, transitions of cross section and often in the innermost coil turn. Since improvements and design options can conveniently be studied with short magnets an intensive fabrication program of 1 m models of the new design started at CERN to provide the required input to the long magnet program presently in progress in industry.

2 AIMS OF THE MODEL PROGRAM

Priority was given to the fabrication, collaring and testing of model coils in a single aperture structure (MBSMS) to optimise specific components and procedures, implement new design features, check expected performances and provide a facility for testing cable performance and possible variants. Present rate of model

completion is about one per month with a lead time of 3 to 4 months. This will allow in due time, to incorporate in the design lessons learned from the cold test results, to refine the assembly techniques and to accumulate statistics. Double aperture models will be built using the same collars and yoke laminations as for the long magnets. The first such model is planned for measurements in early autumn of this year.

3 DESIGN AND FABRICATION

The cross-section and main parameters of MBSMS models are shown in Fig. 1 and Table 1 below.

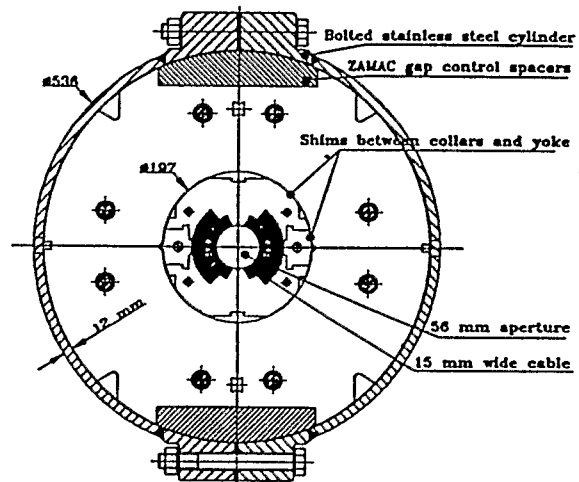


Fig. 1: Cross section of MBSMS type model.

Coil inner diameter	56 mm
Quenching field	9.6T @ 1.9K and 13240 A
Nominal current (Inom) @ 8.368T	11460 A
Ratio of peak field to central field	1.05
Overall coil length	1080 mm
Length of magnetic steel in the yoke	560 mm
Magnetic length	862 mm
Total inductance	3.2 mH
Magnetic forces per quadrant @ Inom	$\Sigma F_x = 1650$ N/mm $\Sigma F_y = -820$ N/mm
Total axial force @ Inom	19 tons
Cu/SC ratio of inner / outer strands	1.60 / 1.90
Ic / dIc/dB of inner cable @ 1.9 K, 10 T	≥ 13.75 kA / 4.8 kA/T
Ic / dIc/dB of outer cable @ 1.9 K, 9 T	≥ 12.95 kA / 3.65 kA/T

3.1 Magnetic design

The coils have two layers, made of graded cables with 28 strands of 1.065 mm diameter for the inner layer and with 36 strands of 0.825 mm for the outer layer. The conductors are distributed in five blocks giving the required low content of b7 and b9 components. The load line and expected quench fields calculated with the programs ROXIE and POISOPT are shown in Fig. 2. Three basic types of coil end geometries were designed (see Table 2). Around both coil ends the yoke is made of non-magnetic laminations to reduce the peak field in these critical regions.

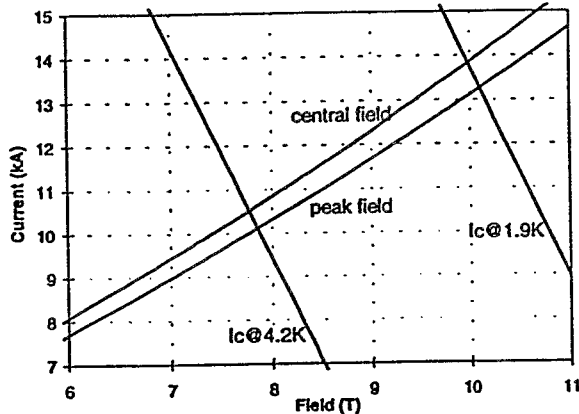


Figure 2. Load line for MBSMS magnets.

3.2 Mechanical design

The MBSMS design concept [4] is to allow reproducible coil testing conditions. The collars, made of Al-alloy type 5083, are of similar mechanical rigidity as adopted for the double aperture design and locked by stainless steel rods. The yoke is vertically split with an open gap at room temperature, controlled by spacers made of Zn/Al alloy, and held together by an outer stainless steel bolted shrinking cylinder which has an azimuthal pre-stress of 150 MPa at ambient temperature. This design permits an easy assembly and reuse of the yoke structure. The gap control spacers limit coil compression during assembly but shrink away at cold so that under nominal conditions the gap is closed and the collars are just in contact with the yoke laminations. The gap remains closed up to 9.7 T. The design azimuthal coil pre-stress after assembly at room temperature is 50 MPa, both for inner and outer layers. Since longitudinal compression of the coil heads may improve their training performance, so-called end-cages consisting of a glued collar pack, a flange at the coil end and four tie rods, allow to pre-load the coil heads up to 8 tons.

3.3 Instrumentation and Protection

The following instrumentation, needed to monitor assembly and testing is implemented in all models:

- Two spot heaters, to trigger transitions for quench studies, are placed one in the inner and one in the outer layer between the cable and the innermost end spacer, made of 50 μm thick stainless steel foil glued between two 25 μm thick polyimide foils.
- Voltage taps, in total 50 per model, for quench detection and location.
- Special collars with strain gauges near the poles to monitor the coil pre-stress in the inner and outer layers during magnet assembly, cool down and power testing. New capacitive pressure transducers were developed and already used at ambient and LN2 temperatures for collaring tests and coil modulus tests.
- End cage tie rods with strain gauges.
- Bullet gauges to monitor the force exerted by the coils against the thick end plates of the magnet.
- Gap opening transducers to monitor the status of the vertical yoke gap.

The mechanical instrumentation is calibrated both at ambient temperature and at 1.9 K.

All models have quench heater strips placed between the outer layer and the ground insulation, to rapidly heat up the conductors in case of a quench. The calculated hot spot temperature assuming full energy dissipation in the coils and triggering of the heaters is around 240 K.

3.4 Fabrication and description of variants

Coil winding of MBSMS1 started in June 1995 and since then seven models have been manufactured and some reworked (see Table 2). The cable insulation is all in polyimide and composed typically of two layers of 25 μm thick tapes each overlapped by 48%, and a third 70 μm thick adhesive-coated layer, spaced by 2 mm to provide channels for helium penetration inside the coils. After winding, each layer is heated in a mould to 185°C for 30 minutes gluing turns firmly together. The coil heads are then impregnated with a heavily charged resin and the layers assembled with a grooved G11 sheet placed in-between them. The inside cable ends are joined with an AgSn alloy, reinsulated and glued back onto the coil blocks in an operation referred to as reconditioning. The size and modulus of each layer are then measured to define pole and coil head shimming for collaring. Different collaring variants have been tried. Typically the coils are compressed to about 120-130 MPa, the collaring rods are inserted and the external pressure released. The residual coil pre-stress is about 50-60 MPa on both layers. The end cages have been tightened generally after collaring. For models 4 and 5 the coils have been stretched during collaring with an internal mandrel reacting against the innermost coil end spacers and the cages pre-tightened at an intermediate stage to better distribute the longitudinal stresses inside the coil heads. An example of such a collaring is shown in Fig. 3.

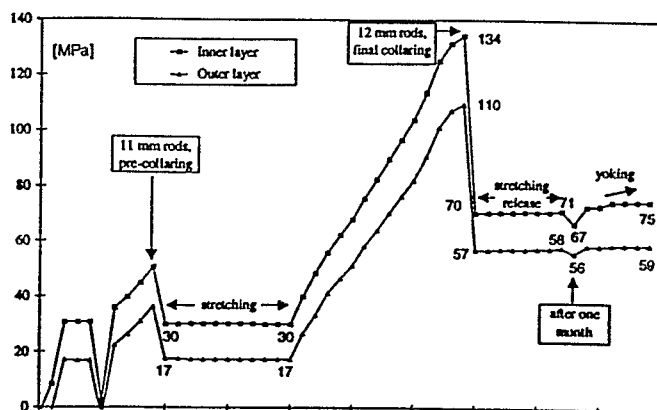


Figure 3. Collaring procedure for magnet MBSMS5

The main fabrication variants are shown in Table 2.

Table 2: fabrication variants from MBSMS1 to MBSMS8		
<i>magnet</i>	<i>coils*</i>	<i>assembly</i>
MBSMS1.V1	non coated cable type I end spacers	reconditioning @ 185°C end cage tightened
MBSMS2.V1	non coated cable type II end spacers	reconditioning @ 100°C end cage tightened
MBSMS2.V2		end cage untightened
MBSMS3.V1	non coated cable type II & IIa end spacers	reconditioning under pressure end cage tightened
MBSMS1.V2		reconditioning under pressure
MBSMS4.V1	tin coated strands cable type II end spacers	reconditioning under pressure end cage tightened coil stretching
MBSMS5.V1	tin coated strands cable type II end spacers no interturn spacers 150 mm layer jump	reconditioning under pressure end cage tightened collaring under stretching
MBSMS6.V1	tin coated strands cable type II end spacers no interturn spacers 150 mm layer jump	reconditioning under pressure end cage tightened different collar material
MBSMS7.V1	tin coated strands cable type III end spacers 150 mm layer jump coil end not impregnated	reconditioning under pressure end cage not present
MBSMS8.V1	tin coated strands cable type II end spacers no interturn spacers 150 mm layer jump different cable insulation	reconditioning under pressure end cage tightened

*end spacers types, design principle and material:
 type I : minimum deformation energy, G11; type II : isoperimetric, G11;
 type IIa : isoperimetric, PEI; type III : Fermilab design with "shoes", G11;

4 TEST RESULTS

Cold tests have been made so far up to model 4. Power test results are reported in Ref. [5]. All models are cooled directly to 1.9 K, first training quenches were at 8.2, 8.65, 8.87, and 8.67 T for models 1, 2, 3, and 4 respectively, followed by slow training, occurring mainly in the first turn and in the transition between straight part and heads of the inner layer. Quenching with full energy deposition gives a hot spot temperature of 240 K and shows a safe magnet protection scheme. After a few

such quenches, training switches to the first turn of the outer layer at a lowered unstable field level which remained however above 8.9 T for model 3. After a thermal cycle, this model retrained at 9.05 T, showing limited training memory, but no longer exhibited the unstable behaviour mentioned above and reached 9.5 T after 16 training quenches.

5 CONCLUSIONS

A full test and result analysis program is going on for these models. In parallel the fabrication of new models and relevant variants is underway. Results show that general improvements in the assembly have brought better training behaviour, model 3 reaching 92% of the calculated short sample limit on its first quench, proving the validity of the basic design. So far there are no clear indications of a significant advantage coming from a particular end spacer type or from a specific technique like end cage tightening or coil stretching. Known weak spots have been addressed in later models: e.g. increased length of layer jump region, suppression of interturn spacers in the coil ends, smoother transitions between straight part and coil ends. Important aspects of dipole performance like dynamic behaviour and field quality are also addressed by the model program and will steer specific design solutions and fabrication techniques.

6 ACKNOWLEDGEMENTS

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LOW TEMPERATURE QUENCH PERFORMANCE OF FERMILAB LOW- β INSERTION QUADRUPOLES

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Abstract

The recently revived superconducting magnet program at Fermilab is currently focused on the development of high gradient quadrupoles for possible use in the Large Hadron Collider (LHC) interaction regions at CERN. In order to provide input for the new quadrupole design which will operate in superfluid helium, we have tested a Fermilab Tevatron low- β quadrupole cold mass and compared its low temperature performance to a newly assembled heavily instrumented version which was mechanically modified to take advantage of the gain in critical current.

1 INTRODUCTION

Fermilab has successfully operated low- β quadrupoles for the Tevatron D0/B0 interaction regions at an operating gradient of 141 T/m at 4.5 K. Future low- β insertions, in particular, the one proposed for the LHC[1], require a 50% higher operating field gradient. Fermilab in cooperation with Brookhaven National Laboratory (BNL) and Lawrence Berkeley National Laboratory (LBNL), is designing a higher gradient quadrupole suitable for the LHC low- β insertions which uses NbTi conductor in superfluid helium[2]. A first step in this program has been to evaluate the feasibility of using a mechanically modified version of the existing Fermilab low- β quadrupole for superfluid operation. We report on results from tests of two Fermilab-style superconducting high gradient quadrupoles.

2 MAGNET DESCRIPTION

The magnets for this study are 1.4 m long Tevatron low- β quadrupoles. Details of the design have been described elsewhere[3,4]. This cold iron superconducting quadrupole has a 2-shell, cos 2θ coil with a 76 mm aperture and an outer cold mass diameter of 276 mm. The inner and outer coils are made from 36 strand Rutherford cable. The strands are 0.528 mm in diameter and contain 13 μ m filaments. There is a copper wedge in the inner coil whose primary purpose is to minimize the geometric 12- and 20-pole harmonics. Four inner to outer coil splices are located in the magnet lead end radially beyond the outer coil and are made through pre-formed solder-

filled cable originating from the lead end pole turn. The coils are supported in the body by aluminum collars. The splice and the coil lead and return ends are clamped with a 4 piece G-10 collet assembly enclosed in a tapered cylindrical can. Iron yoke laminations surround the coil in the body region, and stainless steel laminations surround the end region cylindrical can. A welded stainless steel skin surrounds the yoke.

The two magnets differ in mechanical support and in instrumentation. One magnet (LBQ5425) was built as a spare for the Tevatron and as such has the nominal construction features and instrumentation for a production magnet. There are voltage taps across each quadrant (inner-outer coil pair). The aforementioned cylindrical end cans are made of stainless steel.

A finite element analysis mechanical model of the nominal production magnet at 1.8 K and full current excitation predicted inadequate coil support for both the body and the end regions[5]. Thus a second magnet (R54001) was built with the same tooling but with enhanced mechanical support. To increase the magnet end prestress at low temperature, the stainless steel end cans were replaced with aluminum cans. Kapton pole shims were inserted to increase the coil azimuthal prestress. These shims also increase the effective collar diameter to assure that there was an interference fit between the iron and collared coil. R54001 had several voltage taps in the inner coil concentrated near the pole turn and the copper wedge. Strain gauge transducers in the collars monitored changes in coil stress during manufacturing, cool down and excitation.

3 MAGNET TESTS

LBQ5425 was tested at the Lawrence Berkeley National Laboratory Magnet Test Facility in a horizontally oriented liquid helium dewar. The facility is designed to support superfluid helium at 1.8 K and at 1 atmosphere. R54001 was tested at the Fermilab Technical Support Section horizontal magnet test facility[6]. The test stand was originally designed to test SSC dipole cryostated cold masses at 4 atmospheres and temperatures from 4.6 K to 1.8 K. The outer diameter of our low- β quadrupole is roughly the same as an SSC dipole, thus we were able to build a special shorter length SSC-style cryostat to accommodate this magnet.

3.1 Magnet Training

Both magnets were first trained at 4.3 K prior to superfluid testing. This allows one to compare the training of these magnets with previous low- β quadrupoles[7] and to observe the change in magnet training between normal and superfluid helium. The training histories for LBQ5425 and R54001 are presented in Figures 1 and 2.

At 4.3 K LBQ5425 achieved 4700A on the first quench and required 4 quenches to reach its plateau of 5150 A. The magnet was warmed to room temperature and then cooled to 1.8 K. The first quench was above 6000 A, significantly higher than the 4.3 K plateau quenches. After two more training quenches, the quench current fell to near 5000 A and did not increase in the next 3 quenches. Two quenches (not shown) at 3.8 K

yielded similar results. The magnet was then warmed to 4.3 K and quenched again at just above 5000 A. Finally, it was again cooled to 1.8 K and came within 100 amps of reaching its previous quench current maximum. The 5000 A quenches at 1.8 K, as well as the low quench at 4.3 K (#5 in Fig. 1), are likely due to insufficient coil support.

R54001, with enhanced mechanical structure, exhibited improved quench performance relative to LBQ5425, as seen in Fig. 2. At 4.3 K it came within 200 amps of its quench plateau in three quenches. The training quenches largely occurred within the pole turn return end. After 9 quenches, but before it reached its plateau, the magnet was cooled to 1.8 K. Here it exhibited significant training, but with monotonically increasing quench current. The training quenches occurred predominantly in the coil ends. After the 1.8 K testing the 4.3 K quench plateau was achieved. The plateau quenches occurred near the inner-outer coil splice.

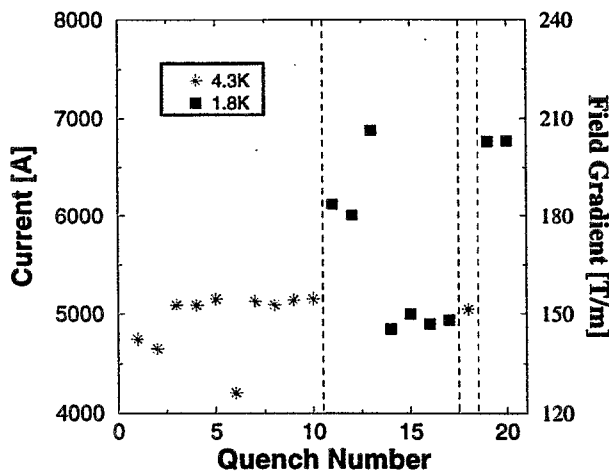


Figure 1: Quench training history for LBQ5425.

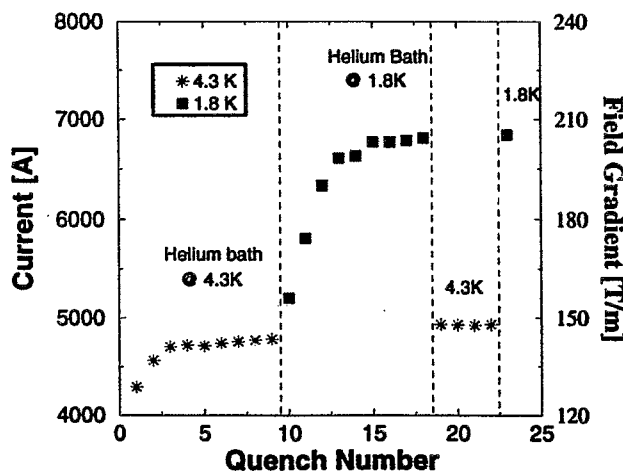


Figure 2: Quench training history for R54001.

3.2 Temperature Dependence of Quench Current

After 9 quenches at 1.8 K, R54001 was quenched at several temperatures between 1.8 K and 4.3 K as shown in Fig. 3. There was a monotonic decrease in quench current with increasing temperature as expected for a magnet which is not limited by mechanical instabilities. However, the shape of the quench current vs. temperature, particularly near the λ -point, is not as predicted by temperature dependence of the conductor critical current. The deviation of the observed temperature dependence from the theoretical prediction can be explained by resistive heating in the coil, which increases as the current rises. The change in the curve shape at temperatures less than 2.17 K is likely due to the improvement of the coil cooling condition in superfluid helium.

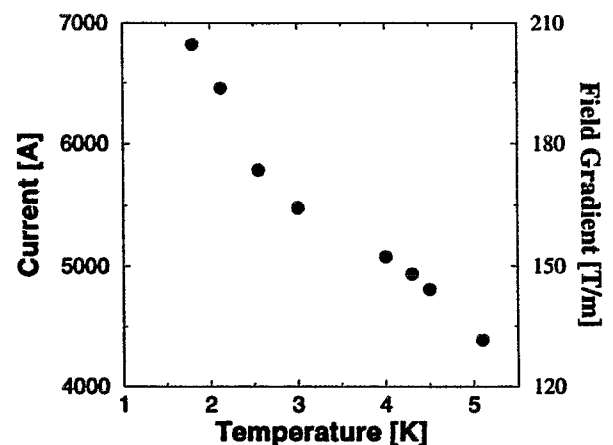


Figure 3: Quench current as a function of magnet temperature (nominal ramp rate of 16 A/s).

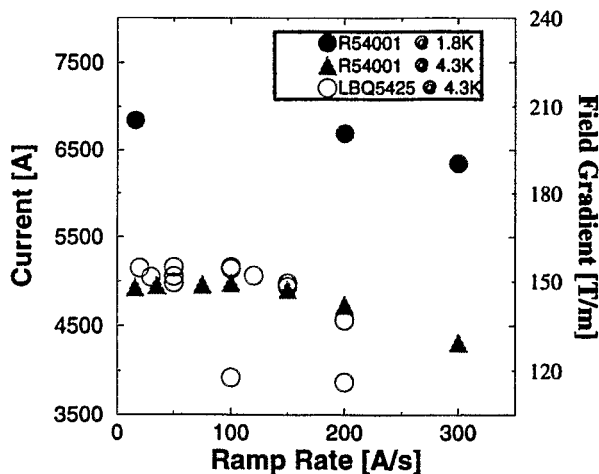


Figure 4: Quench current as a function of ramp rate. (Data were not taken at 1.8 K for LBQ5425.)

3.3 Ramp Rate Dependence of Quench Current

Figure 4 shows the dependence of the magnet quench current vs. the ramp rate for LBQ5425 and R54001. For both magnets, the quench current does not decrease with increased ramp rate up to 150 A/s. At higher ramp rates, cable heating due to AC losses decreases quench current. As above, we observed erratic quench behavior with LBQ5425 during these ramp rate tests.

Figure 5 shows the quench current normalized to the 16 A/s value vs. ramp rate for R54001 measured at 4.3 K in normal helium and at 1.8 K in superfluid helium. An improvement in the quench current at high ramp rate in superfluid helium is evident due to improved coil cooling conditions.

4 CONCLUSION

Two Fermilab low- β quadrupoles have been tested in superfluid helium. LBQ5425, a production spare for the Tevatron, reached 200 T/m gradient at 1.8 K but exhibited erratic quench behavior. R54001, a magnet from the same design but with improved coil mechanical support, also reached 200 T/m at 1.8 K and had significantly better quench behavior. Most quenches in R54001 above 1.9 K were generated by resistive coil heating. We also observed a significant improvement in the coil cooling condition and quench performance of this magnet in superfluid helium.

5 ACKNOWLEDGMENTS

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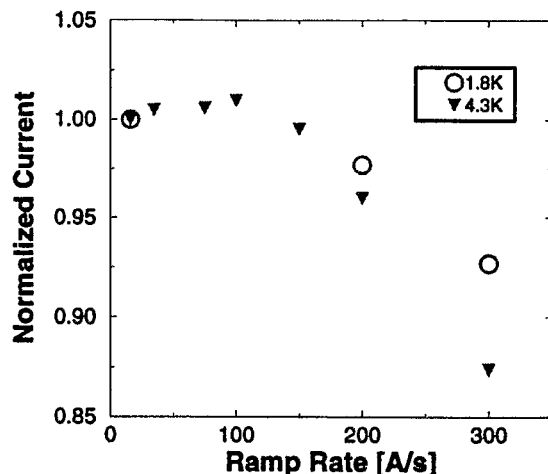


Figure 5: Magnet quench current normalized to the 16 A/s value vs. current ramp rate.

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FIELD QUALITY OF THE MAIN DIPOLE MAGNETS FOR THE LHC ACCELERATOR

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ABSTRACT

Short and long dipole model magnets for the LHC are measured in detail to optimize the field quality to the different operation phases of the accelerator. We will report on recent progress to understand and quantify effects related to the magnetization of the superconductor, to time dependent effects at fields corresponding to beam injection and to behaviour during acceleration. A parametrization for the measurements is proposed. The contributions from any misalignment of the dipoles and correctors magnets is compared to expected field quality of the dipoles.

1 INTRODUCTION

The LHC accelerator will produce head on collisions between beams of 7 TEV protons [1]. Beam optics calculations indicate that the performance of the LHC will be limited by the field quality mainly at low energy through a limitation of the dynamic aperture. Experience from the Tevatron [2] and Hera [3] machines have shown that time dependent effects on the field errors can lead to unstable beams during the injection and at the beginning of the acceleration where an abrupt "snap-back" to the field errors at start of injection takes place. Field ramps necessary for the beam acceleration create coupling currents between the strands of the superconducting cables of the winding, resulting in degradation of the field quality.

The field expansion used is relative to the main field, B_1 , of the magnet at $x = R_{ref} = 10$ mm. Here $n=1$ is a dipole field, $n=2$ is a quadrupole field etc. The b_n and a_n represent the normal and skew relative field errors.

$$B_y + iB_x = \sum_{n=1}^{\infty} C_n \left(\frac{z}{R_{ref}} \right)^{n-1} = B_1 \sum_{n=1}^{\infty} \frac{(b_n + ia_n)}{10^n} \left(\frac{z}{R_{ref}} \right)^{n-1}$$

This paper will concentrate on field errors that could dominate the behaviour of the LHC at low energy.

2 ALIGNMENT OF THE DIPOLES AND THE END CORRECTORS

The harmonics allowed by the symmetry of the dipole are large at injection field due to the cable magnetization. A misalignment of the dipole containing a multipole of order $n+1$ produces an additional one of order n . Table 1 compares the errors generated by a misalignment of 0.5 mm to the expected random errors

$\sigma(b_n)$ and $\sigma(a_n)$ of the dipole field. Moreover, the decapole and sextupole correctors mounted at each end of the main dipoles will produce similar quadrupole and octupole terms if misaligned with respect to the dipoles. Some care will therefore be needed to avoid systematic misalignment of the end correctors with respect to the main dipoles.

Table 1. Comparison for even harmonics (in units of 10^{-4} of injection field at 10 mm radius), between the errors resulting from a misalignment of 0.5 mm of the dipole and the standard deviation (σ) of the field errors expected in the dipoles.

n	b_{n+1} expected in the dipole	$\sqrt{b_n^2 + a_n^2}$ for 0.5mm mis.	$\sigma(b_n)$	$\sigma(a_n)$
2	-3.9	0.4	0.4	1
4	0.25	0.05	0.1	0.1
6	-0.026	0.008	0.06	0.01

3 INTERSTRAND EDDY CURRENTS

Current ramps induce field distortions for all harmonics. These effects are dominated by coupling currents inversely proportional to the interstrands resistance of the superconducting cables [4]. Table 2 shows, for three different 10 m long models, that the deviations are important for all harmonics. Systematic field errors due to this effect could happen for instance due to the bending of the dipoles to the radius of curvature of the LHC.

A tight control of the LHC dynamic aperture is required. A systematic study based on these measured values was performed on the resulting loss of dynamic aperture at the beginning of the acceleration [5]. It makes the assumption that an exponential increase of the acceleration rate gives a constant perturbation all over the acceleration.

Table 2. Field errors due to cable coupling currents (in units of 10^{-4} of the injection dipole field at 10 mm) at 8T/20min ramp-rate, in four 10 m long dipole models. Also reported the average interstrand resistance as deduced from loss measurements.

n	MTP1A1	MTP1A2	MTP1A3	MTP1N2
b_1	88	10.8		9.8
b_2	-0.96	-1.11	-0.052	1.32
a_2	10.8	0.89	1.052	1.06
b_3	4.6	0.95	0.3	0.55
a_3	-0.16	-0.19	-0.035	0.14
b_4	0.45	-0.13	0.019	-0.047
a_4	0.82	0.08	0.13	0.24
b_5	2.08	0.05	0.002	-0.09
a_5	-0.34	-0.02		0.026
$R_{int}(\mu\Omega)$	1.6	6.8	14	6.7

This study shows that a spread comparable to the imperfections of MTP1A1 would significantly reduce the LHC dynamic aperture during the ramp. What matters is really the spread of all harmonic terms for the whole population of dipoles, the average b_3 and b_5 can indeed be corrected with the end correctors. A spread comparable to the imperfections of MTP1A2 or even better of MTP1A3 would not degrade significantly the dynamic aperture. MTP1A2 might be seen as an acceptance limit at the expense of a significant increase of the already long ramp time. These conclusions should be confirmed by further tracking studies.

4 MAGNETIZATION DECAY AND "SNAP-BACK"

LHC superconducting dipoles are characterized by a significant drift of the magnetic field when the current is constant, with typical time scales in the order of several minutes to several hours [4]. At the restart of the ramp (beam acceleration phase) the field quickly bounces back, "snaps back", reaching within 25 to 40 A (approximately 20 to 30 mT of field change) the original value at the start of the injection plateau. This can be of concern for the start of beam acceleration. All harmonics are subject to the field drift. On allowed harmonics the drift is systematic in the direction decreasing the contribution from the filaments magnetization, while it is random on non-allowed harmonics.

In accordance with previous measurements on Tevatron[2], HERA[3] and SSC[6] magnets, we have found that the amount of decay on all harmonics (and thus the snap-back) depends on the time spent at high current, and on the number of operating cycles. Typically, for the LHC dipoles, the decay requires very long flat-tops at high current to saturate (times longer than ½ hour). A few cycles (of the order of 5) to nominal operating current produce an asymptotically stable situation. Quenching resets previous history effects. For machine operation this implies that magnets who have followed different histories will have a different behaviour at injection field.

We speculate that the long term drift and the snapback are associated with the interaction between the changing cable internal field associated with the strand current distribution and the magnetization of the single strands. We have started measurements of typical field decays and snap-back amplitudes, investigating on the influence of several parameters such as pre-cycles and temperature drifts in order to establish suitable procedures to minimize this effect. The main findings of these studies are listed below:

- Snap-back is associated with field change rather than with field change rate. This is shown in Fig. 1 where measurements in a 10 m long prototype are shown in the case of different ramp-rates after the injection plateau. The magnitude of the snap-back is approximately the same in all cases, but the time span for the snap-back scales inversely with the current ramp-rates.
- A pre-injection plateau at a lower current reduces the amount of field decay and snap-back at injection. A $\approx 25\%$ smaller snap-back was measured on b_3 after a 15 minutes pre-injection plateau. The reduction is not a critical function of the pre-injection current.
- The magnitude of the snap-back on the allowed harmonics at different temperatures scales as the ratio of the contributions of DC magnetization to the harmonic at the same temperatures. This demonstrates the relation between field decay, snap-back and DC magnetization.

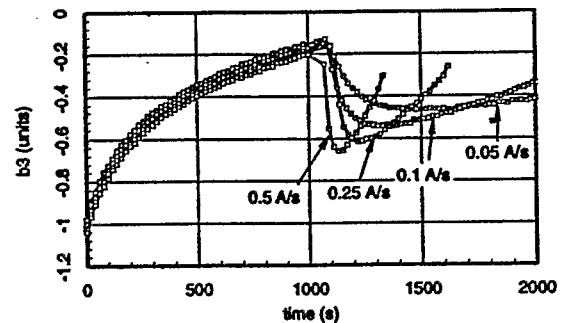


Figure 1. Snap-back in a 10 m long dipole prototype (MTP1N2) at different ramp-rates after the injection plateau.

In summary, a pre-injection plateau before the injection phase decreases the snap-back. The residual magnitude can be accommodated by slow ramps giving time for the corrector scheme to react. Typical magnitudes for the snap-back of the first harmonics in the LHC 10 m long prototypes, measured after several days of operation, are reported in Tab. 3. Systematic effects are expected only on the allowed harmonics.

Table 3. Change of multipoles due to snap-back for the LHC dipoles, estimated from the measured snap-back in 10 m long prototypes. All quantities in units of 10^{-4} of the main field, at 10 mm reference radius. The standard deviations expected from geometrical errors in the magnets series are quoted as a reference.

component	Δ (Snap-Back)	σ (geometric)
b_2	± 0.1	0.4
a_2	± 0.5	1
b_3	+0.8	0.5
a_3	± 0.2	0.15
b_4	± 0.5	0.1
a_4	0.0	0.1
b_5	-0.1	0.05
a_5	0.0	0.04

5 PARAMETRIZATION OF THE MEASUREMENTS

In a most general way, we can write that the field \mathbf{B} produced by a magnet in arbitrary operating conditions is approximated by:

$$\mathbf{B} = \mathbf{B}_{geom} + \mathbf{B}_s + \mathbf{B}_{\Delta x \Delta y} + \mathbf{B}_{AC} + \mathbf{B}_M + \mathbf{B}_{\Delta M}$$

where we have evidenced the various contribution of winding geometry \mathbf{B}_{geom} (including the linear iron contribution), iron saturation \mathbf{B}_s , winding movements $\mathbf{B}_{\Delta x \Delta y}$, coupling current \mathbf{B}_{AC} , filaments magnetization \mathbf{B}_M and long term magnetization drift and snap-back $\mathbf{B}_{\Delta M}$. For each of them we are testing suitable parametric dependencies based on existing models or empirical fits.

The present scaling for the coefficients C_n is given in Tab. 4, where C_n indicates the main field component.

Table 4. Scaling of the field contributions proposed for a parametric representation.

	Main field	Field errors ($n \geq m+1$)
\mathbf{B}_{geom}	$C_n^{geom} = \gamma_n I$	$C_n^{geom} = \gamma_n C_n$
$\mathbf{B}_s + \mathbf{B}_{\Delta x \Delta y}$	$C_n^s = \sigma_{n,1} I^2 + \sigma_{n,2} I^3 + \sigma_{n,3} I^4$	$C_n^s = \sigma_{n,2} C_n^2 + \sigma_{n,3} C_n^3 + \sigma_{n,4} C_n^4$
\mathbf{B}_{AC}	$C_n^{AC} = \tau_n \frac{dI}{dt}$	$C_n^{AC} = \tau_n \frac{dC_n}{dt}$
\mathbf{B}_M	$C_n^M = \mu_{n,0} + \frac{\mu_{n,1}}{I}$	$C_n^M = \mu_{n,0} + \frac{\mu_{n,1}}{C_n}$
$\mathbf{B}_{\Delta M}$	$C_n^{\Delta M} = \Delta_n [1 - e^{-(t-t_0)/\tau}]^{(n)}$ $C_n^{\Delta M} = \Delta_n [1 - (I(t) - I_0) / \Delta I]^{(n)}$	$C_n^{\Delta M} = \Delta_n [1 - e^{-(t-t_0)/\tau}]^{(n)}$ $C_n^{\Delta M} = \Delta_n [1 - (I(t) - I_0) / \Delta I]^{(n)}$

⁽ⁿ⁾ During a $I = I_0$ current plateau starting at $t = t_0$

⁽ⁿ⁾ At ramp restart, between starting current I_0 and current $I_0 + \Delta I$

6 IMPLICATION FOR THE SERIES MEASUREMENTS

A coils assembly able to measure the integral of the field is under construction [7]. It is composed of 14 measuring lengths separated for mechanical reasons by gaps equal to the twist pitch of the cable. It will allow much faster measurements for the series. The errors avoided by an integrated measurement are listed below.

- The residual error from a measuring coils having in average the length of the twist pitch that varies slightly between cable manufacturers is larger than the standard deviation from magnet to magnet.
- Time dependent effects have a larger amplitude when approaching the ends of the magnets.
- Point like variations of eddy currents have been measured.
- The randomness of the snap back effects is not yet understood. We believe that the standard deviation measured along the 15 m of the magnets can be as high as the deviation between magnets.

The request to be able to measure the dipole axis and verify the centering of the spool piece correctors is not possible with a integral shaft. These measurements will be performed during warm tests of the final assemblies with the help of a short coil. Differences between the axis positions in warm and cold conditions will be verified on a reduced number of magnets.

7 CONCLUSION

Exploratory measurements of both short and long models of the LHC superconducting dipole magnets indicate the importance of methodically separating different effects. The reproducibility of the time dependent effects and dynamic effects may dominate the low field quality with respect to superconductor magnetization and harmonics due to the conductors locations. We have defined a model for the magnetic measurements in order to quantify and eventually correct these effects that could limit the dynamic aperture of the LHC machine at low energy.

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Dynamic Aperture Limitations of the LHC in Physic Conditions due to Low-Beta Insertions

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Abstract

The dynamic aperture of LHC in collision is determined by the multipole field errors of the low-beta triplet quadrupoles. Their effect, combined with that of the systematic and random magnetic imperfections expected in all machine magnets is analysed here. The effect of crossing angle is taken into account.

1 INTRODUCTION

The effects of the multipole components in the low- β quadrupoles was investigated some time ago [1] on a preliminary version of LHC, with a short term tracking. It was noticed that the errors in the triplet dominate the non-linear dynamics. As the field errors are now better known, it is important to remake the analysis with long term tracking including synchrotron oscillations. This is the subject of this paper, it is shown that the triplet errors still dominate with the latest LHC design.

2 THE LHC LATTICE.

The LHC lattice we consider is that described in the conceptual design report [2]. It is labelled version 4.2 and has four physics insertions where β^* has a value of 0.5m in physics conditions in both planes, in order to achieve a luminosity of about $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$. This low β^* value is obtained by means of highly focusing quadrupoles triplets [2][3]. As the free space between the interaction point and the first low- β quadrupole is about 20m, the peak β -value in the triplets is of the order of 4000m. This gives an enormous importance to the multipole components at this location as shown below.

3 FIELD ERRORS IN THE LHC MAGNETS.

The field errors are divided into :

- average or design errors, associated with a given magnet type. They obey symmetry rule, e.g. in the dipoles, all b_n with n even change sign when going from an inner aperture to an outer one.
- uncertainty on average errors. This is indicated by a \pm sign in the tables [2]. They are randomly assigned to each LHC arc.

- random errors from magnet to magnet, distributed according to a Gaussian law cut at 3σ , except for the low- β quadrupoles where the cut is at 2σ .

An input for the MAD program [4], used for the optics calculations and trajectory tracking, has been made to specify all errors at a time, so that the distribution of a given error remains the same whatever the others. This procedure is a little heavy but it makes it possible to obtain the effect of any multipole component in any context.

3.1 Arc cell magnets.

The multipole errors in the dipoles and the arc quadrupoles are those given in the "yellow book" [2]. It can be shown with equation 1 below that the quadrupole errors are only a few percents of the dipole errors and hence negligible, except b_8 which will be studied separately.

3.2 Other LHC magnets.

Quadrupoles similar to the arc cell quadrupoles are used to match the insertions and dispersion suppressors. They behave similarly to the arc cell quadrupoles.

The two LHC beams are brought in collision by means of dipole magnets sitting between the dispersion suppressors and the insertions. Under physics conditions the value of the β functions in these dipoles is large. Their effect on the dynamic aperture is then important. As they have not yet been designed, no error table is available and they are not considered here.

In the cleaning, dump and RF insertions there are conventional quadrupoles. Their field errors are determined by the design and are known within one percent. Their effect is examined below.

3.3 Low- β quadrupoles.

The multipole errors in the low- β quadrupoles are given in table 1 [3]. In order to compare them with the field errors in the dipoles, an equivalent dipole error $b_{n,q,d}$ can be defined :

$$b_{n,q,d} = b_{n,q} \frac{K r l_q}{\theta} \quad (1)$$

where K is the normalised quadrupole gradient (0.00972 m^{-2}), l_q the quadrupole length (5.5m), θ the angle of the dipole (5mrad) and $b_{n,q}$ the field error in the dipoles [2]. It can be noticed that the mean b_8 and b_{10} , which are the only important components, are smaller

Table 1: Multipole errors in the low-beta quadrupoles in physics. The errors are field errors at 10mm from the quadrupole axis, in units of 10^{-4} . b11, b12 and b13 are zero.

n	Mean		Random	
	b_n	a_n	b_n	a_n
3	0	0	0.68	0.68
4	0	0	0.23	0.23
5	0	0	0.08	0.08
6	-0.008	0	0.028	0.028
7	0	0	0.004	0.004
8	0	0	0.0009	0.0009
9	0	0	0.0002	0.0002
10	-0.005	0	0.00005	0.00005

than those in the main dipoles by one order of magnitude. However, in order to compare their effects, the non-linear aberrations, which scale with $Krl_q\beta^2$ for b_n have to be compared.

For the low- β quadrupoles, β_x is 4000m in half of them and 1000 in the others with opposite sign. The value of $b_n Krl_q\beta_x^2$ merely summed over all triplet quadrupoles is $4.3 \times 10^6 m^3$. For the dipoles the similar quantity $b_n\theta\beta_x^2$ is only $4.1 \times 10^4 m^3$. Such a calculation is extremely crude but it shows that the effect of the quadrupoles on the non-linear betatron oscillations is expected to be larger than that of the dipoles by about two orders of magnitude.

The accurate estimation of the effect of these errors is done below by long term tracking of particle trajectories.

4 DYNAMIC APERTURE.

The dynamic aperture is defined as the maximum action in both planes simultaneously for which transverse betatron oscillations remain stable over 10^5 turns. Tracking the particle trajectories is done with the MAD program [4]. The trajectories start at the interaction point IP1 where both β^* are 0.5m. Both starting coordinates are equal, the initial relative momentum deviation is 3.6×10^{-4} and all canonical momenta are zero. Under these conditions, the initial amplitude (the same value in x and y) has to be multiplied by $\sqrt{2}$ in order to give the maximum radius in the {x,y} plane.

The tunes are $Q_x=63.28$, $Q_y=63.31$, $Q_s=0.001$ (16MV RF voltage for an energy of 7TeV).

Neither alignment errors nor quadrupole errors are included ($a_2=b_2=0$). The b3 and b5 compensators, the so called "spool pieces" in the dipole ends, are turned on. Both tune derivatives are set to 1.0 after the errors are introduced in order to take into account the tune modulation associated with the synchrotron oscillations. A physical aperture limitation has only been introduced only for the case of individual multipoles (section 7).

5 DYNAMIC APERTURE DUE TO THE ARC CELL ERRORS.

The dynamic apertures as well as their average and standard deviation obtained for 10 random distributions of the errors

in the arc cell magnets are listed in table 2. Under physics

Table 2: Dynamic aperture associated with the dipole and arc quadrupole errors, both systematic and random.

seed	10^4 turns		10^5 turns	
	x at QF/mm	n_σ	x at QF/mm	n_σ
eopt				
1	11.9	39.2	9.6	31.7
2	12.7	42.0	11.6	38.3
3	11.9	39.2	11.3	37.3
4	13.0	42.9	11.9	39.2
5	10.7	35.5	9.1	29.9
6	12.7	42.0	11.6	38.3
7	11.9	39.2	11.0	36.4
8	15.6	51.3	14.7	48.5
9	13.9	45.7	11.3	37.3
10	11.0	36.4	10.7	35.5
x_{av}	12.5	41.4	11.3	37.3
σ_x	1.3	4.4	1.4	4.7

conditions the primary collimators limit the aperture to 6σ at most and the betatron motion must be regular enough for the particles scattered by the primary collimators and reaching 10σ [6]. This is largely fulfilled with a dynamic aperture of 10σ at 10^4 turns.

The results in table 2 show that we are well above these limits. If the systematic multipoles in the warm quadrupoles are added, the reduction of the dynamic aperture is of the order of 5%, i.e. negligible.

6 DYNAMIC APERTURE DUE TO THE LOW- β QUADRUPOLES

6.1 No crossing angle

Table 3: Average dynamic aperture associated with the dipole, arc quadrupole and low- β quadrupole errors, both systematic and random.

	10^4 turns		10^5 turns	
	x at QF/mm	n_σ	x at QF/mm	n_σ
x_{av}	4.0	13.3	3.8	12.5
σ_x	0.2	0.6	0.2	0.7

The summary of the results of the computation of the dynamic apertures obtained for the same 10 random distributions of the errors as in the previous section is shown in table 3. Its variability with the distribution of the random errors is considerably reduced compared with that in table 2: It is clear that the multipoles in the low- β quadrupoles dominate the dynamic aperture.

6.2 Crossing angle

The LHC beams cross at a full angle of $200\mu rad$. This is achieved with a local closed orbit distortion made independently in both rings. The angle of $100\mu rad$ in each ring makes a closed orbit excursion of about 5mm in the low- β quadrupoles. Because of the multipole components in the quadrupoles, lower order multipoles are created (this is often referred to as "multipole feed-down"). The summary of

Table 4: Dynamic aperture associated with the dipole, arc quadrupole and low- β quadrupole errors, both systematic and random. Crossing angle scheme on. If the warm quads are added, x_{av} increases by 3%.

	10 ⁴ turns		10 ⁵ turns	
	x at QF/mm	n_σ	x at QF/mm	n_σ
x_{av}	3.2	10.6	3.0	9.9
σ_x	0.2	0.6	0.2	0.6

the results of the computation of the dynamic aperture obtained by turning the crossing angle scheme on, with the errors described in the preceding section, is shown in table 4. A reduction by a factor 0.8 is observed.

In fact, turning the crossing angle scheme on produces little effects on the linear optics. The tunes-shifts are 1.1×10^{-4} in the horizontal and -3.3×10^{-3} in the vertical plane, they are compensated by means of auxiliary quadrupoles [7] which produce a negligible optics perturbation. The chromaticities do not change. A negligible linear coupling is due to the feed-down effect of the skew multipole components (the resulting width of the coupling resonance is 3.4×10^{-3} , it is not compensated). The observed effect must be then due to the "feed-down" of low-order multipoles. Indeed the anharmonicity increases by about 20% on average when the crossing angle is turned on.

7 EFFECT OF INDIVIDUAL COMPONENTS IN THE LOW- β QUADRUPOLES.

The dynamic aperture was computed for each multipole component in the low- β quadrupoles in presence of the chromaticity sextupoles only. The results are shown in table 5. If the aperture limitation due to the vacuum chamber is introduced both in the arc quadrupoles and in the triplet, the results are identical, with the present amplitude steps.

The dominating multipole is b10. This can be easily understood by computing the field errors at the maximum amplitude in the low- β quadrupoles. For the case where the maximum horizontal amplitude is 3.4mm in the arc where β_x is 181m, it is 16mm in the low- β quadrupoles where β_x is 4000m. The vertical field error at 16mm is then -0.084×10^{-4} for b6 and -0.34×10^{-4} for b10. It is clear that the reference radius of 10mm does not give a right idea of the relevant errors in the low- β quadrupoles.

8 SOLUTIONS TO THE MULTIPOLE PROBLEM.

In order to obtain a dynamic aperture of 10σ , it is probably not possible to re-optimize the field error in the triplet, as a reduction of b10 by a factor of about 5, which may not be possible, gives a just acceptable dynamic aperture (table 5). A safe solution is to use compensators for the most dangerous component and this is presently under study. With the present field errors, three b10 compensators placed close to

Table 5: Dynamic aperture associated with low- β quadrupole systematic errors and chromaticity sextupoles only (upper), random b6 and b10 are added in the lower part (the r.m.s. deviation is indicated with \pm). Crossing angle scheme on \hat{x} is the maximum amplitude at QF. The dynamic aperture according to the Liapounov exponent is given in Liap.

mult.	10 ⁴ turns		10 ⁵ turns		Liap.
	\hat{x} /mm	n_σ	\hat{x} /mm	n_σ	
b6	7.1	23.3	7.1	23.3	23.3
b10	3.4	11.2	3.4	11.2	11.2
b6+b10	3.7	12.1	3.4	11.2	11.2
b6+.2b10	4.8	15.9	4.5	14.9	
b6+r	4.8±0.4	15.9	4.3±0.3	14.3	
b10+r	3.4±0.0	11.2	3.3±0.1	11.0	
b6+b10+r	3.4±0.1	11.2	3.3±0.1	10.7	

each quadrupole are probably necessary.

Eventually increasing β^* reduces the maximum β 's in the low- β quadrupoles. The dynamic aperture computed with a β^* of 1m in presence of b6, b10 and chromaticity sextupoles increases to 28σ , which shows its sensitivity to β^* .

9 CONCLUSION.

The multipole errors in the low- β quadrupoles limit the dynamic aperture under physics conditions because of the large values of the β -functions in these quadrupoles. They are one of the main contributors to the lower limit on β^* . A compensation system, necessary to reach very small β^* is now under study.

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FIELD QUALITY IN SUPERCONDUCTING MAGNETS FOR LARGE PARTICLE ACCELERATORS*

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Abstract

The expected field quality in superconducting magnets for large particle accelerators has improved over a period of time due to the development and application of a number of techniques. These design techniques will be described and as an example the expected harmonics in an improved design of the 50 mm aperture dipole magnets for the Superconducting Super Collider (SSC) will be presented. This field quality is based on experience with the Relativistic Heavy Ion Collider (RHIC) and SSC magnet programs. The initial results of a design approach will also be presented where the first design itself is adopted to produce good field quality in RHIC D0 insertion magnet.

1 INTRODUCTION

The purpose of this paper is (a) to give a brief outline of the concepts which either have been used or can be used in designing and then assuring good field quality in accelerator magnets, (b) to present a design concept which avoids the need of a normal magnetic design iteration which is both time consuming and expensive, (c) to present methods which can be used to control and match the field quality in the magnets built by several vendors, (d) to present techniques which overcome the influence of normal errors in parts and assembly and (e) to apply these concepts in a quantitative way to estimate field quality in a magnet which can be used as a reference for the next generation machines.

The field quality in accelerator magnets is characterized in terms of the normal and skew harmonics, b_n and a_n . They are defined in the following expression

$$B_y + iB_x = 10^{-4} \times B_{R0} \sum_{n=0}^{\infty} (b_n + ia_n) [(x + iy) / R]^n,$$

where B_x and B_y are the components of field at (x,y) and B_{R0} is the magnitude of the field due to fundamental harmonic at a "reference radius" R .

2 METHODS FOR CONTROLLING FIELD QUALITY

The methods for controlling field quality are only discussed briefly here (see references [1] to [9] for details).

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2.1 Allowed Harmonics and Pre-stress on Coils

Two allowed harmonics can be adjusted by adjusting (a) the coil-to-midplane gaps (or midplane shims) and (b) the coil pole shims. This is an efficient technique for small adjustments in lower order harmonics (e.g., b_2 and b_4 in dipoles). It may be used either for normal cross-section iteration or to compensate for the differences in field quality of the magnets built by different vendors based on the same magnetic design.

The pre-stress on the collared coil, however, will change if there is a net change in the combined thickness of the pole and midplane shims. A small variation in pre-stress may be tolerated but if it is larger than a few kpsi, then to avoid it one should adjust the coil curing pressures to change the cured coil size. The first magnet of a series usually requires larger adjustment in field harmonics. This is further complicated by the experience that the desired pre-stress on the coils may also not be obtained by the nominal size pole shims. A change in pole shims in order to get the desired pre-stress also changes the field quality.

To deal with the above difficulties, an approach has been developed for the RHIC 100 mm aperture insertion dipole D0 [5] with the goal that the first design itself can be adopted for production by applying small adjustments which are part of the initial design. A third parameter to provide the desired pre-stress on the coils is obtained by increasing the effective thickness of one (or more) selected wedge(s) by changing the number of layers of insulation on it. This resulted in low field harmonics in the body of the magnet as good as those expected after a number of cross-section iterations. Also the magnet had the desired pre-stress on the coils. The first D0 magnet is being used in the RHIC machine. Based on measurements in the first magnet, small adjustments were made in the thickness of the midplane insulation caps and pole shims to compensate the end harmonics and 0.4 mm radius magnetic rods were inserted in the saturation control holes to reduce small values of saturation induced b_2 and b_4 harmonics.

A regular cross-section iteration generally requires a large mechanical change in several wedges and is associated with (a) a change in tooling and (b) a change in the end design. Both of these are time consuming and expensive and are avoided in the above approach. Moreover, this approach requires only small mechanical changes and therefore has a better chance of succeeding.

2.2 Integral Transfer Function

A difference in the integral transfer function (or effective length, $\delta L/L$) between the magnets built by different vendors has been observed in HERA magnets [10]. A similar situation was observed in RHIC when different length magnets were built for the insertion regions. Even though the magnets were built by the same vendor (Northrop Grumman Corporation) and on the same design, the integral transfer function was off by a small amount from the expected values.

In most magnets the coil ends are enclosed by non-magnetic stainless steel laminations instead of magnetic low carbon steel in order to reduce the field on the conductor. It is proposed here that the axial location where the transition between the low carbon steel laminations to stainless steel laminations takes place be used as a parameter to adjust the integral transfer function. This adjustment may be applied after the first few magnets are built and measured. A rough estimate suggests that a 50 mm adjustment on each end should be adequate in most cases.

An adjustment in the change in transfer function between low field and high field due to iron saturation can be obtained by changing the iron packing factor (i.e. the number of yoke laminations for the same yoke length). This technique has been used in adjusting the integral transfer function in the different length RHIC magnets.

2.3 Skew Quadrupole in Dipole Magnets

The presence of a skew quadrupole harmonic in dipole magnets reflects a top-bottom asymmetry. A difference in the yoke length between the top and bottom halves can similarly be used to compensate for this asymmetry [8]. The current dependence in the variation of this harmonic can be adjusted by adjusting the yoke packing factor between the two halves [7].

2.4 Twist in Field Angle of the Magnets

A change in field angle (or twist) along the axis of RHIC magnets has been reduced by applying slanted welds on the outer diameter of stainless steel shell. The weld moves the magnet in the direction it is applied. Slanted weld in opposite directions on the two sides of the shell effectively generate a torque which takes out the twist of the magnet.

2.5 Harmonic Correction after Initial Assembly

In a perfect design, the field quality is limited only by the errors (tolerances) in parts and assembly. To overcome these limitations, a "tuning shim method" is being used in 130 mm aperture RHIC insertion quadrupoles [6]. Eight tuning shims of variable iron thickness are inserted to cancel out the measured values of eight harmonics. The test results of the first five magnets show

that the harmonics can be corrected within their repeatability and measurement errors.

2.6 Current Dependence in Field Harmonics

At high field the harmonics change as a function of current due to iron saturation and also due to Lorentz forces on the coil. A larger variation is expected in RHIC type magnets where the iron contribution to the total field is large. However, the saturation induced harmonics can be minimized by forcing the yoke iron to saturate uniformly with the help of holes, cutouts, etc. in the yoke geometry. A good parameter to examine [9] the iron saturation is $(\mu-1)/(\mu+1)$. The value of this parameter is ~ 1.0 at low field (no saturation or $\mu > 1000$) and zero at very high field (complete saturation or $\mu = 1$). The variation in this parameter as a function of azimuthal angle should be small, particularly near the yoke inner radius. In RHIC magnets, despite a large contribution from the yoke, the saturation induced harmonics are small (see references [2] to [5]).

3 EXPECTED FIELD QUALITY IN SSC-TYPE DIPOLE MAGNETS

The expected field quality presented in this section for a 50 mm aperture SSC-type dipole is significantly better than previously assumed in SSC beam tracking studies [11]. The improvement comes from (a) the use of the above methods for controlling field quality (b) a feedback from the measurements in the SSC prototype magnets and (c) the use of revised methods of estimating field errors. The methods that are commonly used for estimating the field harmonics tend to over-estimate the expected field errors. This has been discussed in detail in reference [4] where the possible sources for this over-estimate are discussed and the experience with RHIC and SSC magnets is presented. The expected field harmonics are described by (i) the expected mean (ii) the uncertainty in mean and (iii) the expected RMS width or sigma (σ).

The expected integral values of field harmonics in a 50 mm aperture, 15 meter long, 2-layer SSC-type dipole magnet built along the lines discussed here are given in Table 1 at 10 mm reference radius for an operating range of 2000 A (2 tesla) to 6600 A (6.6 tesla). At injection (660 A), the allowed harmonics are dominated by the effects of the persistent currents. Since the issue of persistent current is not discussed here, the harmonics are given only at or above 2000 A, where their influence is negligible. The last row gives the expected systematic difference in the integral transfer function or effective length ($\delta L/L$) between the magnets built by two or more vendors; the σ is for the whole accelerator. For reference, SSC specifications at high field are also given.

The use of the RHIC insertion quadrupole type tuning shims [6] is not assumed here. The tuning shims may

reduce those harmonics selected for correction by about a factor of five, as long as they are reproducible.

Table 1. The expected normal (b_n) and skew harmonics (a_n) in 50 mm aperture SSC type dipole magnets at a reference radius of 10 mm in an operating range of 2000 A to 6600 A. $\langle b_n \rangle$ and $\langle a_n \rangle$ are the expected mean, $d(b_n)$ and $d(a_n)$ are the uncertainty in the mean and $\sigma(b_n)$ and $\sigma(a_n)$ are the expected sigma. The last two columns show the SSC tolerances/specifications for the mean (ssc $\langle \rangle$) and sigma (ssc σ) at high field. Since $\langle b_n \rangle$ and $\langle a_n \rangle$ are zero $d(b_n)$ and $d(a_n)$ should be compared to the ssc $\langle \rangle$. The last row gives the expected systematic difference in the effective length ($\delta L/L$) between the magnets built by two or more vendors and σ for the whole accelerator.

n	SSC Expected Harmonics			SSC Tolerances	
	$\langle b_n \rangle$	$d(b_n)$	$\sigma(b_n)$	ssc $\langle \rangle$	ssc σ
1	0.0	0.04	0.2	0.04	0.5
2	0.0	0.2	0.4	0.8	1.15
3	0.0	0.002	0.03	0.026	0.16
4	0.0	0.02	0.05	0.08	0.22
5	0.0	0.001	0.002	0.016	0.02
6	0.0	0.002	0.005	0.013	0.02
7	0.0	0.0004	0.001	0.01	0.01
8	0.0	0.0005	0.001	0.02	0.0075
n	$\langle a_n \rangle$	$d(a_n)$	$\sigma(a_n)$	ssc $\langle \rangle$	ssc σ
1	0.0	0.04	0.5	0.04	1.25
2	0.0	0.02	0.15	0.032	0.35
3	0.0	0.01	0.07	0.026	0.32
4	0.0	0.002	0.02	0.02	0.05
5	0.0	0.002	0.008	0.016	0.05
6	0.0	0.001	0.005	0.013	0.01
7	0.0	0.0004	0.001	0.01	0.01
8	0.0	0.0004	0.001	0.008	0.0075
$\delta L/L$	--	0.0003	0.0004	--	0.0006

4 CONCLUSIONS

The expected field quality shown in Table 1 is generally a factor of 2 to 5 better than that previously assumed in SSC beam tracking studies [11]. Such improvements should have some influence on the expected performance of the machine. Most concepts presented/reviewed in this paper have already been successfully tested in RHIC or SSC magnets. The spirit behind the techniques presented here is that simple mechanical adjustments can make significant improvements in the performance of a series of magnets, if they are planned ahead and made part of the initial design.

5 ACKNOWLEDGMENTS

Discussions with A. Jain, S. Kahn, G. Morgan, P. Thompson, P. Wanderer and E. Willen on the field quality issues have been very useful and contributed to

developing these concepts. The engineering efforts in incorporating these ideas into real magnets from M. Anarella, J. Cozzolino, E. Kelly and J. Schmalzle are also acknowledged. The feedback on field quality requirements in RHIC magnets was provided by S. Peggs and J. Wei. G. Morgan, P. Wanderer and E. Willen made a careful review of this manuscript.

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DESIGN OF A HIGH GRADIENT QUADRUPOLE FOR THE LHC INTERACTION REGIONS

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Abstract

A collaboration of Fermilab, Lawrence Berkeley National Laboratory and Brookhaven National Laboratory is currently engaged in the design of a high gradient quadrupole suitable for use in the LHC interaction regions. The cold iron design incorporates a two-shell, $\cos 2\theta$ coil geometry with a 70 mm aperture. This paper summarizes the progress on a magnetic and mechanical design that meets the requirements of maximum gradient ≥ 250 T/m, operation at 1.8K, high field quality and provision for adequate cooling in a high radiation environment.

1 INTRODUCTION

Fermilab, Lawrence Berkeley National Laboratory (LBNL) and Brookhaven National Laboratory (BNL), have formed a consortium to provide components for the Large Hadron Collider[1] to be built at CERN in Geneva, Switzerland. A proposed U.S. contribution is the high gradient quadrupoles (HGQ) for the interaction regions. These magnets present some formidable challenges. In addition to the large aperture and high gradient, they must operate at superfluid helium temperature with a beam induced heat load of up to 40 W per magnet. The design effort has been underway for less than a year and is far from complete. This paper describes the status of the project as of this date.

2 DESIGN

2.1 Magnetic Design

The HGQ has a two-layer, graded, $\cos 2\theta$ coil with a 70 mm diameter bore, magnetic length of 5.5 m and a gradient in excess of 250 T/m. The two-layer design was chosen in preference to a 4-layer geometry[2]. Relative to a 4-layer design, quench protection is simplified because the inductance is low and the geometry allows for uniformly heating both coils during a quench. This design also rests solidly on our past experience and makes use of existing tooling. The cable uses existing SSC strand, which has a nominal J_c of 2,750 A/mm². The SSC conductor will be used for model magnets and will

eventually be replaced with improved conductor that is now under development. The goal of the program is to produce strand with a current density of 3,400 A/mm². Figure 1 shows the expected short sample (ss) performance using the existing conductor, and Table 1 summarizes the expected magnet performance with existing and improved strand. Table 2 lists the values for the design allowed harmonic coefficients.

Table 1. Short sample (ss) performance parameters.

J_c @ 5T, 4.2K	G_{SS} (T/m)	I_{SS} (A)	B_{SS} (T)
2,750 A/mm ²	251.6	13,874	10.01
3,300 A/mm ²	260.7	14,517	10.39

Table 2. Allowed harmonic coefficients (units 10⁻⁴).

Component	Collision	Injection
b6	0.006	0.01
b10	0.002	0.0015
b14	0.0015	0.001

The two-layer cross section design requires large aspect ratio cables. Samples of both the inner and outer cables have been produced by LBNL and winding studies with the new cables have been successful. Parameters of the improved conductor and cable are given in Tables 3 and 4.

Table 3. Strand parameters for the improved conductor.

Alloy	NbTi
Inner Strand Diameter	0.808 mm
Outer Strand Diameter	0.648 mm
Filament Diameter	9 μ m
Cu/SC Inner	1.3:1 / 1.1:1
Cu/SC Outer	1.8:1
Twist Pitch	10 mm
Critical Field @ T = 4.2K	10.4 T
Critical Field @ T = 1.8K	13.2 T
Critical Current Density	
T = 4.2 K & B = 5/8 T	3,400/1,500 A/mm ²
T = 1.8 K & B = 8/11 T	3,700/1,800 A/mm ²

Table 4. Cable Parameters

	Cable 1 (Inner)	Cable 2 (Outer)
Number of Strands	38	46
Cable Width	15.4 mm	15.4 mm
Keystone Angle	0.99 degrees	0.68 degrees
Mid Thickness	1.457 mm	1.146 mm

The load lines for each coil layer are shown in Fig. 1, indicating the operating margin for a nominal gradient of 235 T/m for the SSC conductor performance parameters. The improved conductor will increase the margin to approximately 11%.

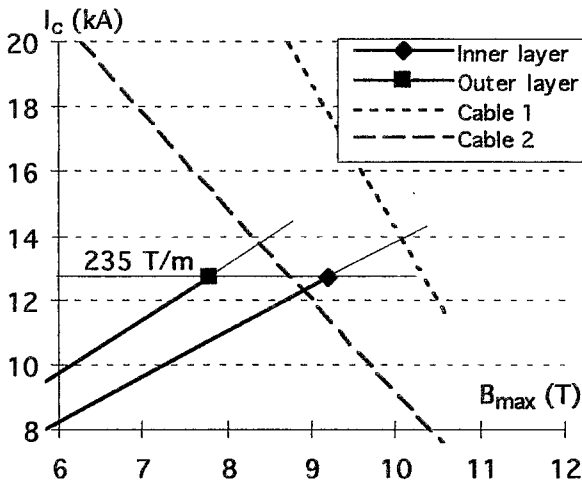


Figure 1. Load lines and critical surfaces for existing strand.

2.2 Mechanical Analysis

The HGQ design incorporates the following mechanical features. A strong 20 mm wide aluminum collar provides significant preload and defines the coil geometry so as to allow warm magnetic measurements to be made. A 2-piece iron yoke surrounds the collared coil and is closed both warm and cold. There is collar-yoke contact under all conditions at the 4 coil mid-planes to reduce coil motion under excitation. Contact is ensured by the larger collar than yoke diameter and the collar deflections due to coil preload. The use of yoke-collar shims is planned to adjust for manufacturing tolerances. A welded stainless steel shell, prestressed to about 200 MPa, provides helium containment and compressive load to maintain yoke closure. Following cooldown the compressive force between iron halves reaches at least 2×10^6 N/m. This is larger than the radial Lorentz force of 1.14×10^6 N/m, so the yoke gap remains closed.

Finite element analysis of this design has been performed[3] using ANSYS to determine limits on the initial collared coil preload and collar-to-yoke spacing. Figure 2 shows azimuthal compressive stress on the coils

during assembly, cooldown to 1.8 K, and excitation to maximum gradient for the lower and upper bound cases having a collar-to-yoke interference (warm, undeflected) of 50 μ m. Figure 3 shows lower and upper bounds for the coil prestress as a function of the yoke-collar spacing. For a given spacing, the lower bound is set by the requirement that the coil stress at a gradient of 250 T/m be ≥ 14 MPa; the upper bound corresponds to the case when the yoke

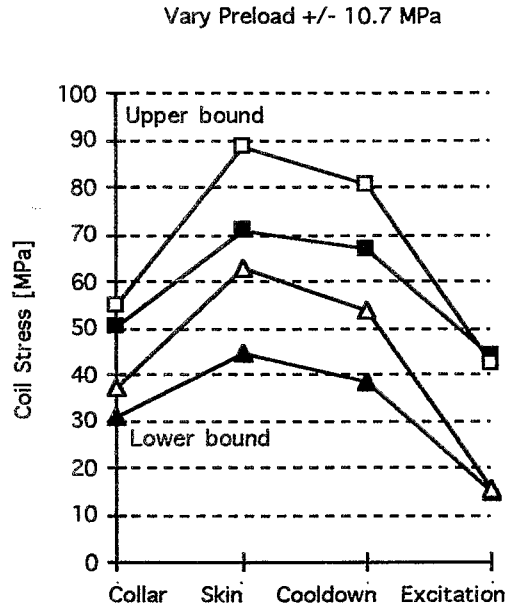


Figure 2. Azimuthal compressive stress on coils during assembly (collaring), cooldown and excitation. Inner (outer) coil stress is shown by open (filled) symbols.

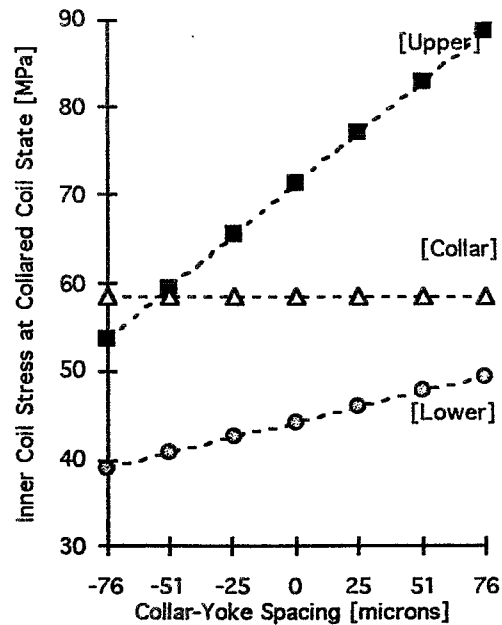


Figure 3. Upper and lower bounds for various collar-yoke spacing and collar stress limit.

mating surface begins to open at room temperature. The final bound is set by the limit on the peak stress in collar material in the initial collared coil state; it is independent of collar-to-yoke spacing. The range of acceptable preload is between the 3 limits. For the design value of 50 μm yoke-collar interference, the range of acceptable collared coil preload is about ± 10 Mpa. For a given preload, as the collar-yoke spacing increases the collar to yoke force decreases which increases the compressive force between iron halves. The upper bound is a function of that force, so the larger the collar-to-yoke spacing, the larger the upper bound point. Also, all lower bound cases have collar-iron contact cold.

2.3 Beam Induced Heat Load

The quadrupole coils in the LHC low- β insertions are subject to a constant heat load up to 40 Watts per magnet and a peak energy density as large as 1.2 mW/g at the coil mid-plane due to secondary particles from beam-beam collisions at the nominal luminosity[4]. The HGQ cables are wrapped with 25 micron thick Kapton tape with a 50% overlap for electrical insulation purposes and one

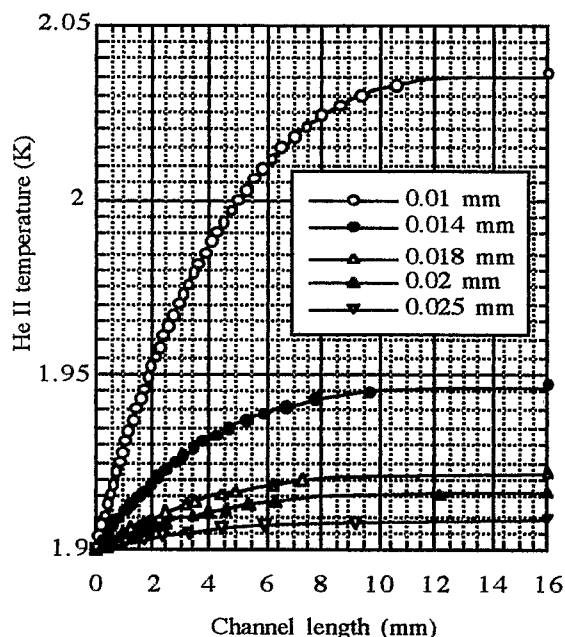


Figure 4. Temperature profile in He II with 2 mm wide cooling channels and varying channel length and height with a surface heat flux of 4.5 mW/cm².

layer of spiral wrapped Kapton for mechanical protection and to provide He II cooling channels. The maximum heat load is at the coil mid-plane and the surface heat flux can be calculated based on the surface area in contact with the He II. The cooling channel formed by the second layer of Kapton insulation film will be deformed during the curing process. The temperature profile within the He II cooling chamber for varying channel gap and length is shown in Fig. 4. A minimum gap of 14 microns is required to keep the He II temperature below 1.95 K.

3 CONCLUSIONS

The project is still within its first year and much progress has been made. Several design issues remain to be considered. Tests of the first model are planned for the Spring of 1997.

4 ACKNOWLEDGMENTS

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SUPERCONDUCTING COIL COMPRESSION BY SCISSOR LAMINATIONS

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ABSTRACT

A new system of coil compression [1] has been designed which uses iron laminations to transfer the pressure from an outer shrink ring to the coil. The laminations are simple circular discs around the coil with the peculiarity that the rim is slightly eccentric as compared to the coil. Successive laminations are mounted with different angular orientations to oppose their eccentricities. The shrink ring pushes these discs inwards against the coil creating compression by a scissor movement. Tests on mechanical models of single as well as multiple aperture magnets have shown it to work as expected. The system has already successfully been applied to several corrector magnets for LHC. The advantages are low cost (suppression of the usual collars), increased coil compression in particular from cooling down, and field enhancement from having the iron close to the coil.

1 INTRODUCTION

Coils of superconducting magnets for accelerators need to be compressed to avoid wire movements which would cause quenches. The prestress is in general obtained by collaring the coils under a press or by shrinking rings around the coils. The iron yoke can help increase the pre-compression if pushed inwards by an outer shell but at the expense of subdividing the yoke into different segments. Scissor laminations offer a simple and economic alternative to regular collar systems.

2 PRINCIPLE

The principle of the scissor laminations is shown in figure 1. The yoke is made of a single type of ring-shaped laminations stacked around the magnet coil. They have the peculiarity that the circular periphery is designed to be slightly off-centre as compared to the magnet centre (~0.5 mm). These laminations are mounted with different angular orientations. The outer shrinking shell presses on the wider side of each lamination only and the latter transmits the compressive force to the coil. The wide side of each lamination simply acts as a local spacer between the shrink ring and the coil and the rest of the circular lamination is free. In general the laminations are oriented per pair of opposing eccentricity each pair compressing the coil thanks to the "scissor" effect. Cooling the magnet to cryogenic temperatures enhances the pre-compression of the coil if the shrink ring is of a strongly shrinking

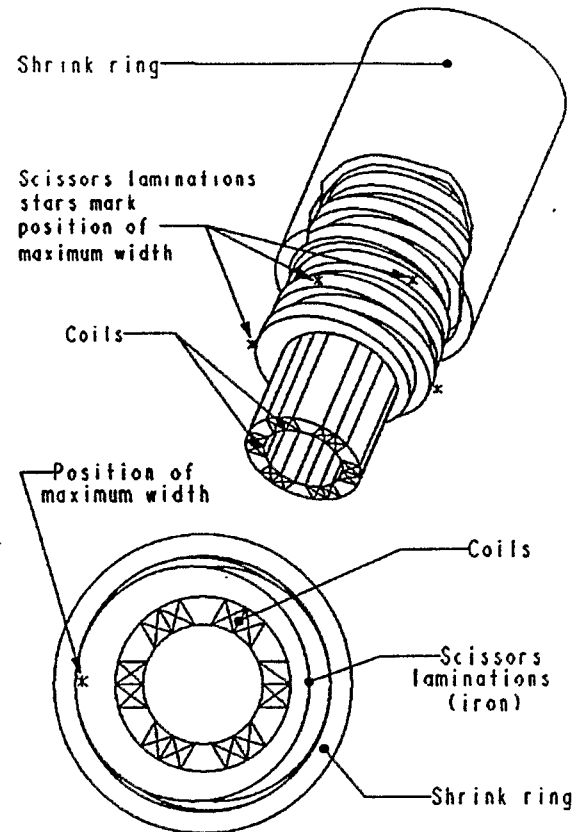


Figure: 1 Principle of Scissor laminations. Shown is sextupole corrector.(eccentricity strongly exaggerated)

material such as aluminium or stainless steel. If desired, one can introduce keys in the laminations to stop the scissor action beyond a certain cooling temperature. Another way to stop it is to choose such an eccentricity that the peripheries of the laminations line up at a well defined cooling temperature. This can be used to build up a circular compression on the laminations that stiffens the yoke; the Lorentz forces that would tend to move the laminations outwards must first overcome that pre-compression before any elastic movement can occur. The scissor system is readily applicable to magnets with more than one aperture, as shown in sections 4, 5.

3 SINGLE APERTURE APPLICATIONS

The method has first been applied to some small superconducting corrector magnets for LHC. Originally their design consisted of a coil shrunk inside an aluminium shrink ring and this assembly was centred inside an iron yoke by means of keys [2]. This has now been simplified by packing the coil inside scissor laminations and shrinking the ring around this assembly.

The advantages are that the iron has been brought close to the coil boosting the field, that the shrink ring acts now at a larger diameter and is therefore easier to shrink and also yields more prestress when cooling down, and finally that it suppresses the expensive machining of keyways. Three different corrector magnet prototypes have been built along these lines, one sextupole and one decapole by INDAR, Spain and another sextupole by CERN [3]. Table 1 gives their data. All three magnets use the same enamelled wire of 0.61x1.13 mm metal cross section. The first magnets are made with vacuum

Table:1 Parameters of the three corrector prototypes

		INDAR	INDAR	CERN
		sextupole	decapole	sextupole
Magnet length	mm	150	75	150
Nominal current	A	383	346	625
Critical current (4.2K)	A	765	945	993
Peak field at crit. current	T	5.2	4.1	3.26
Coil Inner Diameter	mm	56	56	56
Coil Outer Diameter	mm	76	71	61
Number of radial layers		8	6	2
Number of turns		112	48	26
Diameter over coil insulation	mm	85.5	78.9	66
Scissors lamination I.D.	mm	86.2	79	66.2
Scissors lamination O.D.	mm	115	95.2	89.4
Largest lamination width	mm	15.1	8.4	12
Smallest lamination width	mm	13.8	7.7	11.2
Lamination thickness	mm	0.8	0.8	1
Angular orientations		6	10	6
Aluminium shrink ring I.D.	mm	115.4	95.3	89.88
Shrink ring O.D.	mm	140	106	100
Interference on diameter	mm	0.35	0.44	0.12
Average coil prestress 300 K	MPa	-35	-46	-31
Expected coil prestress 4.2 K	MPa	-53	-47	-36

impregnated coils with several layers of winding. The last magnet is made in a more economical way with a two layer coil wetted with epoxy during the winding process. The coils of these magnets are protected with a 2.5 mm thick glassfiber epoxy layer and the laminations are directly placed around this insulation layer. The design did not incorporate a mechanical stop. The sextupole laminations were oriented in six angular positions. The circular coil assembly will therefore undergo a slightly hexagonal deformation calculated to be less than 0.03 mm on the radius which does not compromise the field quality. If necessary one can shape the opening in the laminations to compensate for this. This also implies that the coil is supported by local forces repeated every 4.8 or 6 mm over the length. The decapole laminations were oriented in ten angular positions and the pitch between every tenth lamination is 8 mm. The magnets have been trained at 4.2 K at CEDEX in Spain (Fig. 2). The first two magnets show one quench at or below the nominal current. The third magnet starts training well above the nominal current. Compared to other designs with similar impregnated coils the training is similar or even quicker which shows that the laminations give a successful pre-stress.

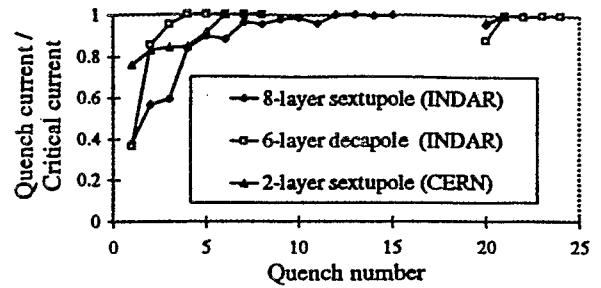


Fig: 2 Training and re-training of the three corrector prototypes

4 MULTI APERTURE MODEL

The method can readily be applied to magnets with a multitude of apertures. A dummy assembly of a twin aperture model has been built and tested to measure the application to multi aperture magnets (Fig.3). In contrast to the single aperture case, one cannot turn the scissor lamination to any desired orientation. However, if the multiple apertures are symmetric in the X and the Y planes a single type of lamination can be used in four different positions: a starting position, an opposing position obtained by turning the plate over 180 degrees, and a third and fourth position by turning the first two plates upside-down. The angle of eccentricity can be chosen to 45 degrees as we did in this model for equal prestress in all four directions or it can be chosen differently to obtain a dominant pre-stress in one particular direction. The plate thickness was 5 mm and aluminium dummy coils were used as well as an aluminium shrink ring. The stresses as found from diameter measurements show (Fig.4) that the coil compression increases strongly during the cooldown and the alignment of the plates occurs at about 130 K. Beyond this temperature the shrinkage of the shrink ring is not acting on the coil anymore and the coil is slightly

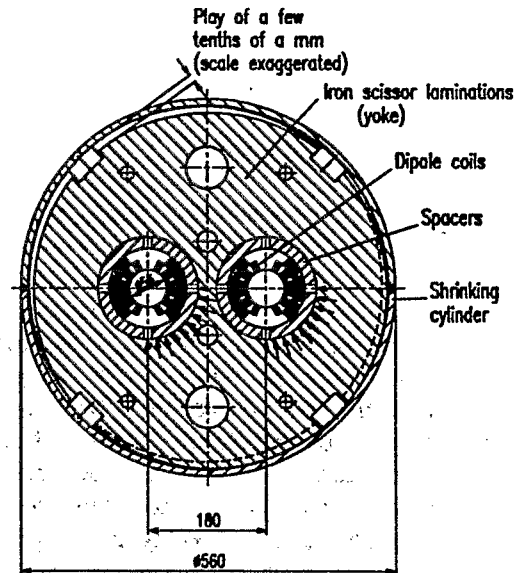


Fig:3 Mechanical model of twin aperture magnet

de-compressing. Several heat cycles showed that the results are repeatable. The pressure on the coil before and after a heat cycle are practically identical ($\pm 10\%$) which means that frictional effects are very small. The model has several times successfully been disassembled and re-assembled, once using a real coil.

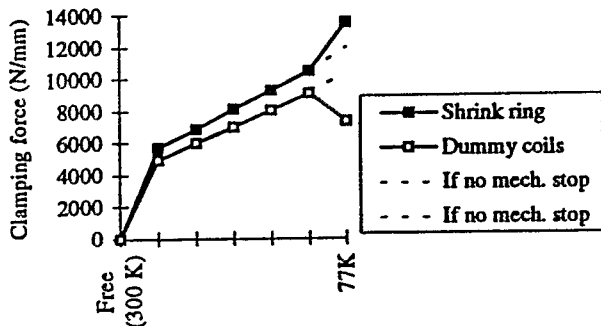


Fig:4 Clamping force in the twin aperture model (corresponding to coil stress of 40 MPa at 300K and 60 MPa at 77K for chosen interference)

5 TWIN APERTURE DESIGN

As a possible application to a multi aperture yoke a tentative design has been made for the LHC main dipole. Two questions turned up: Can one slip the laminations over the not yet collared coil which has dimensions a few millimetres larger than the compressed dimensions and how close can one bring the iron to the coil without losing the field homogeneity due to saturation of the iron at the very high field of 8.4 Tesla? To satisfy these demands we use the fact that each lamination only needs to touch the coil over the 90 degree angle where the pressure is exerted. The shape of the rest of the hole in the lamination can be at a larger radius to obtain the play necessary for the assembly. Magnetic calculations showed that the multipoles caused by local saturation in the iron can be reduced to less than 10^{-4} of the main field over the whole range of excitation by shaping the hole of the iron lamination elliptical and by adding holes to the lamination in the correct positions (Fig.5). The calculations showed that the field homogeneity is not influenced by the iron quality at least less than in the present design. A cost estimation showed that the suppression of the usual collars may lead to an interesting cost saving on such a magnet. The assembly can be made by securing packages of laminations in the open position using keys, slip the packages over the coils, take the keys out, pre-compress the coils moderately under a press and maintain the compressed position introducing corresponding keys. The skin can then be welded around the magnet and the keyways are such that they accept the additional compression during cooldown loosening the keys.

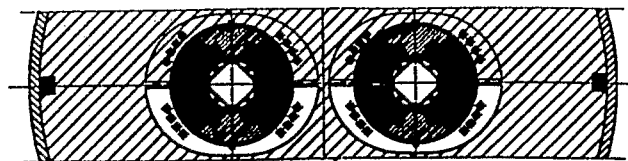


Fig:5 Elliptical lamination holes for field homogeneity

6 CONCLUSIONS

It has been shown that the scissor laminations form a practical and elegant way to transmit the compressive force from the shrink ring to the coil body. The yoke is built with a single type of lamination and does not need a subdivision in segments. The system allows to choose freely the desired direction of the prestress on the coil by orienting the laminations accordingly. It can replace the usual collars leading to non-negligible savings. It permits to bring the iron close to the coil enhancing the magnetic field and allowing to reduce the outer diameter of the yoke. Finally it allows to "collar" the coil at a relatively low compression because an additional compression can be obtained from the differential contractions during the cooldown thus avoiding damage to the insulation by over-stress at room temperature and reducing the size of the necessary press. The system has been tested on mechanical models and has already successfully been applied to several types of corrector magnets.

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Construction and Tests of a Model of the LHC Superconducting Corrector Magnet MDSBV

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Abstract - A full-scale model of the 1.25 m long MDSBV (Magnet Decapole Sextupole Bending Vertical) correction magnet for the Large Hadron Collider (LHC) [1] has been constructed and is currently being tested. The model contains the desired dipole and sextupole but not the decapole which was decided upon later. The magnet was built in a very compact way by placing the dipole coil around the sextupole coil. The two coils were vacuum impregnated and prestressed by shrink-fitted aluminium rings. The design took into account the high positional accuracy requirements for the coils and incorporated manufacturing techniques which are compatible with mass production methods, as approximately 800 of these magnets will be required for the LHC. The model is being tested in liquid helium at the temperature of 4.2 K and will be tested later at 2.0 K. The paper describes the construction, the experience gained during assembly, the test conditions and gives the first test results.

I. INTRODUCTION

The LHC lattice foresees two correction magnets in the standard half cell of each of the two rings, a first one called MDSBH (Magnet Decapole Sextupole Bending Horizontal) or MDSBV (*idem* Bending Vertical) is placed close to the quadrupoles and a second one, MDOS (Magnet Decapole Octupole Sextupole) in the centre of the half cell between the dipoles. The MDSBH and the MDSBV contain respectively a horizontal and a vertical 1.5 T dipole coil for orbit correction. In addition, both contain a 4500 T/m² sextupole coil to compensate the sextupole errors from the main dipole magnets and to correct the lattice chromaticity, and a 3.2×10^7 T/m⁴ decapole coil to correct the decapole errors from the main dipole magnets. The other corrector magnet (MDOS) will be a combination of a sextupole of 1600 T/m², an octupole of 1.8×10^5 T/m³ and a decapole of 1.4×10^8 T/m⁴; the design of which is discussed elsewhere [2].

The MDSBV model has been designed with a sextupole of 8000 T/m² and does not yet contain the decapole. Its cross-section is shown in Fig. 1. In addition to the coils and the shrink-fitted rings one sees the system of keys which centres the coil assembly in the solid iron yoke and absorbs the differential contraction between the two components.

II. CONSTRUCTION OF THE COILS

The design parameters of the sextupole and the dipole coils are given in Table 1.

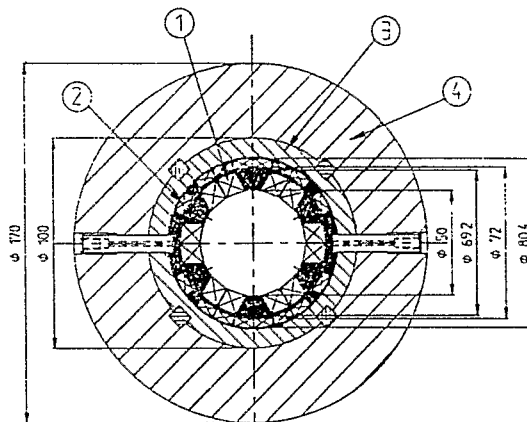


Fig. 1. Cross-section of MDSBV model magnet
 1. Sextupole coil, 2. Dipole coil,

The sextupole coil was wound with a superconducting wire of rectangular cross-section, around a copper central island which had previously been insulated with 0.2 mm of G-10. End spacers were put in after every second layer to dilute the end field and to absorb the space taken by the joggles. It was then impregnated with epoxy resin in a mould. From the tests carried out on trial coils it appeared that the PVA insulation on the wire expanded during the heating of the mould thereby blocking the interturn spaces and preventing the resin from fully penetrating

TABLE 1
 COIL PARAMETERS OF MDSBV MODEL MAGNET

	Sextupole	Dipole
Wire dimension (metal)	0.58 mm x 1.08 mm	Diam. 0.29 mm
Wire dimension (overall)	0.70 mm x 1.20 mm	Diam. 0.35 mm
Copper/NbTi ratio	1.6	4.5
Insulation (enamel)	0.06 mm PVA	0.03 mm PVA
Nominal current	458 A	47 A
% of short sample	61	83
Number turns/coil	104	1332
Stored energy	6.2 kJ	5 kJ
Inductance	59 mH	5.7 H
Length straight sectn.	1 m	1 m

the coil. The impregnation and curing of the actual coils was therefore carried out at a lower temperature to avoid this problem.

The coil wire resistance was measured and the ground insulation tested, before and after winding the coil and also after impregnating with resin.

The dipole coil was wound with a superconducting wire of round cross-section (diam. 0.29 mm), again around a copper central island. Winding trials showed that it was difficult to wind such a thin wire in a regular pattern. The wire was therefore prefabricated, on a special machine, into a twelve parallel wire ribbon, bonded together with a meniscus of epoxy resin between each wire. To ensure that subsequent turns did not nest in the grooves of the previous turn, the ribbon was flattened between rollers reducing the overall thickness from 0.35 mm to 0.32 mm. This produced flats on the wires and ensured that the turns would stack in a regular manner. As the insulation on the ribbon was very thin, it was important to carry out the work in a clean area to avoid dirt particles getting between the layers and puncturing the insulation.

The coil basically consists of two blocks, separated by a copper spacer. Three other sub-blocks had to be created to allow the coil to follow the curvature of the cylinder. These were created by wedge-shaped longitudinal spacers. The latter and the end spacers were fabricated from filled epoxy resin.

The series connections of the twelve wires in the ribbon were made after the coil had been impregnated with epoxy resin. This had the advantage of being able to carry out inter-wire insulation tests at all stages during the fabrication.

The connections were made by fanning out the ribbon at the ends into individual wires and removing the PVA insulation. Connecting wires were then twisted together and soft-soldered. Each joint was then soldered to a printed circuit board to provide mechanical support and additional heatsink by the soldered tracks. Table 2 gives the resistances of different connection samples, measured at 4.2 K. Fig. 2 show the execution of the series connections of the dipole coil pancakes.

TABLE 2
JOINT RESISTANCES OF PAIRS OF DIPOLE WIRES

Measured results for a joint length of 10 mm. (Only one sample of each combination was measured, therefore can only be used as guideline figures)

Method of bonding	Joint layout		
	Twisted	Twisted and wire-wrapped with Cu wire	Twisted and wire-wrapped with s.c. wire
Pre-tinned wire, soldered	$3.6 \cdot 10^{-7} \Omega$	$2.6 \cdot 10^{-7} \Omega$	$1.0 \cdot 10^{-7} \Omega$
Bare wire, soldered	$1.2 \cdot 10^{-7} \Omega$	$0.5 \cdot 10^{-7} \Omega$	$0.6 \cdot 10^{-7} \Omega$
Bare wire, not soldered	$2.2 \cdot 10^{-4} \Omega$	$2.0 \cdot 10^{-4} \Omega$	$6.6 \cdot 10^{-4} \Omega$

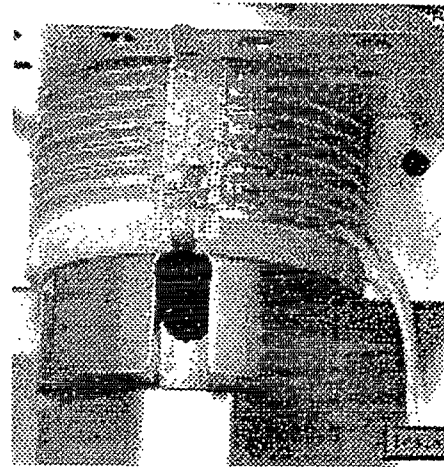


Fig. 2. The soldered joints of the dipole pancakes

III. ASSEMBLY OF THE COILS

The sextupole coils were assembled on a retractable mandrel using dowel pins. The inter spaces between the coils were filled with glass fibre and, at seven places over the length of the coil, bands of pre-impregnated glass fibre were wound around the assembly and cured at room temperature. The bands were clamped in accurately located mould tools such that after curing the outer surface of the bands provided a precise diameter onto which the dipole coils were located concentrically with the sextupole coils within the required tolerance. The bands also guaranteed the angular positions of the sextupole coils such that the dowel pins could be withdrawn. This was necessary to allow the final extraction of the central mandrel. The series connections were then made between the sextupole coils by soldering the connecting pairs of wires into u-shaped copper channels, the latter forming an additional heatsink.

The sextupole coil assembly was finished by filling the spaces between the bands with glass fibre tape which was then filled with resin during the final impregnation of the complete coil assembly.

The dipole coils were assembled on the outside diameter of the sextupole and banded in the same way. Their angular position relative to the sextupole coils was maintained with tooling mounted on the two ends. Once these bands were cured the tooling was removed, the series connections were made, and the interspaces between the bands filled with glass fibre tape. Finally the whole assembly was impregnated in a mould. The diameter of the mould had been made such that the coil diameter would be precise enough for the fitting of the aluminium shrink rings without any machining of the coil assembly. It appeared to be necessary, however, to polish the outside of the coil assembly as the roughness of the surface was not acceptable for the shrink fitting process. (The precision obtained during assembly is typically ± 1.5 milliradian for the angles of the pole mating faces and ± 0.05 mm on the radii.)

The shrink fit was done after measuring the diameter of the impregnated coil assembly and machining the aluminium shrink rings to the desired interference fit. Table 3 gives the measured fit at room temperature as well as at 77 K and also lists the deduced pre-stress and contraction coefficients at liquid nitrogen temperature.

Finally, the keyways were machined in the aluminium rings to centre the coil assembly in the yoke. The latter was planned to be made of a simple thick-walled iron tube. However, it was not possible to machine the keyways over the one-metre length within the desired precision of 0.05 mm and therefore the yoke had to be subdivided into short lengths. For a series production a laminated yoke made with precise stampings will therefore be more appropriate.

The power connections to the current leads were made on the periphery of the magnet by clamping the lead against the coil wire (Fig. 3).

The input and output wires of the sextupole coil were brought out of the coil assembly up to the clamps on the periphery and were backed up with a copper strip. The input and output wires of the dipole were first soldered to a thicker wire of the type used for the sextupole coil and then brought from the coil to the clamps as in the case of the sextupole. Shunt resistors, to protect the magnet against an accidental open circuit, were mounted on the outside of the yoke and were directly connected to the coil side of the connection clamps.

IV. THE TESTS

A first test has been carried out at Rutherford Laboratory at a temperature of 4.2 K. This test, where the magnet was suspended in a vertical cryostat, included the measurement of the magnetic field (by means of Hall probes), the voltages on each individual sextupole coil and the voltages on each dipole pancake. As can be seen on the load lines (Fig.4 and Fig. 5), the sextupole was expected to reach its nominal field even at this temperature but the dipole was not. The first part of the test consisted of running up the sextupole magnet and the dipole magnet separately. In the

TABLE 3
PRE-STRESS FROM SHRINK FIT AND COOLDOWN

Shrink fit	
O.diam. impregnated coil (R.T.)	82.482 ± 0.009 mm
I.diam. aluminium rings	82.360 ± 0.010 mm
Idem after shrink fit	82.422 ± 0.010 mm
Pressure between ring and coil	9.8 N/mm ²
Cooldown to 77 K	
I.diam. aluminium rings	82.103 mm
I.diam. free aluminium ring	82.066 mm
Thermal contraction aluminium	3.6 × 10 ⁻³ mm/mm
Thermal contraction coil	4.1 × 10 ⁻³ mm/mm
Remaining pressure between ring and coil	6.8 N/mm ²

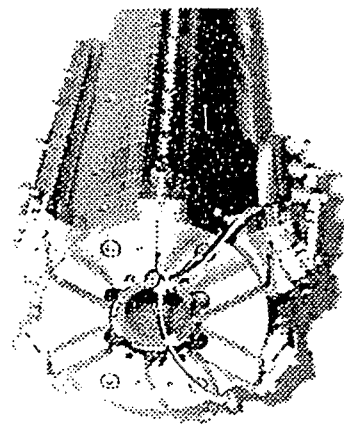


Fig. 3. The power connections of the magnet

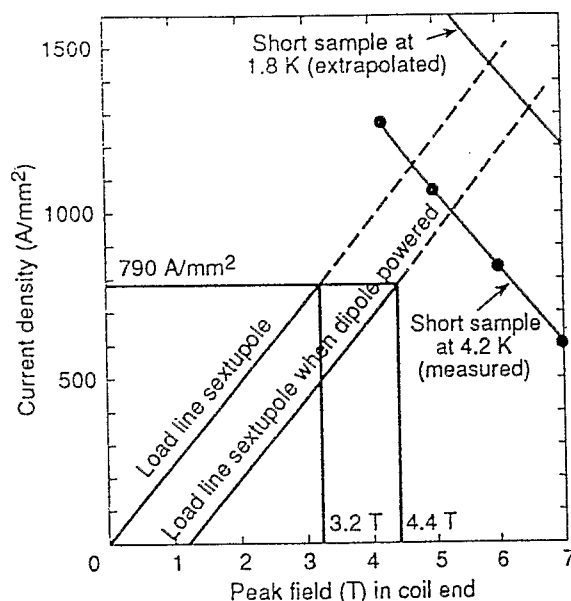


Fig. 4. The load lines of the sextupole. (Conductor parameters given in Table 1)

second part one magnet was powered and the other run up to see the limitations introduced by the coexistence of the two magnetic fields. This second part was repeated after changing the polarity of one of the magnets. In this case the field distribution changes to its mirror image and the peak fields will therefore be found in locations symmetrical to those in the previous condition. The dipole performance is shown in Fig. 6. In stand-alone operation, it trained from 44% to 80% of short sample current (corresponding to a field of 1.44 T) in 32 quenches and was still improving. With an additional sextupolar field, the peak field in the dipole coil increases and occurs in another location. Therefore it trained again but more rapidly. After changing the polarity, it seemed to continue the same training curve. The sextupole performance is shown in Fig. 7. In seven quenches it trained from 52% to 74% of short sample current (corresponding to a gradient of 8500 T/m²). The training was stopped at this point to avoid overstressing

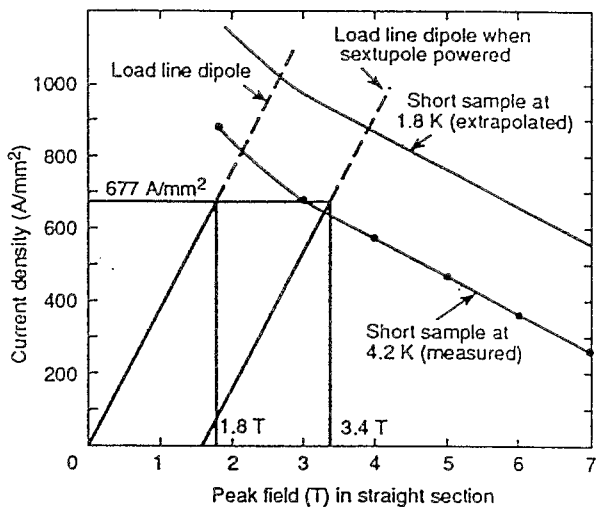


Fig. 5. The load lines of the dipole. (Conductor parameters given in Table 1)

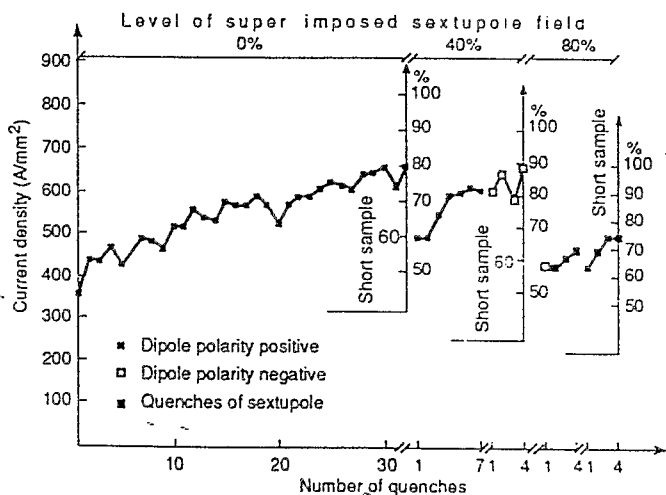


Fig. 6. Training of dipole at 4.2 K

the coil as its design gradient was 8000 T/m². An additional dipole field caused a new training to approach a similar point as compared to short sample current. A second test will follow at CERN and will include the measurement of the field harmonics in the straight section and in both ends. Although this test will finally be done at the LHC operating temperature of 1.8 K, the magnet will first be tested at 4.2 K to see whether it will directly reach the present performance again or whether it will re-train.

VI. CONCLUSIONS

The MDSBV magnet model, built by "nesting" the sextupole coil inside the dipole coil, has been constructed and successfully tested at 4.2 K. The dipole, designed for a field of 1.5 T at a temperature of 1.8 K was brought to 1.44 T in stand-alone mode and to 1.08 T when the sextupole was powered at 80% of nominal field. The sextupole, designed for 8000 T/m² at 1.8 K, was brought

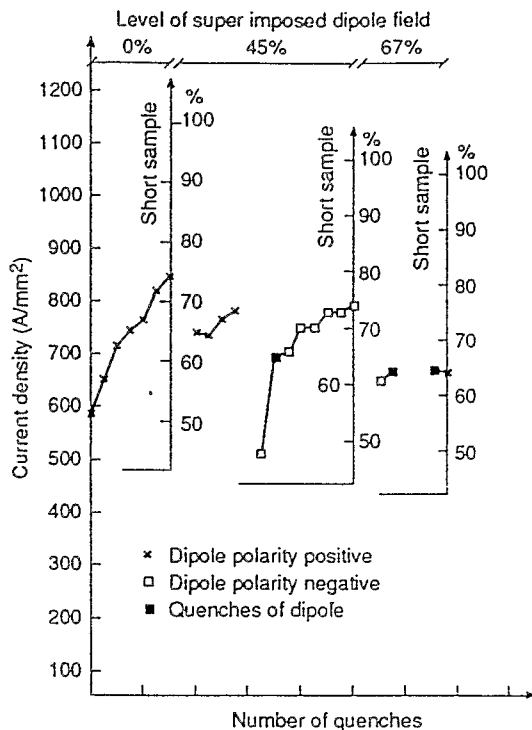


Fig. 7 Training of sextupole at 4.2 K

to 8500 T/m² in stand-alone operation and to 6800 T/m² when the dipole was powered to 67% of its nominal field. The dipole showed more training than the sextupole, maybe due to the large number of turns (1332) per pole. The result seems to confirm that the magnets, after training, may run at 70% of short sample current. However, a second test at 4.2 K is necessary to check whether the magnets will re-train. A final test at the LHC operating temperature of 1.8 K is planned in the near future.

ACKNOWLEDGMENTS

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AUTOMATED DESIGN OF A CORRECTION DIPOLE MAGNET FOR LHC

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ABSTRACT

A correction dipole magnet, with a horizontal dipole nested inside a vertical dipole has been designed and optimized linking together different electromagnetic software and CAD/CAM systems. The necessary interfaces have recently been established in the program ROXIE [1] which has been developed at CERN for the automatic generation and optimization of superconducting coil geometries. The program provides in addition to a mathematical optimization chest, interfaces to commercial electromagnetic and structural software packages, CAD/CAM and databases. The results from electromagnetic calculations with different programs have been compared. Some modeling considerations to reduce the computation time are also given.

1 INTRODUCTION

The low- β dipole corrector, MCBX, is a single-bore, 0.6 m long magnet, whose main parameters are given in Table 1. It features two nested single-layer dipole coils, the inner coil yielding a vertical dipole field and the outer coil a horizontal dipole field. The coils are individually powered with 600 A power supplies. The coils are wound from a NbTi rectangular wire bound together as a flat cable of 9 or 7 wires in the inner and

Table 1: Parameters of the LHC low- β corrector dipole prototype, MCBX.

		Inner/Outer Coil
Operating dipole field	[T]	3.3 / 3.3
Integrated dipole field	[Tm]	1.341 / 1.273
Peak field in the cond.	[T]	4.431 / 4.749
Margin on the load line		50.51% / 45.54%
Operating current	[A]	511 / 599
Magnetic length	[m]	0.41 / 0.38
Overall length	[m]	0.6
Coil inner diameter	[mm]	90 / 123.7
Coil outer diameter	[mm]	119.7 / 146.8
Yoke inner diameter	[mm]	100
Yoke outer diameter	[mm]	450
Wire dimension (ins.)	[mm]	1.65 x 0.97
Wire dimension (metal)	[mm]	1.53 x 0.97
Cu/Sc-ratio		1.6
Operating temperature	[K]	1.9
No. of turns per coil		414 / 406
Stored energy	[kJ]	21.162 / 29.725
Self inductance	[mH]	162 / 166

the outer coil respectively, and cooled at 1.9 K. The design field integral is 1 Tm in any direction for an excitation current ranging from 360 to 600 A. Due to the short length of the magnet, the end fields contribute more than 50 % to the field integral. Therefore, an optimization in 2D and scaling with the magnetic length is not sufficient. Particular attention has to be paid to the lead end where the transitions from one block to another are made in addition to the leads entering over the top of the end blocks. Figure 1 illustrates the approach to an integrated automated design comprising the following 8 steps: 2D coil design including mathematical coil optimization, 3D coil design with mathematical coil optimization, transfer of model file to Opera-3D[®][3] for calculations including the iron saturation, transfer of file to CAD for the mechanical drawings, transfer of file to the CNC-machine for machining of the end spacers.

2 ELECTROMAGNETIC CALCULATIONS

2.1 Coil Optimization in 2D

Based on some design constraints, which includes the space limitations and the operating current, a preliminary conductor and coil lay-out can be chosen. With only a few lines of input data the cross-section is generated by means of the ROXIE-program. ROXIE

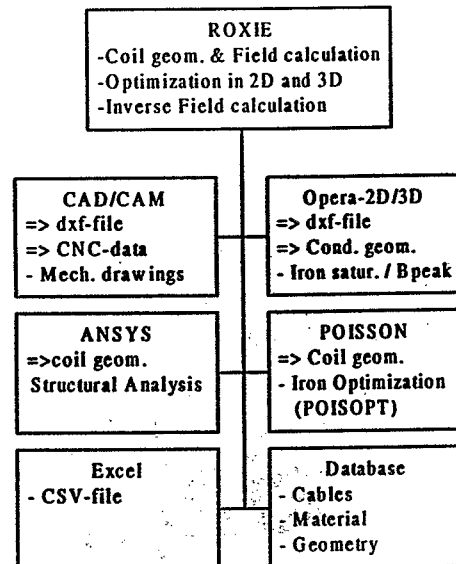


Figure : 1. Approach to an integrated automated magnet design

Table 2 : MCBX-magnet, multipole content in 2D.
(bi = Bi/B1)

	ROXIE-2D linear iron	Opera-2D [®] linear iron	Opera-2D [®] non-linear iron
B1 [T]	3.2989	3.3001	3.2972
b3 x 10 ⁴	0.0007	0.0250	-0.0446
b5 x 10 ⁴	0.0580	0.0565	0.0557
b7 x 10 ⁴	0.0143	0.0141	0.0141
b9 x 10 ⁴	0.0001	0.0001	0.0001

applies Biot Savart's law on line currents and the iron with linear or infinite permeability is taken into account by imaging. After a few manual iterations an automatic design optimization can be carried out. ROXIE includes many different optimization algorithms [2]. All the design parameters can be addressed as design variables for the optimization. One of the most robust algorithm is EXTREM, which was used for the optimization of the LHC low- β dipole, MCBX, described here. The objectives were low peak field in the conductor and minimized multipole content.

2.2 Optimization of the iron circuit

Once the coil cross-section is optimized the geometry is exported in DXF-format to commercial FE-software. The Opera-2D[®] package is used to calculate the saturation effects in the iron circuit. Different inner and outer radii were investigated for variation of the harmonic content, b3 and b5 in particular, at current levels ranging from injection to 120% of the nominal current. Table 2 compares the calculated field values from ROXIE with those from Opera-2D[®] linear and non-linear models. The main difference is the b3 due to the simplified conductor blocks in the FE-model. In total 23000 isoparametric, 6-node-triangles were used. The

Table 3 : Integrated multipole content in 3D
(bi = Bi/B1)

	ROXIE-3D linear iron	Opera-3D [®] linear iron	Opera-3D [®] non-linear iron
B1 [Tm]	1.3627	1.3421	1.3423
b3 x 10 ⁴	-0.3642	1.9654	1.9639
b5 x 10 ⁴	0.1700	0.1775	0.1775
b7 x 10 ⁴	0.0372	0.0381	0.0381
b9 x 10 ⁴	0.0013	0.0014	0.0014

harmonic content was evaluated at a 30 mm radius and scaled down to 10 mm radius.

2.3 Coil Optimization in 3D

Most of the quenches occur in the coil ends. Therefore, particular attention is paid to the optimization of the coil end geometry with its constant perimeter ends. In ROXIE the geometry is created with a few additional parameters: big half axis of the ellipse, angle of the cable in the yz-plane, and the axial shift in the z-direction. Before the electromagnetic optimization the end geometry is optimized to maximize the minimum radius of curvature. The user can choose from a normal ellipse or a hyper-ellipse form. Coil blocks can be aligned on the winding mandrel or on the outer radius with so-called shoes. The connection end with transitions from one block to another can be created with an asymmetric model. The leads entering from outside are modeled as 8-node-bricks from pre-defined cut planes. The field calculation in 3D cannot take into account the iron except for the integrated multipole content, when the calculation is carried out from the symmetry plane far enough along the z-axis. The field optimization is done by varying the axial position of the blocks.

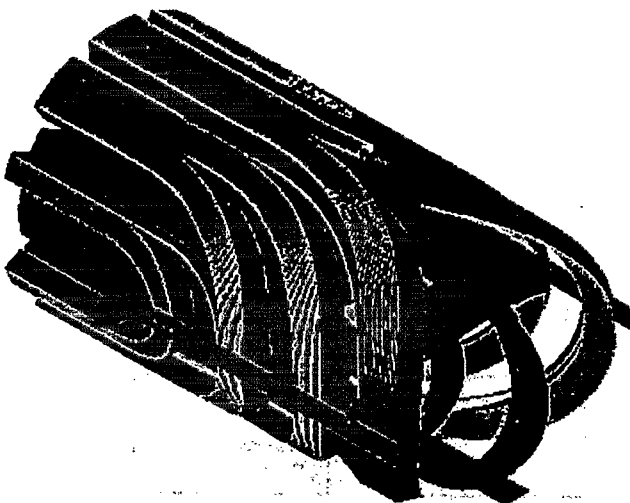


Figure : 2. MCBX-magnet lead end from ROXIE. Note the modelling of cross-over conductors

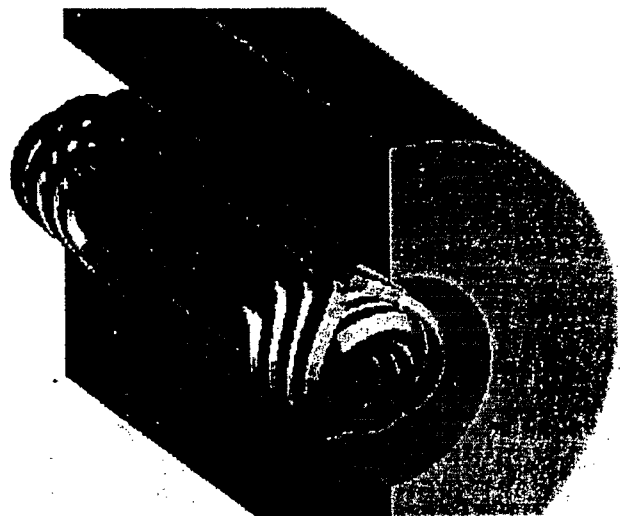


Figure : 3. MCBX-magnet Opera-3D[®] model. Only the return ends modelled for simplicity.

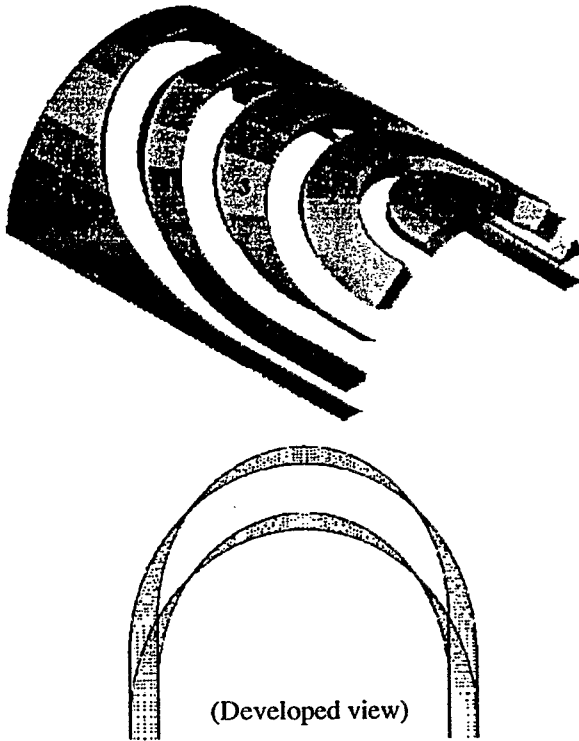


Figure : 4. End-spacers and the 3D-polygons for CNC-machining of spacer no. 3.

2.4 FE-analysis in 3D

The influence of the non-linear iron on the field quality and the peak fields is studied with Opera-3D[®]. The optimized conductor geometry can be imported from ROXIE as 8-node-bricks or as 20-node-bricks. This speeds up considerably the modelling time, since these geometries cannot be created in Opera using the built-in constant-perimeter-end coil primitive. Reduced scalar potential was used for all the regions inside the iron. Table 3 compares the integrated multipole content of ROXIE with linear iron ($\mu_r=2000$) to that of the Opera-3D[®] model. A set of dummy conductors, which are duplicates of the 'active' ones without imaged parts, and with zero current, were modelled to speed up the peak-field calculation. Only the part of the geometry, which was solved in Tosca[®] had to be activated. A considerable enhancement of the peak field was found in 3D: 4.75 T with respect to 4.23 T in 2D.

3 MECHANICAL DESIGN

3.1 Structural analysis

Microscopic movements of the conductors or micro-cracking of epoxy can dissipate locally enough energy to cause a super conducting magnet to quench. Therefore a sufficient coil pre-compression is essential for the magnet performance. Too high pre-stress, however, can damage the insulation or cause copper to yield. The

ROXIE conductor geometry was imported to ANSYS[®] [5]. An electromagnetic model was first created to determine the Lorenz forces, which were then transferred to the mechanical model. Five load steps were considered: assembly at 293 K, cooldown to 1.9 K, inner or outer coil powered, and the combined field.

3.2 CAD-CAM

The coil cross-section, developed view (sz-plane) and cut through the end (yz-plane) can be exported from ROXIE to most CAD programs in DXF format. The end spacer geometry is defined by 9 polygons, presented in global xyz-coordinates, over the radius and can be imported to CAD and CAM packages. These rather complicated shapes are then machined with a 5-axis CNC machine.

3.3 Mechanical tolerances

The sensitivity analysis for the mechanical tolerances is carried out in ROXIE by giving the tolerances as upper and lower bounds for the number of geometrical design variables. The upper and lower quotient per unit displacement are then calculated, and the Jacobian error matrix can be taken into a spreadsheet program for further processing.

4 CONCLUSIONS

A number of tools to speed up the design process and decision making have been established in the ROXIE program as an approach towards an integrated design of superconducting magnets. The optimized conductor geometry can be directly taken into commercially available FE-programs. Links to CAD/CAM-packages allow creation and changes of the mechanical drawings to be done in parallel with the design optimization. Short correction magnets, which in addition to optimization require studying of the manufacturing tolerances in 3D, can be efficiently designed by these means.

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PROGRESS IN THE DEVELOPMENT OF THE 1-m MODEL OF THE 70 mm APERTURE QUADRUPOLE FOR THE LHC LOW- β INSERTIONS

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Abstract

Within the LHC magnet development program Oxford Instruments has built a one metre model of the 70 mm aperture low- β quadrupole. The magnet features a four layer coil wound from two 8.2 mm wide graded NbTi cables, and is designed for 250 T/m at 1.9 K. The magnet has previously been tested between 4.5 K and 2.3 K. In this paper we review the magnet rebuild and the subsequent tests. Results on magnet training at 4.3 K and 1.9 K are presented along with the results related to quench protection studies.

1 INTRODUCTION

As part of the LHC magnet development program Oxford Instruments has built and tested a one metre model of the 70 mm aperture quadrupole for the low- β insertions. The design and construction of this magnet has been reported previously along with the results of the first tests [1, 2]. The magnet [3] has a graded 4-layer coil with the transition between the two cable types in the middle of the 2nd layer. Thin collars 10 mm wide ensure accurate coil location during assembly. Final preload is applied using a set of aluminum collet force rings, whose load is transmitted to the collared coil through the 4-piece iron yoke.

During the first tests in March 1995, the magnet reached a maximum current of 3780 A at 4.5 K and of 4112 A at 2.5 K. After this point the performance became erratic and it was not possible to train further. It was thought that the quenches were caused by conductor motion. The magnet was warmed to room temperature and the coil preload increased by tightening the force rings. The magnet was tested again in August 1995, but repeatedly quenched at currents between 3045 and 3180 A. The quench current was independent of temperature. The quenches were all in layers 3 and 4 of quadrant D, but the instrumentation was insufficient to locate the quenches within this coil pair. The magnet had protection resistors fitted across each pole which prevented further diagnostic testing without removal of the magnet from the test dewar.

In this paper we describe the rebuild of the magnet and present the results of the subsequent tests. The magnet has been energised to a gradient above the LHC operating point and measurements have been made of the quench velocities and peak temperatures.

2 DISASSEMBLY AND REBUILD

With the magnet at room temperature and the support structure disconnected, the protection resistors across each pole were removed and the external joints broken. Measurements of resistance and inductance for each coil indicated the presence of a multi-turn short in quadrant D. After the disassembly of the magnet, the location of the short was determined by applying a current to the coil and measuring the voltage to each turn. A multi-turn short caused by scissoring was found in the ramp between layers 3 and 4. About half the strands of one piece of cable were damaged. In hindsight, this region was probably first damaged during energisation to 4112 A in March 1995.

The size and location of the short precluded any attempts at repair and a replacement coil was wound with modified geometry in the layer ramp region. This modification was also made to all the other coils to prevent similar damage. Due to availability of material, the rewound coil incorporated a modified insulation system which was 9 μm thicker per turn than in the other coils. To compensate for the increased thickness the sizes of the copper wedges at the pole were reduced. The magnet was rebuilt with a modified set of voltage taps on all coils and all the ground plane insulation replaced. The coil prestress at room temperature was reset to the same value as previously. In all other ways the magnet was identical to the first build.

3 MAGNET TESTS

The magnet was equipped with a set of spot heaters located between turns, one in each of the inner three layers of each quadrant. The heaters could be energised individually to trigger a quench, and all were energised once a quench was detected. Bypass resistors of $R = 2.3 \text{ m}\Omega$ were connected across each quadrant of the coil in an effort to provide additional protection. Although complicating analysis, these resistors were refitted, since the previous tests had not produced sufficient data to allow confidence in their removal.

3.1 Training History

The training history is displayed in Fig. 1. This figure also shows the estimated conductor limits (2-4% uncertainty), the calculated gradients (including iron saturation) and the operating gradients G_{op} of two versions of this magnet. The high luminosity low- β insertions utilize a single aperture quadrupole (MQX), operated at $G_{\text{op}} = 225 \text{ T/m}$ and 1.9 K.

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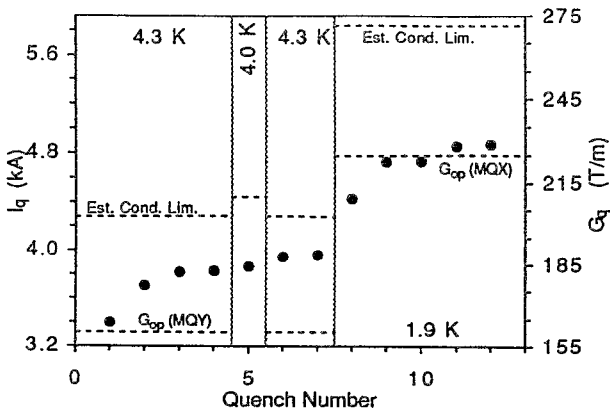


Figure 1: Quench training history.

In the dump insertion, a two-in-one version (MQY) operating at 4.5 K with $G_{op} = 160$ T/m is envisaged.

The magnet was initially trained at 4.3 K and the first quench occurred at 3408 A. This was about 150 A lower than the first quench of the first test. The next three quenches were at 3712 A, 3828 A and 3833 A, slightly above the plateau of the first test. In view of possible conductor limitation at 4.3 K, the magnet was pumped to 4.0 K and energised with a quench at 3860 A. The increase of 30 A was less than the expected gain by cooling and further training was continued at 4.3 K. The following quenches at 3953 A and 3958 A were deemed to represent a relatively stable plateau, and the magnet was cooled to superfluid helium temperature. Throughout this stage quenches occurred in all the coils except in the rebuilt coil D, although coil C (layers 2 and 4) showed a propensity for quenching.

The first quench in superfluid helium occurred at 4441 A. The next two quenches were recorded at 4746 A and 4743 A, with a change in quench location. On the next two runs the current increased to 4862 A and 4879 A. Further training was discontinued due to test equipment failure.

At both test temperatures, the magnet has achieved a plateau above the operating gradients required for the LHC. However, the plateaus are achieved somewhat below the computed conductor limits. The last quench at each temperature was in layer 2, where conductor limited quenches are expected to occur, but the instrumentation is insufficient to determine if these quenches occur at the high field point. The coil pressure is measured to be 74 MPa at zero field and decreases by only 13.3 MPa up to a current of 3750 A, suggesting that the coil remains adequately preloaded to well above the highest quench current achieved. Further testing will be required to verify if the magnet is mechanically limited or if the actual conductor limit is lower than that computed.

3.2 Quench Protection Studies

An important objective of this test was to measure the dependence of the peak conductor temperature T_{peak} and the initial quench velocity v_q on the magnet current I_q , in or-

der to help in the design of a quench protection system for a full-length magnet.

For these studies, special voltage taps were used. They are placed adjacent to the spot heater (HA2), which is located near the lead end of the second layer of quadrant A between turns of the small cross-section cable. By measuring the final resistance of the 125 mm long cable segment which contains HA2, T_{peak} can be determined. Similarly, by measuring the time necessary for the quench front to propagate to another tap 585 mm away v_q can be determined. The quench velocity can be independently measured from the initial resistance growth dR/dt and the measured resistance per unit length of the cable (including magneto-resistance). The measured $\int I^2 dt$ can be converted to peak temperature using the known conductor properties and the magnetic field at the point of the quench.

Figure 2 shows the v_q measured by time of flight (Δt) versus I_q . For spontaneous quenches which do not originate in quadrant A, the quench at HA2 is still triggered by that heater and v_q can be determined. Spontaneous and heater-induced quenches are displayed independently in Fig. 2; they clearly display the same trend. The v_q grows from about 5 m/s at low current to 65 m/s for the highest current 4.3 K quench. At 1.9 K, v_q grows more slowly with current, but quenches at a similar fraction of the conductor limit at the two temperatures have similar velocities.

The quench velocity can also be measured from the initial dR/dt . Comparison is made between the velocities determined by the two methods in Fig. 3. For $v_q \leq 20$ m/s the two methods agree well; however for larger v_q the dR/dt method gives a systematically lower value. Since dR/dt also has a component due to resistivity growth with temperature, the velocity measurements are based only on the first 6-8 ms after the quench enters the cable segment adjacent to HA2. The Δt method, on the other hand, the velocity averaged over the time the quench front propagates between the two taps at the end of the 585 mm segment adjacent to HA2. For two quenches, with Δt velocities of 20 and 45 m/s, v_q ob-

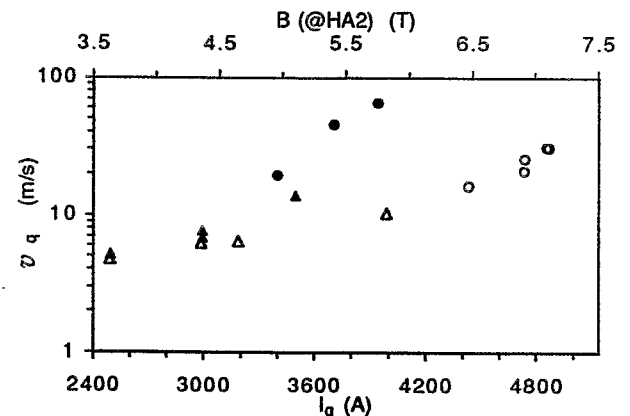


Figure 2: v_q vs. I_q for heater HA2 induced (triangles) and spontaneous (circles) quenches. Quenches at 4.3 (1.9) K are shown as filled (open) symbols.

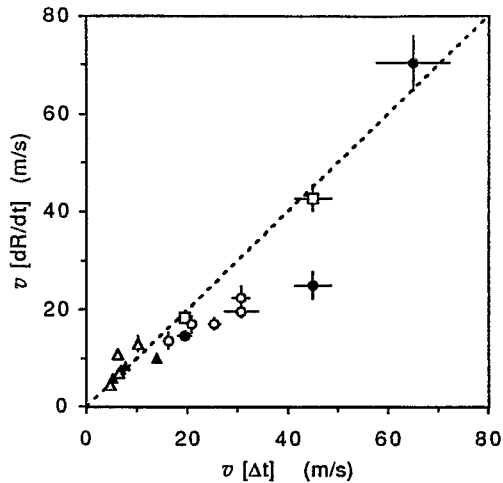


Figure 3: v_q measured by initial dR/dt vs that by Δt for heater HA2 induced (triangles) and natural (circles) quenches. Quenches at 4.3 (1.9) K are shown as filled (open) symbols. The open squares give the v_q measured by average dR/dt for the same quenches plotted directly below them.

tained on the basis of the average dR/dt were calculated, and are shown in Fig. 3. For the highest v_q point Δt is too short to measure initial and average dR/dt independently and the two calculations coincide. The velocities obtained from the average dR/dt agree at both low and high v_q with those from the Δt method. This strongly suggests there is an acceleration of the quench front, at least during the time interval covered by these data.

Figure 4 displays T_{peak} vs. I_q for heater induced quenches at two operating temperatures. At 4.3 K, T_{peak} has a maximum of about 250 K which occurs between 3000 and 3500 A, 70-80% of the estimated conductor limit. T_{peak} is substantially higher at 1.9 K and no maximum is evident up to 4000 A. However, if the highest temperature occurs at 70-80% of the conductor limit, the 4000 A point is close to the maximum. Based on this consideration, it was deemed safe to train the magnet at 1.9 K.

To estimate T_{peak} for spontaneous quenches, $\int I^2 dt$ is computed and displayed in Fig. 5 for the heater induced and the spontaneous quenches. For the heater quenches, the pattern of $\int I^2 dt$ and T_{peak} vs I_q are similar, as expected. The induced and natural quenches follow similar trend in $\int I^2 dt$ vs I_q . At 1.9 K, $\int I^2 dt$ achieves a maximum near 4400 A, about 75% of the conductor limit, and decreases substantially at higher current. The T_{peak} for the highest point is computed to be 465 ± 40 K, where the error includes uncertainties in the start time of the integration and in the value of B at the quench location.

The coil resistance $R(t)$ as a function of time and the energy deposited in the coil $\int I^2 R dt$ were computed. It was found that only about 0.1% of the stored energy was deposited in the bypass resistors, the rest being absorbed by the coil. Since the magnet absorbs its own energy safely, the protection resistors can be removed for subsequent tests.

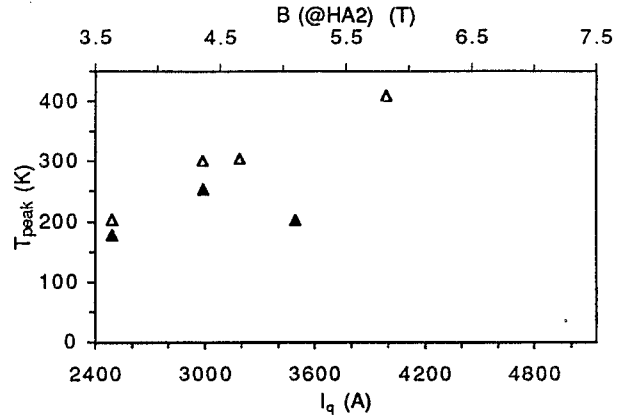


Figure 4: T_{peak} vs. I_q for quenches induced by HA2. Quenches at 4.3 (1.9) K are shown as filled (open) symbols.

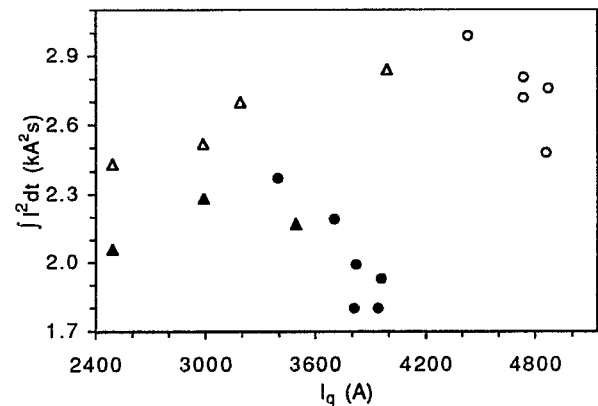


Figure 5: $\int I^2 dt$ vs. I_q for spontaneous (circles) and heater HA2 induced (triangles) quenches. Quenches at 4.3 (1.9) K are shown as filled (open) symbols.

4 CONCLUSIONS AND FUTURE PLANS

The one metre model of the 70 mm aperture quadrupole for the LHC low- β insertions built by Oxford Instruments has achieved the operating current of 4790 A (225 T/m) with only three training quenches in superfluid helium. In comparison with other quadrupoles tested for the LHC, the highest gradient times aperture has been achieved. The mechanical structure of the magnet was shown to be appropriate, and the construction was achieved without the use of expensive tooling. A further test of the magnet is planned in order to measure the field quality, examine the training behaviour following a thermal cycle, and collect additional data relevant for quench protection.

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OPTIMIZATION OF THE LHC INTERACTION REGION WITH RESPECT TO BEAM-INDUCED ENERGY DEPOSITION

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Abstract

Energy deposition in the superconducting magnets by particles from p-p collisions is a significant challenge for the design of the LHC high luminosity insertions. We have studied the dependence of the energy deposition on the apertures and strengths of insertion magnets and on the placement of absorbers in front of and within the quadrupoles. Monte Carlo simulations were made using the code DTUJET to generate 7x7 TeV p-p events and the code MARS to follow hadronic and electromagnetic cascades induced in the insertion components. The 3D geometry and magnetic field descriptions of the LHC-4.1 lattice were used. With a quadrupole coil aperture ≥ 70 mm, absorbers can be placed within the magnet bore which reduce the peak power density, at full luminosity, below 0.5 mW/g, a level that should allow the magnets to operate at their design field. The total heat load can be removed by a cooling system similar to that used in the main magnets.

1 INTRODUCTION

The Large Hadron Collider (LHC) [1] is designed to produce p-p collisions at $\sqrt{s}=14$ TeV and $L = 10^{34} \text{cm}^{-2} \text{s}^{-1}$. The interaction rate of $8 \times 10^8 \text{s}^{-1}$ represents a power of almost 900 W per beam, the large majority of which is directed towards the low- β insertions. Previous studies [2, 3, 4] have identified this as a potentially serious problem. The quadrupole fields sweep the secondary particles into the coils preferentially along the vertical and horizontal planes, giving rise to local peak power density P_{max} as much as an order of magnitude larger than the average. Tests of porous cable insulation systems [5] cooled by a 1.9 K helium bath have shown that for typical cable dimensions up to about 1 mW/g of heat can be removed while keeping the helium inside the cable below 2.2 K. This is the most important point of this study, since too large P_{max} could prevent the low- β quadrupoles from reaching their required gradient.

To minimize P_{max} , we have studied its dependence on the aperture of the front absorber and of absorbers placed inside the magnets, the low- β quadrupole coil diameters, and the beam separation/recombination dipole length. Solutions which satisfy the requirements with a reasonable safety margin have been found.

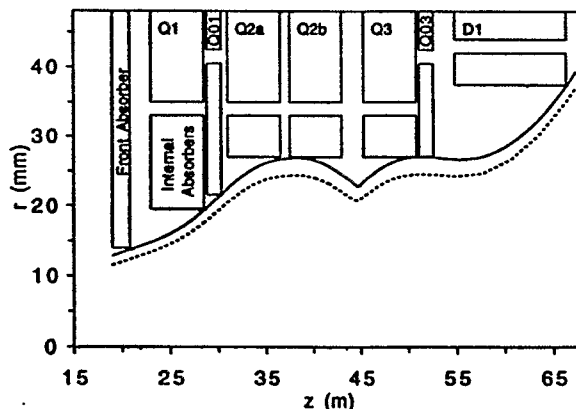


Figure 1: The LHC low- β insertion including absorbers which reach the 10σ limit (injection/collision optics: solid/dashed line) for $d_{coil}=70$ mm quadrupoles.

2 COMPUTER MODELLING

Fig. 1 shows the LHC low- β insertion [1, 6]. The inner triplet is made of four identical high-gradient quadrupoles with coil inner diameter $d_{coil}=70$ mm [7] (Q1 and Q3 focusing and Q2a and Q2b defocusing), which are powered in series and operate at a maximum gradient of 227 T/m at the high luminosity IRs. Two independently powered "trim" quadrupoles of $d_{coil}=85$ mm and maximum operating gradient of 120 T/m (Q01 and Q03) provide the additional strength required by Q1 and Q3 and allow tuning of the triplet. Behind the triplet are the dipoles D1 (single aperture) and D2 (twin aperture). They have $d_{coil}=85$ mm and an operating field of 4.3 T. Their length (11.5 m) is set by the required strength at the combined experimental and injection insertions (points 2 and 8), where space is more limited than at the high luminosity IRs (points 1 and 5).

Alternate IR designs with quadrupoles of $d_{coil}=60$ mm and 80 mm have also been considered. The gradients were scaled with d_{coil} (a little more slowly than $1/r$ for thick shell quadrupoles) and corresponding length changes were made. Optics with minimal perturbations to the baseline were computed for injection (450 GeV, $\beta^*=6$ m) and collision (7 TeV, $\beta^*=0.5$ m) conditions and are summarized in Table 1. Relative to the baseline 70 mm case β_{max} for 80 (60) mm quadrupoles changes by +6.6% (-5.1%) yielding changes in the maximum beam size of +3.3% (-2.6%), considerably less than the change in d_{coil} . Thus increasing the aperture should improve the field quality over the region occupied by the beam and allow more shielding inside the magnet bore,

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Table 1: Characteristics of the IR optics.

d_{coil} L_{mag}	60 mm		70 mm		80 mm	
	5.1 m		5.5 m		6.0 m	
Stage	Coll.	Inj.	Coll.	Inj.	Coll.	Inj.
G (T/m)						
Q1-Q3	251	15.5	227	14.0	202	12.4
Q01	70	4.7	80	5.3	92	6.1
Q03	105	8.5	101	7.7	95	6.1
β_{max} (m)	4204	358	4431	377	4724	402

Table 2: Minimum inner radii of absorbers.

d_{coil} Clearance	60 mm		70 mm		80 mm
	10σ		10σ	8σ	10σ
Collimator	14.0	14.0	12.0	14.0	14.0
Q1	19.0	19.5	17.0	20.0	20.0
Q01	21.5	21.5	18.5	22.0	22.0
Q2-Q03	26.5	27.0	23.5	28.0	28.0
D1	36.5	37.5	34.0	38.0	38.0

while decreasing the aperture will have the opposite effect.

A 1.8 m long copper absorber is placed in front of the triplet and stainless steel absorbers are placed within the magnet bores to minimize the energy deposition in the coils. The LHC design requires [1] that the physical aperture, including effects of dispersion, closed orbit errors, construction and alignment tolerances, and the crossing angle in the IRs, be everywhere at least 10σ (except at the beam cleaning collimators), where σ is the *rms* beam size. Fig. 1 shows the 10σ limit for injection and collision conditions and absorbers with inner radius r_{in} at this limit. The cusp between Q2b and Q3 is where the maximum β changes from one plane to the other. The outer radius of the internal absorbers is 2 mm less than r_{coil} . Table 2 gives r_{in} of the absorbers for the three quadrupole diameters. To allow the effectiveness of the absorber to be evaluated versus thickness for a fixed insertion design, the r_{in} for 8σ are given for the 70 mm quadrupoles. As the r_{coil} grows from 30 to 40 mm the 10σ limits increase by only 0.5 mm (Q1-Q01) and 1.5 mm (Q2-Q03), allowing an increased absorber thickness. However, the D1 absorber decreases from 5.5 to 4 mm thick.

The p-p collisions and showers in the IR components are simulated with the DTUJET93 event generator [8] and the MARS code [9], version 13(96) respectively. Charged particles are tracked through the lattice and the fields within each magnetic element. The cut-off energies are 1 MeV (charged particles), 0.2 MeV (photons) and 0.5 eV (neutrons). Magnet coils are modeled with 4 radial bins of 8.5 mm depth, azimuthal bins varying from 5° at the horizontal and vertical planes to 15° between, and axial bins between 1.1 m (Q1) and 3.8 m (D1) long. The magnet coils, which are a mixture of NbTi, copper, insulation and helium, are simulated as a homogeneous material with $A=50$, $Z=23$ and $\rho=7 \text{ g/cm}^3$. Details such as cooling channels in the yoke and coil ends

are not included. Statistical errors on the Monte Carlo calculation are estimated to be $\pm 15\%$ for P_{max} , $\pm 6\%$ for the energy deposited in each magnet, and $\pm 1\%$ for the total power in the inner triplet, based on comparison of results from different runs with independent random seeds.

3 RESULTS

Fig. 2 shows P_{max} vs z for the IR with 70 mm quadrupoles and no internal absorbers, absorbers at 10σ and 8σ , and all quadrupole absorbers of a uniform radius at the 10σ limit in Q2-Q03. The front absorber aperture is set at 10σ for the case of no internal absorbers. With no internal absorber $P_{\text{max}} = 1.2 \pm 0.2 \text{ mW/g}$, at or above the allowable limit. With individually sized 10σ absorbers the peak is a factor of 3 smaller, giving a reasonable safety margin. Use of an 8σ absorber reduces P_{max} in Q1, but there is little overall improvement. Increasing r_{in} of the absorbers in Q1-Q01 to match the other quadrupoles results in a 25% increase in P_{max} . However, this increase may not be statistically significant and further study will be required to determine if it is necessary to use different absorbers in Q1 and Q01 than in the rest of the triplet.

An unacceptably large P_{max} is observed at the back of D1 even with a 10σ absorber. However, at the high luminosity IRs it is possible to move the outer dipole D2 up to an additional 90 m farther from D1. This would reduce the length of D1 to one-third its present value, corresponding to the first bin in Fig. 2, which has an acceptable power density. Alternatively the integrated strength would be low enough to allow use of conventional magnets for D1 and D2, eliminating the problem altogether.

The cases of three quadrupole diameters with 10σ absorbers are compared in Fig. 3. P_{max} is 40% larger (30% smaller) with 60 mm (80 mm) quadrupoles than the baseline 70 mm case. The reduced margin with $d_{\text{coil}}=60$ mm makes this option unattractive. The 80 mm case has a significantly larger margin, which could be used, if required, to provide additional physical aperture. The larger β_{max} is unlikely to be a problem since the field quality in the region occupied by the beam would be better with a larger aperture magnet.

Shown also is the case in which all quadrupoles, including Q01 and Q03, have the same $d_{\text{coil}}=70$ mm and the absorbers have a uniform r_{in} . P_{max} is 80% larger than the case with 85 mm trims and individually sized absorbers (Fig. 3) and 45% larger than with uniform absorbers (Fig. 2). Apparently it is unacceptable to have a continuous annular gap between the absorber outer and coil inner radii.

Table 3 summarizes the total power deposited in the magnets and internal absorbers. The quadrupoles and the dipole D1 are considered separately, since the actual dipole configuration will probably be different than that considered here. There is little difference among the cases with internal absorbers. Up to half the power is deposited in the absorbers, and it is tempting to consider cooling them at a higher temperature. However, the insulating space between the absorber and the vacuum pipe would reduce the absorber

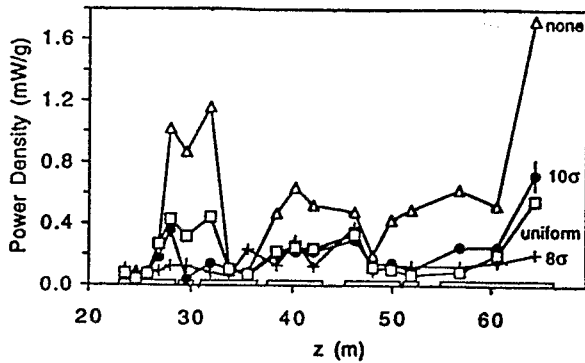


Figure 2: P_{max} vs z for 70 mm quadrupoles with several absorber configurations.

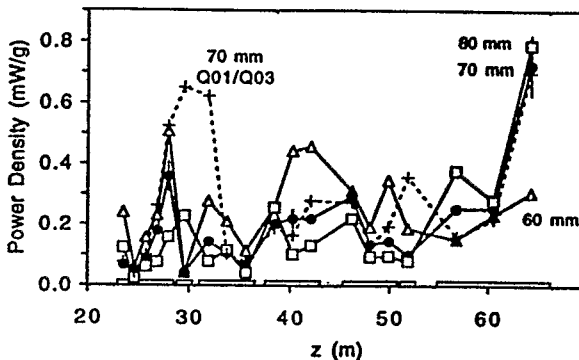


Figure 3: P_{max} vs z for main quadrupoles of three d_{coil} with 85 mm trim quadrupoles, and 70 mm main with 70 mm trim quadrupoles. The data are plotted at the z values for the 70 mm quadrupoles to ease comparison.

thickness, making this option impractical except possibly with 80 mm magnets. The longitudinal power distribution is shown in Table 4 for 70 mm quadrupoles with 10σ absorbers. Averaged over each element, the power density varies from 3 W/m (Q2a) to 10 W/m (Q01). However, the variation within one magnet can be as large as a factor of 8 (Q1 with uniform diameter absorbers).

4 CONCLUSIONS

Energy deposited in the superconducting magnets is an important issue in the overall design of the LHC IRs. Reducing P_{max} to an acceptable level requires the use of internal absorbers at least 5-6 mm thick. Quadrupoles with $d_{coil}=70$ mm, the current baseline design, are large enough to accommodate such liners and leave a 10σ physical aperture. A larger d_{coil} would allow use of a thicker absorber, greater physical aperture for the same absorber thickness, or possibly cooling the absorber at a higher temperature than the magnet. P_{max} in D1 at the high luminosity IRs is unacceptably large if dipoles of the baseline length are used. However, here the dipoles can be moved farther apart reducing their length by up to a factor of 3 or allowing the use of conventional magnets.

Table 3: Total deposited power (W).

d_{coil} (mm)	70	70	70	70	60	80
Absorber (σ)	none	10	10	8	10	10
			unif			
Quadrupoles	115	82	86	66	98	69
Absorbers	—	61	52	73	37	78
Total	115	143	138	139	135	147
D1	45	26	34	19	24	25
Absorber	—	10	16	15	12	9
Total	45	36	50	34	36	34

Table 4: Total power (W) deposited in each magnet for 70 mm quadrupoles with 10σ absorbers.

	Q1	Q01	Q2a	Q2b	Q3	Q03
Magnet	15	6	13	17	26	4
Absorber	18	10	5	9	13	7
Total	33	16	18	26	39	11

5 ACKNOWLEDGEMENTS

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AMPLITUDE DEPENDENT TUNE SPREAD AND FIELD ERRORS OF SUPERCONDUCTING LOW- β QUADRUPOLES

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Abstract

The nonlinear effects in the low- β insertions are studied on the basis of the amplitude dependent tune spread. Several methods of estimating the tune spread in superconducting low- β quadrupoles are compared. The main feature of these magnets is that the random errors dominate the multipole spectrum. The proposed methods allow an analysis of the final focus layout in the early stages of the insertion design, and point to the critical magnets and the dominant multipole errors.

1 INTRODUCTION

The large β -functions in the low- β insertions which are associated with small values of β^* , chosen for optimising the luminosity, are generally considered a limiting factor of hadron collider performance. High- β conditions inside the final focus quadrupoles are the source of chromatic and amplitude dependent effects. While both are related to the precise layout of the low- β section, the field quality of the low- β quadrupoles drives non-linear resonances, which together with the beam-beam collisions, limit the lifetime of the colliding beams. In order to estimate the relevance of the large amplitude motion, the contribution of each low- β quadrupole should be considered by taking into account the details of the local optics, in particular of the variation of the β -function and of the central orbit inside the quadrupoles, and of the possible asymmetries of the layout around the collision point. It is also important to study the layout in the early stages of the design, and to determine those features of the low- β quadrupoles which contribute mostly to exciting the non-linearities. An important indicator in this respect is the amplitude dependent tune spread.

In this report we propose a method of estimating the amplitude dependent tune spread of a low- β triplet of superconducting quadrupoles, in which the random errors typically dominate the multipole spectrum. These quadrupoles cannot be treated as thin lenses both because of their length and gradient, and because of the displaced beam trajectories arising from finite crossing angles. Furthermore, the optical functions vary considerably over short lengths, comparable to the extension of the quadrupole end field, where the systematic multipole errors are compensated on the average. With modest computational effort, dominant features of the triplet layout can be determined, and tolerance limits on the multipole errors can be derived.

2 AMPLITUDE DEPENDANT TUNE SPREAD IN PERTURBATION THEORY

The long-term particle motion in non-linear magnetic fields is an area of extensive studies. The long term dynamic aperture is correlated to the tune spread of the beam, with the value of around $\Delta Q = 0.015$, chosen for the LHC [1], considered as the upper limit for a hadron collider. The dominant part of the tune spread budget is attributed to the head-on and parasitic beam-beam collisions. While only a small fraction is related to the single particle motion (typically $\Delta Q = 0.005$), mostly due to large amplitude oscillations in the low- β insertions, the amplitude dependent tune spread beyond this level clearly limits the dynamic aperture.

We base our approach of estimating the amplitude dependent tune spread on perturbation theory, where the non-linear terms of the Hamiltonian are treated as perturbations to the well known linear motion [2]. To first order, the tune shift is obtained by averaging the phase independent part of the perturbation around the ring. As the perturbation is directly related to the multipole spectrum of the guide field, customarily represented as :

$$B_y + iB_x = B_0 \sum_{n=1} (b_n + ia_n) \left(\frac{x + iy}{r_0} \right)^{n-1}$$

the contribution of the k -th multipole error to the tune shift may be written in terms of the optics and multipole field errors of each magnet in a given section of the low- β insertion:

$$\delta Q_x = \int \frac{ds}{2\pi\rho} \beta_x \sum_{n=1}^{10} (C_n^{x,b} b_n + C_n^{x,a} a_n) \quad (1)$$

The most important coefficients for the bending plane are :

$$\begin{aligned} C_4^{x,b} &= (1 - \delta)(1.5\Delta^2 + \frac{J_1}{4}) \\ C_5^{x,b} &= (1 - \delta)\Delta_x J_1 \\ C_6^{x,b} &= (1 - \delta)(2.5\Delta^2 J_1 + \frac{J_2}{6}) \end{aligned}$$

where:

$$\begin{aligned} \Delta_x &= x_c + D_x \delta \\ \Delta_y &= y_c + D_y \delta \\ \Delta^2 &= \Delta_x^2 - \Delta_y^2 \\ J_1 &= 3(J_x \beta_x - 2J_y \beta_y) \\ J_2 &= \frac{15}{2}(J_x^2 \beta_x^2 - 6J_x \beta_x J_y \beta_y + 3J_y^2 \beta_y^2) \end{aligned}$$

luminosity insertions [1]. The “Gaussian beam” tune spread is considered a baseline method, but is itself time consuming and inappropriate for examining a large number of different situations. A very precise upper limit is obtained by using the “ 4σ ” method, which is presented in Fig. 1 by its “four point” footprint $((J_x, J_y) = (0, 0), (0, 4), (4, 0), (4, 4))$. The non-zero average tune shift in these two cases is due to the fact that all random errors are considered, including the strong b_3 term which contributes to the average tune shift but not to the tune spread. The tune spread obtained by the “rms” and the “maximum error” methods, are on the contrary centred at the tune shift produced by the small systematic errors. The “rms” tune spread is by a factor of 1.6 larger than the “ 4σ ” value and requires the smallest computational effort. However, we consider the “ 4σ ” method as the most appropriate as it consistently gives good results within short computational time. The “maximum error” case is larger by a factor of 1.7.

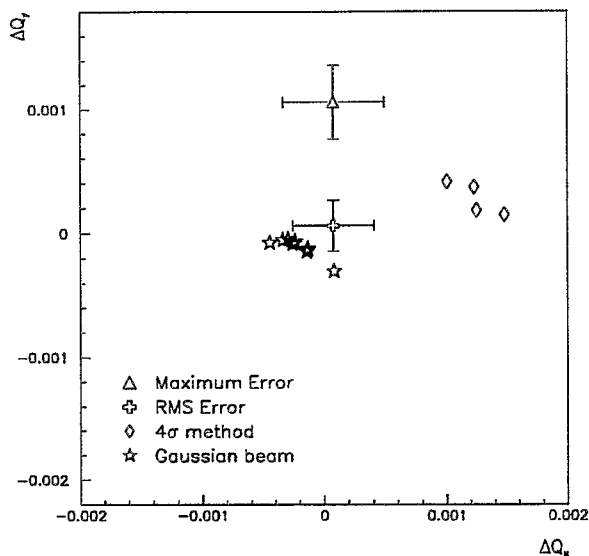


Figure 1: Comparison of the tune spread estimates for the LHC low- β triplet. For clarity, the “maximum error” tune spread has been shifted by 0.001 in ΔQ_y .

4 ANALYSIS OF LOW- β TRIPLET LAYOUTS

The major advantage of a simple method of estimating the amplitude dependent tune spread is that a number of configurations can be studied in the early stages of the insertion design without resorting to time consuming tracking studies. In the case of the LHC low- β quadrupoles [3], the application of this method gave the following results:

- The tune spread of the low- β insertion is dominated by the quadrupoles sitting in regions of peak β -values, which contribute 90% of the total tune spread. The off-centred central orbit ($\leq 4\sigma$) increases the tune spread

by a factor of 2.5. In the LHC, changes in β^* influence the tune spread twice as much as changes of the crossing angle.

- The multipoles which contribute most to the tune spread are the octupole and decapole random, and dodecapole systematic and random errors. The errors of quadrupole ends, in particular the b_6 component of the lead end, contribute to about half of the total tune spread. These errors are systematic, and can be compensated by passive or active methods.
- By choosing the side on which to put magnet connections, the low- β triplet can be made such that the two LHC beams have the same tune spread. The value of the tune spread is reduced from $\Delta Q = 5 \cdot 10^{-4}$ to $3 \cdot 10^{-4}$ by the correct choice of the position of the connection side.
- Assuming the “maximum error” case and four low- β insertions tuned at a β^* of 0.5 m, the LHC tune spread is by a factor 1.6 below the tentative limit of 0.005.

The last result indicates that the dynamic aperture of the LHC at top energy should not be limited by the random multipoles. This is confirmed by recent tracking studies [4], which give a dynamic aperture of 10σ in physics conditions, on the edge of the “good field region” of the 70 mm aperture low- β quadrupoles.

5 CONCLUSIONS

We have compared several methods of estimating the amplitude dependent tune spread in superconducting low- β quadrupoles. The approach addresses the main feature of the multipole spectrum of these magnets, i.e. the dominance of the random errors. Furthermore, the details of the local optics, in particular the rapid variation of the β -function and central orbit deviation, and the asymmetries of the insertion optics are taken into account. The method allows to determine the critical quadrupoles and multipole errors, to investigate the role of magnet orientation and connections, and to set limits on multipole error tolerances.

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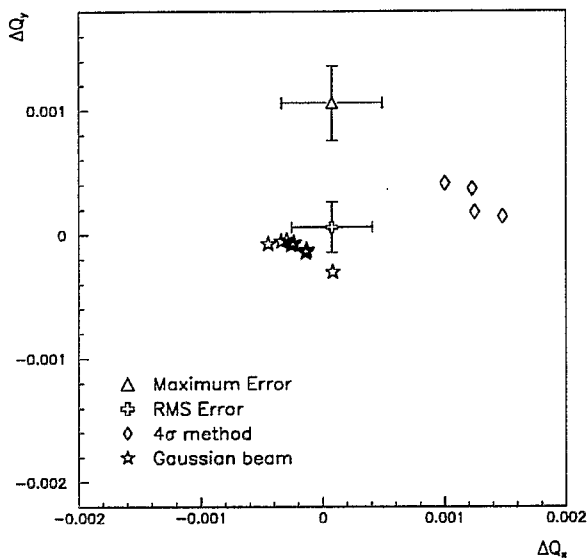


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SUPERCONDUCTING SEXTUPOLE CORRECTOR MAGNET FOR THE LHC MAIN DIPOLES

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Abstract

Each LHC main dipole will be equipped with small sextupole corrector magnets with a field strength of $1970 \text{ x}^2 \text{ T/m}^2$ and a magnetic length of 100 mm designed to correct the sextupole field errors. The paper presents a cosine- θ type of design where much emphasis has been put on the cost reduction because these magnets have to be made in a large series of some 2500 pieces. We describe the design of a two layer coil which can be wound automatically. The winding starts in the middle of the wire with the only joggle, the layer jump, which is housed in a corresponding groove in the end of the central island. The two layers are wound simultaneously turning in opposite directions to find their position without the need of local tooling. The coil ends are closely packed and need no end spacers. The 18 pole perturbation introduced by the ends is corrected by the position of the coil block in the straight part. The yoke is made of iron laminations of the "Scissors type" which transmit the pre-stress from the outer aluminium shrink ring to the coil. This allows the iron to be close to the coil for field enhancement and also boosts the pre-stress in the coil due to the cool down contractions. The paper describes the experience with the magnet construction and gives the first test results.

1 INTRODUCTION

1.1 General

In the LHC ring each bore of the main bending dipole is equipped with a superconducting sextupole (MCS) and decapole (MCD) corrector magnet to correct systematic sextupole and decapole field errors in the dipole field. Sextupole correctors are placed at one end of the dipoles and decapoles at the opposite end. In the first generation sextupole corrector design, coils were wound around a copper central island in eight layers each consisting of fourteen turns [1]. This coil design proved to be laborious and expensive to manufacture, since after each turn the cable must be twisted to be able to jump to the next turn, and after each layer to do a layer jump. During the winding each turn must be clamped by special tooling to prevent unwinding of the coil. The new design drastically reduces the time to

make a magnet. In this paper will be described only the design and the manufacturing of the sextupole magnet however similar design principles will be used for the decapole and octupole magnets. Six of this type of sextupole and decapole and four octupole magnets will be built at CERN by beginning of April 1997.

1.2 Design Criteria

The design criteria of the magnets are given by the beam dynamics and the available space in the dipole ends. The main parameters are given in the Table 1. In the LHC ring the corrector power supplies will be of 600A type.

Nominal strength	$1970 \text{ x}^2 \text{ T/m}^2$
Magnetic / Overall length	104 / 150 mm
Nominal current	625 A
Peak field in the coil 3D	2.12 T
Self inductance	0.672 mH
Inner / Outer diameter of the coil	58 / 61 mm
Inner / Outer diameter of the yoke	66 / 90 mm
Inner / Outer diameter of shrink ring	89.88 / 100 mm
Cable dimensions bare	$1.13 \times 0.606 \text{ mm}^2$
Cu/Sc ratio	1.6
Insulation thickness (PVA)	0.06 mm
Turns per coil	2x13
Working temperature	1.9 K
% on the load line 1.9 / 4.2 K	62.9 / 47.5

Table 1. Magnet parameters

2 MAGNETIC DESIGN

2.1 General

The magnetic design aimed for an economic two layer coil. In the design a combination of 2D and 3D modelling has been used. 3D models have been created using the ROXIE program, which uses Biot-Savart on the line currents [2]. The surrounding iron yoke can only be modelled with linear or infinite permeability. The induction of the iron remains well below the saturation level up to nominal operation level, which was checked with the POISSON program in 2D. ROXIE is not able to calculate 3D peak fields in a coil surrounded by iron, the 2D peak field in the coil is calculated with and without iron and the same ratio is assumed to be valid in 3D.

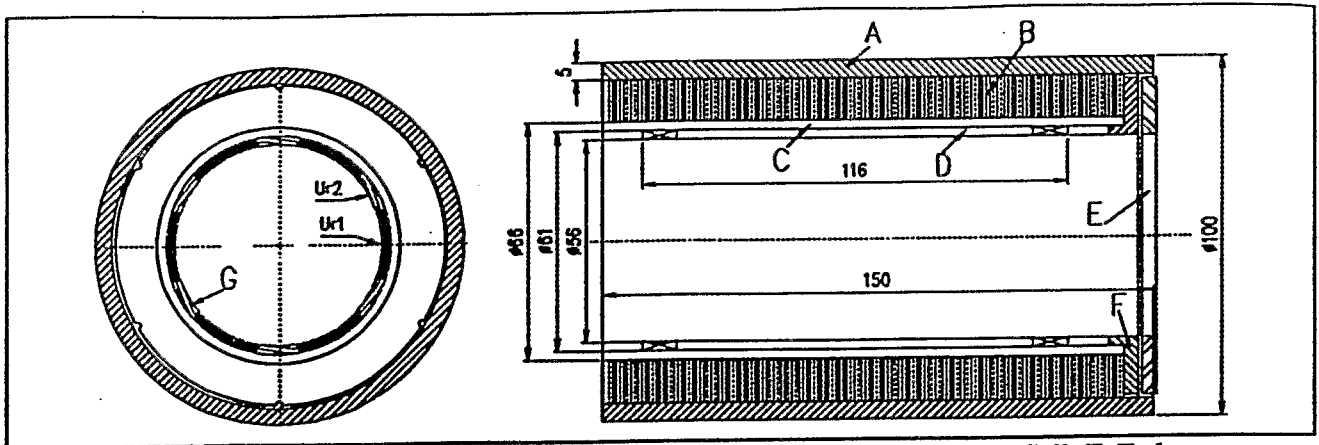


Figure 1. MCS magnet A: Shrinking ring B: Scissor laminations C: Glass fibre bandage D: Coils E: End cover F: End flange for the connections G: Layer jump in one end of the central island

2.2 Coil optimisation

By bringing the surrounding yoke close to the coils for field enhancement, it was possible to reach the required field strength by using a two layer coil. The iron boosts the field by 48.5 %. The coil ends are designed without end spacers to allow for automatic winding. The eighteen pole perturbation induced by the coil ends was compensated by the position of the coil block in the straight part (Table 2). An existing wire was chosen because it's availability. In the LHC ring 154 of these magnets are connected in series and in case of a quench, most of the energy of the 153 other magnets is absorbed in the quenching magnet. In order to survive a quench, for the final design the amount of the copper in the wire will be increased by 40 %, i.e. the Cu/Sc ratio increases from 1.6 to 2.25 [3].

B3 [Tm]	B9/ B3 [$\times 10^{-4}$]	B15/ B3 [$\times 10^{-4}$]
-0.0206249	0.01978	0.000417

Table 2. Integrated harmonics at radius of 10 mm

3 MECHANICAL DESIGN

3.1 Central island material

Glass fibre reinforced epoxy (G11) was chosen for central islands and end spacers material, since it is integrated contraction factor 260×10^{-5} , in the longitudinal direction between 293 and 4.2 K, is close to that of the cables 295×10^{-5} . In addition to which no ground insulation is needed. The end filling pieces and central islands were machined from tubes of this material.

3.2 Scissor laminations

The iron yoke is composed of so called scissor laminations, which transmit the pre-stress from the outer aluminium shrink ring into the coils [4]. Two slightly eccentric laminations, which press from opposite directions on the glass fibre bandage around the coils, form a pair

of "scissors". Over the length of the magnet, successive pairs of laminations are rotated by 60° . As the lamination thickness is 1 mm, every 6 mm pressure is applied in the same angular position.

3.3 FEM-model

A 2D FEM-model has been created, to define stresses in the structure. A 30° sector of the magnet has been modelled using the ANSYS code. The eccentric contact of the scissor laminations to the aluminium shrink ring and the fibre glass bandage has been simulated by using contact elements. Table 5 presents the calculated azimuthal stresses in the coil and in the aluminium cylinder during different load steps.

	293 K	4.2 K	4.2 K (1000A)
σ coil max.	-21	-27	-22
σ coil min.	-49	-68	-74
σ cyl max.	56	107	108
σ cyl min.	-7	21	28
Rad disp Ur1	-0.020	-0.139	-0.133
Rad disp Ur2	-0.0334	-0.168	-0.166

Table 3. Azim. stresses in coil and shrink ring [MPa] and radial displacements of the coil [mm] (Fig 1). Radial interference between yoke and shrink ring is 0.06 mm

4 FABRICATION

4.1 Coil winding

The two layer coil (double pancake) design allows to start winding in the middle of the coil with the only joggle, the layer jump, which has been housed in a groove in the end of the central island. The two layers are wound simultaneously in opposite directions to find their position automatically not needing the clamps normally used to keep the wire in it's plane. The wire is continuously wetted during the winding process by epoxy (Fig. 2). To produce compact coils a constant winding tension of 35 MPa was maintained in the metal

cross section of the cable. Although the wire is thin, when bent over the end of the central post, it has a tendency to bend away from the central post in the straight part, due to its bending stiffness. If the wires are then pushed towards the central island after the whole coil is wound, most of the winding tension will be lost. The situation can be improved, by counter bending the wire around a small pulley, which deforms the wire so that it naturally pushes itself towards the central post. In addition to counter bending, before winding the central island is pre-compressed in the longitudinal direction by 0.1 mm. This is done by applying a load of 350 N the ends, and then locking the central island in place by screws. After the winding the locking screws are loosened so that the pre-stressed central island and the surrounding coil find an equilibrium. After the winding the end fillers are positioned and the coil is clamped in the curing mould, which is doweled through the central island on the winding mandrel. Finally the coil is heated in the oven in order to cure the epoxy. Winding and impregnation take approximately two hours per coil using a simple hand driven machine.

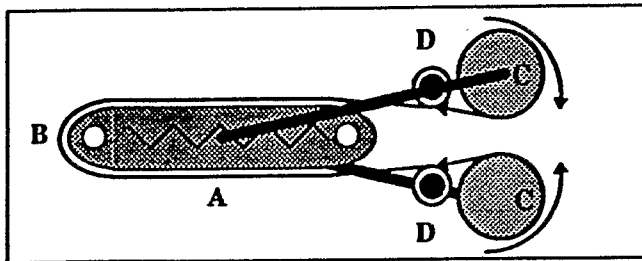


Figure 2. Winding A: Pre-compressed central island B: Layer jump C: Wire reels D: Counter bending pulleys

4.2 Assembly of the magnet

The impregnated coils are positioned by dowel pins around an assembly mandrel and a glass fibre-epoxy end flange is placed at one end of the coils for the series connections. A prepreg glass fibre cloth is wrapped around the coils whereupon it is clamped and cured in an oven. After turning the bandage, holes are drilled in the positions of the dowel pins, and these pins are pushed into the hollow mandrel which is then pulled out. The series connections are made by soldering the wires side by side in grooves on the end flange over a length of 35 mm. The connections are kept in place by a glass fibre cover plate which is screwed on the end flange. The yoke laminations are assembled around the glass fibre bandage as described in the section 3. To facilitate the slipping on of the heated aluminium shrink ring over the laminations, a slightly conical Teflon hat is mounted at the end of the magnet.

5 TEST RESULTS

The training test of the magnet was done at 4.2 K in Madrid, by CEDEX. The first quench of the magnet was

at 734 A which is well above the nominal operation level (Fig. 3). After a heat cycle to room temperature the magnet showed very little re-training i.e. the first quench occurred at 934 A. The current was measured by clamp meters around the bus bars with an accuracy of ± 20 A. Voltage taps were installed over the poles in order to record the quench voltages. Also the current decay was measured in order to calculate the MIT's and hence define the hot spot temperatures which proved to be at maximum around 130 K.

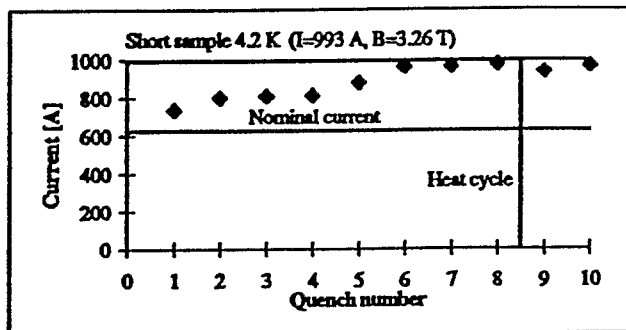


Figure 3. Training of the magnet at 4.2 K.

6 CONCLUSIONS

A second generation cosine- θ type of superconducting sextupole corrector magnet (MCS) has been designed and built at CERN. A training test has been done at 4.2 K by CEDEX. By using scissor laminations, it has been possible to bring the iron close to the coils for maximum field enhancement. Since the iron boosts the field by 49%, the required field strength of 0.197 T at 10 mm radius, is obtained by two layer coil. The aim of the two layer coil design has been to allow for an easy automatization of the winding process.

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